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# The high energy X-ray probe (HEX-P): the most powerful jets through the lens of a superb X-ray eye

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A fraction of the active supermassive black holes at the centers of galaxies in our Universe are capable of launching extreme kiloparsec-long relativistic jets. These jets are known multiband (radio to  $\gamma$ -ray) and multimessenger (neutrino) emitters, and some of them have been monitored over decades at all accessible wavelengths. However, several open questions remain unanswered about the processes powering these highly energetic phenomena. These jets intrinsically produce soft-to-hard X-ray emission that extends from  $E > 0.1$  keV up to  $E > 100$  keV, and simultaneous broadband X-ray coverage, combined with excellent timing and imaging capabilities, is required to uncover the physics of jets. Indeed, truly simultaneous soft-to-hard X-ray coverage, in synergy with current and upcoming high-energy facilities (such as IXPE, COSI, CTAO, etc.) and neutrino detectors (e.g., IceCube), would enable us to disentangle the particle population responsible for the high-energy radiation from these jets. A sensitive hard X-ray survey ( $F_{20-80 \text{ keV}} < 10^{-15} \text{ erg cm}^{-2} \text{ s}^{-1}$ ) could unveil the bulk of their population in the early Universe. Acceleration and radiative processes responsible for the majority of their X-ray emission would be

pinned down by microsecond timing capabilities at both soft and hard X-rays. Furthermore, imaging jet structures for the first time in the hard X-ray regime could unravel the origin of their high-energy emission. The proposed Probe-class mission concept High Energy X-ray Probe (HEX-P) combines all these required capabilities, making it the crucial next-generation X-ray telescope in the multi-messenger, time-domain era. HEX-P will be the ideal mission to unravel the science behind the most powerful accelerators in the Universe.

#### KEYWORDS

blazar, supermassive black hole, jet, high energy astrophysical phenomena, X-ray mission, time domain astrophysics, multimessenger astronomy

## 1 Introduction

The high-energy ( $E > 1$  keV) extragalactic sky is filled with accreting supermassive black holes (known as active galactic nuclei, AGN, Ajello et al., 2008; Ajello et al., 2009; Gilli, 2013; Ueda et al., 2014; Cappelluti et al., 2017; Ananna et al., 2020). These monstrous objects, reaching masses of billions or more Suns, are feeding from the gas that surrounds them, emitting copious amounts of radiation. A fraction of these black hole systems also powers relativistic jets from their core. When the jet is closely aligned to the Earth's line of sight (viewing angles,  $\theta_v < 5^\circ - 10^\circ$ ), these objects are known as blazars; otherwise they are classified as misaligned AGNs or radio galaxies (e.g., Urry and Padovani, 1995; Hovatta et al., 2009; Pushkarev et al., 2017). Although continuous multi-band monitoring of these sources has been ongoing for more than 2 decades (e.g., Urry et al., 1997; Kataoka et al., 1999; Komossa et al., 2021; Weaver et al., 2022; de Jaeger et al., 2023; Otero-Santos et al., 2023), many open questions remain regarding the nature of these extreme jets.

A unique mission such as the High Energy X-ray Probe (HEX-P; Madsen et al., 2023), with its broadband X-ray coverage (0.2–80 keV), unprecedented sensitivity ( $F_{0.5-2 \text{ keV}} < 10^{-16} \text{ erg cm}^{-2} \text{ s}^{-1}$  and  $F_{3-24 \text{ keV}} < 10^{-15} \text{ erg cm}^{-2} \text{ s}^{-1}$  in a 1 Ms exposure, Civano et al., 2023), good timing resolution ( $\Delta t \sim 2 \mu\text{s}$ ) across a wide range of energies and angular resolution in the hard X-ray band 3–4 times better than NuSTAR, will enable us to answer several open science questions regarding blazars. HEX-P will allow us, for the first time, to: (1) spatially resolve extragalactic nearby jets in hard X-rays; (2) detect spectral signatures (e.g., spectral curvatures, exponential cut-off) that will uniquely point us to the preferred high-energy emission mechanisms in jets, and, by extension, their composition; and (3) in synergy with the current and upcoming landscape of experiments (e.g., COSI, IXPE, CTAO, IceCube Gen 2), reveal the preferred particle acceleration and emission mechanism of the brightest objects in the Universe, fulfilling the promises of multi-wavelength and multi-messenger science. In this paper, we explore some of the open issues in blazar science, and more broadly jetted AGN, and how HEX-P will revolutionize our understanding of these sources.

The peculiar jet orientation in blazars allows for extreme relativistic beaming, opening a window for studies of particle acceleration mechanisms and jet structure (Madau et al., 1987; Dermer, 1995; Dondi and Ghisellini, 1995). These sources are

extremely bright multi-wavelength emitters (from radio up to TeV energies; see Massaro et al., 2008; Madejski and Sikora, 2016; Blandford et al., 2019), reaching apparent bolometric luminosities of  $> 10^{48} \text{ erg s}^{-1}$  (e.g., Massaro et al., 2015; Ghisellini et al., 2017; Paliya et al., 2019b; Ajello et al., 2020). Blazars are also likely multi-messenger sources, as they are able to accelerate protons and produce both high-energy neutrinos and  $\gamma$ -rays (e.g., IceCube Collaboration et al., 2018). Phenomenologically speaking, blazars have been divided into several subcategories. Objects with broad and strong emission lines (equivalent width,  $\text{EW} > 5 \text{ \AA}$ ) in optical spectra are classified as Flat Spectrum Radio Quasars (FSRQs). When these lines are weak or absent, these sources are known as BL Lacertae objects (BL Lacs, e.g., Schmidt, 1968; Stein et al., 1976). Their electromagnetic broadband spectral energy distribution (SED) is characterized by two broad “humps”, a low-energy hump peaking in the IR–X-ray range, and a high-energy hump peaking in the MeV– $\gamma$ -ray range. It is now well understood that most of the radiation from radio to optical/X-ray frequencies of these jets is produced by synchrotron radiation of electrons moving at relativistic speeds in magnetic fields in the jets (e.g., Pian et al., 1998; Padovani et al., 2002). This scenario is further strengthened by the high polarization degree measured in these sources at radio to optical frequencies, and now also in X-rays for the two nearest blazars (e.g., Jorstad et al., 2007; Blinov et al., 2021; Liodakis et al., 2022; Zhang et al., 2023; Zobnina et al., 2023). Based on the location of their synchrotron spectral peak (Abdo et al., 2010a), these sources have later been classified as low- (LSP,  $\nu_{\text{synch,peak}} < 10^{14.5} \text{ Hz}$ ), intermediate- (ISP,  $10^{14.5} \text{ Hz} < \nu_{\text{synch,peak}} < 10^{15} \text{ Hz}$ ) and high-synchrotron peaked (HSP,  $\nu_{\text{synch,peak}} > 10^{15} \text{ Hz}$ ) sources. FSRQs usually belong to the LSP class, being the most luminous of the class, detected up to the highest redshifts ( $z > 4$ ). BL Lacs instead span the whole range of classifications, are less luminous, and are typically found at  $z < 1$  (Shen et al., 2011; Ajello et al., 2014; Massaro et al., 2016; Peña-Herazo et al., 2021; Middei et al., 2022; Rajagopal et al., 2023). However, the physical origin of the high-energy part of the SED is still a matter of debate. In the leptonic scenario (e.g., Ghisellini and Maraschi, 1989; Dermer et al., 1992; Sikora et al., 1994), the X-ray to very high energy (VHE,  $E > 100 \text{ GeV}$ )  $\gamma$ -ray emission is explained by inverse Compton scattering by the same population of electrons that produces the synchrotron emission (e.g., Beckmann et al., 2002; Tagliaferri et al., 2002). These electrons scatter the photon fields either within the jet structure (referred to as synchrotron self-Compton; SSC) or those external to

the jet (for example, the accretion disk, the broad-line region, or an infrared-emitting dust torus; this emission is referred to as external Compton; EC). As an alternative to the leptonic inverse-Compton interpretation, the high-energy emission could also be produced by ultra-relativistic protons through proton synchrotron emission and emission from electromagnetic cascades initiated by photon production and Bethe-Heitler pair production processes (e.g., Mannheim, 1993; Mücke et al., 2003).

Clues towards the jets' particle composition lie in the simultaneous characterization of the X-ray spectrum from soft X-rays ( $E \sim 1$  keV) to hard ( $E \geq 10$  keV) X-rays (e.g., Maraschi, 1999). As pioneered by missions such as BeppoSAX (Boella et al., 1997) and Suzaku (Koyama et al., 2007; Takahashi et al., 2007), and later confirmed by blazar studies using combined observations from soft (e.g., Swift/XRT, XMM-Newton, Chandra) and hard (Swift/BAT, NuSTAR) X-ray instruments, simultaneous coverage at these energies is crucial for several blazar classes: it allows characterization of the synchrotron peak of the most extreme BL Lac sources (e.g., Costamante et al., 2001; Reimer et al., 2008), pinpointing the high-energy peak of the most luminous MeV-peaked blazars (e.g., Ghisellini et al., 1999; Tagliaferri et al., 2015; Marcotulli et al., 2020), and characterization of the transition between the low- and high-energy humps of the SED (e.g., Abdo et al., 2010b; Böttcher et al., 2013), the shape of which is fundamental to unveil the nature of the particle population in the jet. The superior broadband sensitivity, imaging capabilities, and lower backgrounds, of HEX-P would allow us to study in detail a larger number of blazars of different types, sampling their SED at flux levels order of magnitude lower than previously achievable by past broadband X-ray missions. Due to its coverage of the full X-ray energy range, HEX-P will be complementary for blazar science to many other recent or near-future experiments, such as the recently launched Imaging X-ray Polarimetry Explorer (IXPE; 2–8 keV, Weisskopf, 2022), the approved Compton Spectrometer and Imager (COSI; Tomsick et al., 2019), and the order-of-magnitude improvement in sensitivity to high-energy  $\gamma$ -rays above 100 GeV provided by the upcoming Cherenkov Telescope Array Observatory (CTAO; Cherenkov Telescope Array Consortium et al., 2019).

Another observational multi-band characteristic of blazars is their variability, with flaring timescales ranging from months down to minutes, flux enhancements of up to several orders of magnitude, and cross-correlated variability patterns (or lack thereof) observed between different energy bands. To explain this behavior, several particle acceleration mechanisms have been proposed to act in these jets, related to magnetic reconnection occurring inside the jet (e.g., Giannios, 2013; Christie et al., 2019; Zhang H. et al., 2022) for the fastest time scales observed, or mildly relativistic shocks (Rees, 1978; Marscher and Gear, 1985) to describe variability on longer time scales. The most rapid variations, which are shorter than 10 min, have been observed at  $\gamma$ -ray energies for a handful of sources (Albert et al., 2007; Aleksić et al., 2014a; Ackermann et al., 2016a; Aharonian et al., 2017; Meyer et al., 2019). In the X-ray energy range from  $\sim 0.5$ –20 keV, variability on similar time scales has also been found for several sources (e.g., Catanese and Sambruna, 2000; Zhu et al., 2018; Zhang Z. et al., 2019; Zhang et al., 2021b; Chatterjee et al., 2021; Mohorian et al., 2022). At optical and near-infrared wavelengths, rapid variability has also been detected within a day or shorter (e.g., Gopal-Krishna et al., 2011; Sandrinelli et al.,

2014; Agarwal et al., 2015). Typically, those rapid variations occur when a blazar is in a bright state. Often, spectral changes in the high-energy regime are correlated with an increase in luminosity, where the brighter emission leads to a harder spectrum in the X-ray band for HSP sources (e.g., Bhatta et al., 2018; Wang et al., 2018; Giommi et al., 2021), and in the  $\gamma$ -ray band for LSP sources (e.g., Abdo et al., 2010c; Williamson et al., 2014). Long-term variability is usually also present across all wavelengths (e.g., Mao and Zhang, 2016; Gaur et al., 2019; Rajput et al., 2020; Safna et al., 2020; Mundo and Mushotzky, 2023). However, correlation studies between several energy bands are still inconclusive in diagnosing the production sites of flares (e.g., single zone versus multi-zone scenarios), the mechanisms of particle acceleration, and the particle species responsible for the flares in different wavelengths. Strictly simultaneous multi-wavelength observations and monitoring of flaring blazars must be performed to constrain jet physics. In particular, simultaneously characterizing the transition between the soft and hard X-rays in blazars belonging to different subclasses will provide unprecedented insight into these acceleration processes.

Flaring can lead to a deeper understanding of the physics in AGN jets in terms of particle acceleration mechanisms and dissipative locations in the jet. Additionally, together with neutrino observations, flaring can be pivotal for understanding the particle content responsible for the photon emission. The association of the  $\gamma$ -ray flaring blazar TXS 0506 + 056 with the IceCube neutrino event IceCube-170922A was the first, and so far most convincing, evidence that blazars can accelerate protons to very high energies (IceCube Collaboration et al., 2018). However, the lack of simultaneous broadband coverage from X-ray to MeV  $\gamma$ -rays has prevented us from constraining the relevant models sufficiently well to probe the transition between the synchrotron and high-energy emission part of the spectrum, where signatures of electromagnetic cascades associated with the neutrino-producing photo-hadronic interactions should emerge (Murase et al., 2018). In particular, studies of a previous neutrino flare potentially linked to TXS 0506 + 056 show that the complete lack of sensitive simultaneous broad X-ray spectral and temporal coverage significantly hampered any conclusion on the emission mechanism (e.g., Reimer et al., 2019; Petropoulou et al., 2020; Zhang B. T. et al., 2020). Simultaneous hard X-ray ( $E > 10$  keV) observations would be able to probe the transition regime between the low-energy and high-energy peak, and the evolution of the spectral shape at different epochs (pre/post neutrino event), hence removing degeneracies on parameters in the tested models. Moreover, in the context of a one-zone hybrid lepto-hadronic scenario, Petropoulou et al. (2020) found that the maximal neutrino flux is strongly correlated with the X-ray flux in the case of the source, which underscores the importance of X-ray measurements for constraining the blazar neutrino output. Simultaneous broadband coverage of neutrino alerts in the direction of blazar sources could be the tipping point for jet models as well as confirming AGN as one of the sources of high-energy cosmic rays.

In terms of population studies, there appears to exist an anti-correlation between the synchrotron peak position and the intrinsic bolometric luminosity of blazars, the so-called blazar sequence (see Fossati et al., 1998; Ghisellini et al., 1998; Ghisellini et al., 2017). If real, this would imply that more luminous sources (found at the highest redshifts) have higher accretion rates ( $L_{\text{disk}}/L_{\text{Edd}} \geq 10^{-2}$ ), hence host more radiatively efficient accretion

disks than less luminous (and lower redshift) ones. Opponents of the sequence, however, argue that this is a mere selection effect (e.g., Nieppola et al., 2008; Giommi et al., 2012; Keenan et al., 2021). A mission with high sensitivity in the broad X-ray band such as HEX-P will be key to narrowing down this issue. On the one hand, HEX-P's superior sensitivity will allow us to detect tens to hundreds of sources beyond  $z \geq 4$ , as predicted by current X-ray blazar evolutionary models (the so-called MeV blazars; see, Marcotulli et al., 2022). Multi-band coverage will enable us to pinpoint the location of the high-energy peak and the intrinsic jet luminosity (probing the low-synchrotron-peak-frequency end of the blazar sequence). On the other hand, HEX-P has the potential to discover a handful of MeV-peaked BL Lac sources in the local Universe. Predicted by the blazar sequence, these jets may have been missed by current surveys due to being too faint or overwhelmed by galaxy emission. If found by HEX-P, they would confirm that the blazar sequence is real and clarify its physical origin.

Finally, jets of active galaxies have been observed to extend up to hundreds of kiloparsecs into the intergalactic medium. They are spatially best resolved in the radio band, using very-long-baseline interferometry (VLBI) observations, where low-energy synchrotron emission can be observed along the entire jet (e.g., Linfield, 1981; Boccardi et al., 2016). Moreover, many jets also show extended components (straight jets or hotspots along the jet flow) in various other energy bands including the X-ray domain (Chartas et al., 2000; Harris and Krawczynski, 2006; Stawarz et al., 2013; Marshall et al., 2018; Meyer et al., 2018), even up to the TeV energy band (e.g., Abramowski et al., 2012; H. E. S. S. Collaboration et al., 2020; Archer et al., 2020). The origin of the X-ray (and high-energy) emission from extended jets is certainly non-thermal (see e.g., Harris and Krawczynski, 2002; Harris and Krawczynski, 2006). However, it is unclear whether synchrotron radiation or inverse-Compton emission dominates. The former process seems to well describe the emission arising from low-power jets while for high-power radio galaxies and QSOs the dominant mechanism could be inverse Compton, in most cases having seed photons provided by the CMB radiation (e.g., Harris and Grindlay, 1979; Stawarz et al., 2003; Atoyan and Dermer, 2004; Zacharias and Wagner, 2016; Meyer et al., 2017; Roychowdhury et al., 2022). Observations with a mission such as HEX-P will be crucial for resolving, for the very first time in the hard X-ray band beyond 10 keV, nearby bright jets and disentangling these scenarios.

The paper is organized as follows: Section 2 describes the HEX-P mission design; Section 3 details the setup for the simulations performed throughout the paper; Section 4 addresses the contribution of HEX-P to the multi-messenger and time domain astrophysics of blazars; Section 5 highlights the prospects for different blazar classes; Section 6 showcases the importance of HEX-P in resolving jets in the hard X-ray band; and Section 7 summarizes the main science cases that HEX-P will be able to address with respect to the most powerful jets in the Universe.

## 2 Mission design

HEX-P (Madsen et al., 2023) is a Probe-class mission concept that offers sensitive broad-band coverage (0.2 – 80 keV) of the X-ray spectrum with exceptional spectral, timing, and angular capabilities.

It features two high-energy telescopes (HETs) that focus hard X-rays, and one low-energy telescope (LET) that focuses lower energy X-rays.

The LET consists of a segmented mirror assembly coated with iridium on monocrystalline silicon that achieves a half-power diameter of 3.5", and a low-energy DEPFET detector, of the same type as the Wide Field Imager (WFI; Meidinger et al., 2020) onboard Athena (Nandra et al., 2013). It has  $512 \times 512$  pixels that cover a field of view of  $11.3' \times 11.3'$ . It has an effective passband of 0.2 – 25 keV, and a full frame readout time of 2 ms, which can be operated in a 128 and 64 channel window mode for higher count-rates to mitigate pile-up and faster readout. Pile-up effects remain below an acceptable limit of  $\sim 1\%$  for a flux up to  $\sim 100$  mCrab ( $2.4 \times 10^{-9}$  erg cm $^{-2}$  s $^{-1}$ , in the 2 – 10 keV band) in the smallest window configuration. Excising the core of the PSF, a common practice in X-ray astronomy, will allow for observations of brighter sources, with a typical loss of up to  $\sim 60\%$  of the total photon counts.

The HET consists of two co-aligned telescopes and detector modules. The optics are made of Ni-electroformed full shell mirror substrates, leveraging the experience of XMM-Newton (Jansen et al., 2001), and coated with Pt/C and W/Si multilayers for an effective passband of 2 – 80 keV, achieving a half-power diameter 3 – 4 times better than that of NuSTAR. The high-energy detectors are of the same type as flown on NuSTAR (Harrison et al., 2013), and they consist of 16 CZT sensors per focal plane, tiled  $4 \times 4$ , for a total of  $128 \times 128$  pixels spanning a field of view slightly larger than for the LET, of  $13.4' \times 13.4'$ .

## 3 Simulations

All the simulations presented here were produced with a set of response files that represent the observatory performance based on current best estimates as of Spring 2023 (see Madsen et al., 2023). The effective area is derived from ray-tracing calculations for the mirror design including obscuration by all known structures. The detector responses are based on simulations performed by the respective hardware groups, with an optical blocking filter for the LET and a Beryllium window and thermal insulation for the HET. The LET background was derived from a GEANT4 simulation (Eraerds et al., 2021) of the WFI instrument, and the HET background was derived from a GEANT4 simulation of the NuSTAR instrument; both backgrounds assume HEX-P's Lagrange point L1 orbit.

The spectral simulations were performed in XSPEC (Arnaud, 1996) via the "fakeit" command with V07<sup>1</sup> response files, selecting an 80% PSF correction. The data are assumed to be extracted from a 10 arcsec region, with the anticipated background for the telescope being positioned at L1. The images and light curves were simulated using the SIXTE toolkit (Dauser et al., 2019), using the same response files as the spectral simulation.

The differential sensitivity curves shown throughout the paper are obtained by calculating for each energy bin (LET:  $\Delta E = 0.01$  keV, HET:  $\Delta E = 0.04$  keV) the minimum detectable flux for a signal-to-noise ratio (SNR) of 5 in exposures of 100 ks and 1 Ms.

<sup>1</sup> <https://hexp.org/response-files/>



## 4 Time domain and multi-messenger astrophysics

### 4.1 Leptonic vs. hadronic populations in blazar jets

#### 4.1.1 Probing the relativistic leptonic population in blazar jets

As discussed in Section 1, the location of the two peaks of the blazar's broadband SED generally depends on its apparent luminosity (e.g., Ghisellini et al., 2017). For HSP sources, the flux in the spectral regime corresponding to the X-ray energy (0.1–100 keV) is typically lower than that of the peaks. Thus in an SED plot ( $\log(E \times F(E))$  vs.  $\log(E)$ ), it appears as a “valley” (although note that the peak of the synchrotron component of HSPs and some LSPs and ISPs during flares can extend to the X-ray bands). Figure 1 shows the multi-epoch broadband SED for the HSP blazar PKS 2155–304, highlighting the data taken by Madejski et al. (2016) during a non-flaring period of the object, illustrating this “valley”. Recent measurements by the IXPE satellite have revealed that the soft X-ray (0.2–8 keV) emission of such sources is also polarized (Liodakis et al., 2022), confirming the synchrotron process as the dominant emission mechanism for the entire low-energy “hump” in blazar SEDs. This suggests that hard X-ray studies probe emission by the most energetic particles in the jet (e.g., Ravasio et al., 2003; Massaro et al., 2004; Massaro et al., 2006), since in the synchrotron process, the radiated photon energy goes as the square of the energy of the radiating particle. As the highest energy particles in the distribution radiate more efficiently, we expect to see a cut-off informing us about the maximum Lorentz factor of the particles in the jet ( $\gamma_{\text{max}}$ ) and, in general, constraining the spectral evolution of the electron distribution (e.g., Hota et al., 2021). However, given the weak  $E \times F(E)$  flux in the X-ray band and a soft X-ray spectrum ( $\Gamma_X > 2$ , Middei et al., 2022), current instruments cannot accurately measure the spectral cutoff of the synchrotron tail during non-flaring states of ISP/HSP blazars. For example, past 20 keV NuSTAR data are dominated by the instrumental background, and hence cannot inform us about the underlying energy distribution beyond that energy. With broadband coverage in both soft and hard X-ray energies, HEX-P will be able to accurately measure the spectral shape of the synchrotron tail, setting important constraints on the acceleration process of the radiating particles. Moreover, in the leptonic scenario, VHE  $\gamma$ -ray photons are produced by the same electrons that generate the X-ray synchrotron emission. Thus, HEX-P will probe the true electron distribution generating (by the inverse Compton process) the intrinsic VHE spectrum of HSPs, as well as flaring LSPs and ISPs that extend to X-rays. These constraints on the shape of the particle population are crucial as VHE photons from blazars are attenuated by  $\gamma$ - $\gamma$  interactions with low-energy photons associated with the extragalactic background light (EBL). In synergy with CTAO, this can provide unprecedentedly strong constraints on the EBL (e.g., Desai et al., 2019).

To see the improvement of a mission such as HEX-P for this scenario, we simulated a 100 ks observation of a broken power-law X-ray spectrum of PKS 2155–304 at the same flux level as the observation shown in Madejski et al. (2016),  $F_{2-10 \text{ keV}} = 1.3 \times 10^{-11} \text{ erg cm}^{-2} \text{ s}^{-1}$ . The spectral index is  $\Gamma_{0.2-20 \text{ keV}} = 2.74$ , the break occurs at  $E = 20 \text{ keV}$

after which the spectrum steepens to  $\Gamma_{20-50 \text{ keV}} = 4.0$ . The simulated spectrum and model residuals for a power-law (PL) and broken powerlaw model (BPL) are shown in Figure 2. From these plots things to notice are that: (i) the low background of the HET would enable us to reach  $E \sim 40 \text{ keV}$ ; and (ii) it would ascertain the presence of such break above  $E > 20 \text{ keV}$  at a  $\sim 5\sigma$  level. The high background of NuSTAR would not enable us to go past  $E \sim 20 \text{ keV}$ , even with a 1 Ms exposure. Moreover, we checked that for a 1 Ms HEX-P exposure, we would be able to detect the same break at a  $2.5\sigma$  confidence level for a source with  $F_{2-10 \text{ keV}} = 9 \times 10^{-13} \text{ erg cm}^{-2} \text{ s}^{-1}$  ( $\sim 92\%$  of ISP/HSP blazars with  $\Gamma_X > 2$  in Middei et al., 2022 have  $F_{2-10 \text{ keV}} > 9 \times 10^{-13} \text{ erg cm}^{-2} \text{ s}^{-1}$ ).

In the case of LSP sources, the X-ray band already probes the onset of the high-energy SED “hump”. In the context of synchro-Compton models for relativistic jets, this part of the spectrum is produced by the lowest-energy part of the distribution of radiating particles. In these sources, therefore, HEX-P will probe the low-energy tail of the particle distribution, possibly revealing the presence of low-energy cutoffs thanks to the long arm provided by broad-band sensitivity out to 80 keV in a single observation (e.g., Fabian et al., 2001; Tavecchio et al., 2007; Sbarrato et al., 2016). The broad-band capability of HEX-P is also important in ISPs sources, where the X-ray band corresponds to the cross-over between the synchrotron and inverse Compton peaks, i.e., the “valley” floor (see e.g., ON 231 and TXS 0506 + 056, Tagliaferri et al., 2000; IceCube Collaboration et al., 2018), and the hard-X-ray range provides constraints on the slope of the low-energy electrons. In both cases, given the observed photon spectrum, it is possible to infer the distribution of energies of radiating particles, and for relativistic jets, the low-energy radiating particles are the most numerous. This, in turn, provides good estimates of the total particle content of the jet, and strong constraints on the total kinetic power of the jet (Madejski et al., 2016), as well as its electron-positron pair content. Thus, for high-luminosity sources, hard X-ray observations by HEX-P will provide estimates of the jet power. We note that, in principle, one could study the low-energy end of particle distributions by examining the lowest energy part of the low-energy “hump”. However, the region producing the high-energy (X-to- $\gamma$ -rays) bump of the SED is assumed compact and close to the central black hole. The synchrotron emission from this region at the lowest radio frequencies is self-absorbed, thus the observed emission is optically thick and the SED at these frequencies is the sum of the compact core plus extended jet (e.g., Ghisellini and Tavecchio, 2009; Massaro et al., 2013; Giroletti et al., 2016; Fan and Wu, 2018; Boula and Mastichiadis, 2022). This hampers robust measurements of the total content of low-energy particles in the jet close to the base of the jet.

#### 4.1.2 Probing the relativistic hadronic population in blazar jets

The high-energy blazar spectral component may consist of one or multiple radiation processes. In the leptonic model, the high-energy component may have contributions from both SSC and EC, while in the hadronic model, the hadronic cascading component can be considerable in the X-ray to MeV  $\gamma$ -ray range. In either case, the hard X-rays often mark the transition from one radiation process to another, which usually manifests itself

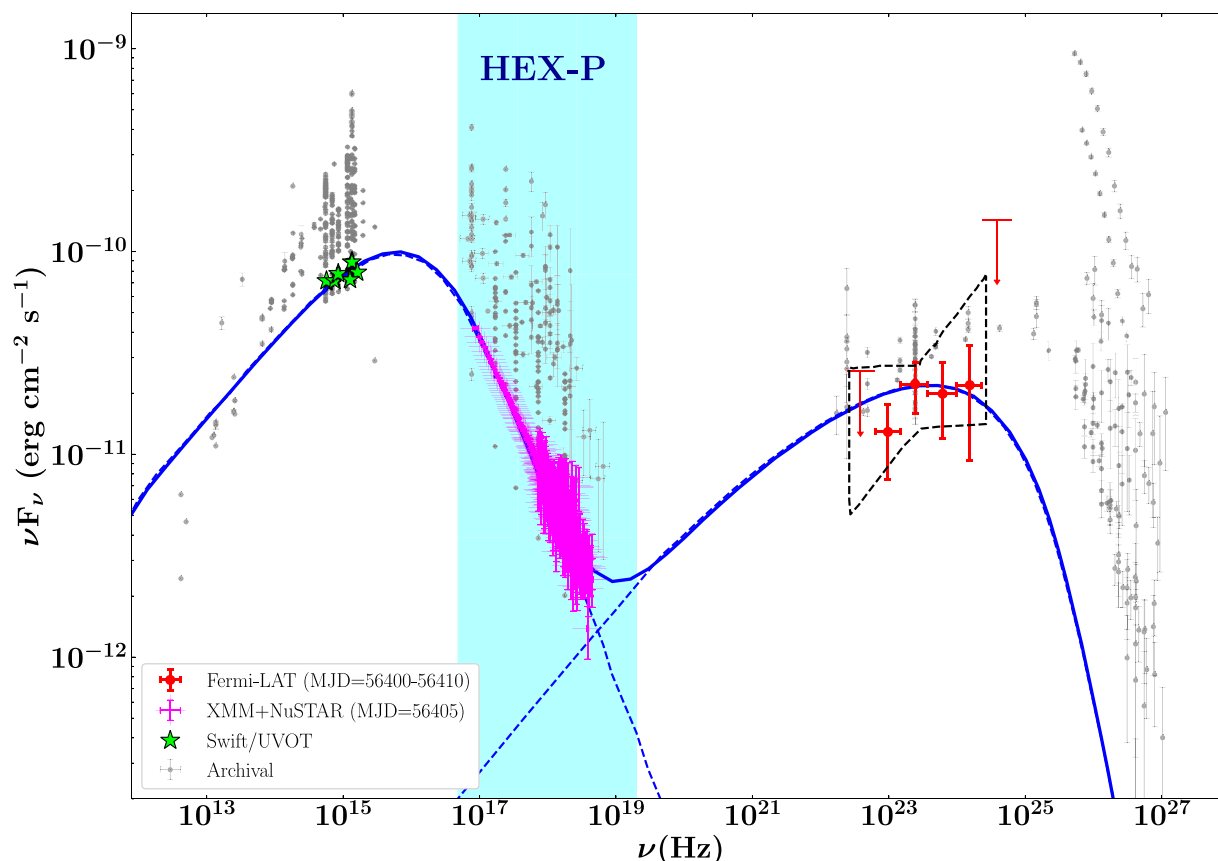


FIGURE 1

Broadband SED of the HSP blazar PKS 2155–304 adapted from [Madejski et al. \(2016\)](#). In gray, we show the multi-epoch multi-band archival data for the source, as an example of the extreme multi-band variability of blazar sources. The low-state was characterized by [Madejski et al. \(2016\)](#), collecting contemporaneous data from: NuSTAR and XMM-Newton-pn (MJD = 56405; magenta points), Swift-UVOT (fluxes corrected for Galactic reddening; green stars), and Fermi-LAT obtained for 10 days centered on the NuSTAR observation (data points in red; broad-band spectral fit “butterfly” in black). Also shown is a basic SED model including synchrotron and SSC components (blue lines). As mentioned in the text, the X-ray data occupy the “valley” in the SED. Highlighted in cyan is the energy band that will be simultaneously covered by HEX-P.

through a change in the spectral shape (e.g., [Böttcher et al., 2013](#); [Petropoulou et al., 2015](#); [Gao et al., 2017](#)). But if the high-energy component is dominated by only one radiation mechanism (such as SSC), we expect a smooth, power-law like, hard X-ray spectrum. With its unprecedented spectral resolution and broadband coverage, HEX-P will be able to identify the presence of more than one radiation mechanism at X-ray energies.

To illustrate the capabilities of HEX-P in identifying spectral breaks due to the transition from leptonic to hadronic emission components, we choose Ap Librae, an intriguing LSP at redshift  $z \approx 0.049$  ([Disney et al., 1974](#)) that has been detected at VHE ([H. E. S. S. Collaboration et al., 2015](#)). It is the only TeV BL Lac object with a detected kpc-scale X-ray jet ([Kaufmann et al., 2013](#)). VHE data in combination with the X-ray and high-energy (HE, 0.1–300 GeV)  $\gamma$ -ray observations with Fermi-LAT, revealed an unexpected broad high-energy hump, spanning more than nine orders of magnitude in energy (from X-rays to TeV  $\gamma$ -rays). [Petropoulou et al. \(2017\)](#) demonstrated that the superposition of different emission components related to photohadronic interactions in the blazar core can explain the broad high-energy component of Ap Librae without invoking external

radiation fields or particle acceleration at kpc scales (for large-scale jet scenarios, see [Section 6](#) and, e.g., [Zacharias and Wagner, 2016](#); [Roychowdhury et al., 2022](#)).

We adopt Model A from [Petropoulou et al. \(2017\)](#) (see in [Figure 3A](#)) to simulate spectra for HEX-P, assuming a 50 ks exposure. The model predicts a broadband X-ray spectrum with leptonic and hadronic contributions in the energy range probed by HEX-P. More specifically, inverse Compton scattered (ICS) emission by accelerated electrons in the jet contributes mostly at hard X-rays ( $> 10$  keV), while synchrotron emission from secondary leptons produced in a hadronic-initiated cascade is dominating in soft X-rays. The simulation results are presented in the [Figure 3B](#). The residuals of a power-law fit to the broadband spectrum have some structure which is improved when the spectrum is modeled with a tabulated leptohadronic model. In the latter case the statistical improvement is of the order of 100 using  $\chi^2$  or c-statistics for 5,800 degrees of freedom (unbinned data). Moreover, the residuals at 2 keV, where the spectral break is expected, are reduced in the leptohadronic model. We caution, however, that similar spectral breaks may also occur in a purely leptonic scenario, at the transition from SSC to EC emission, if a sufficiently dense external radiation field exists.

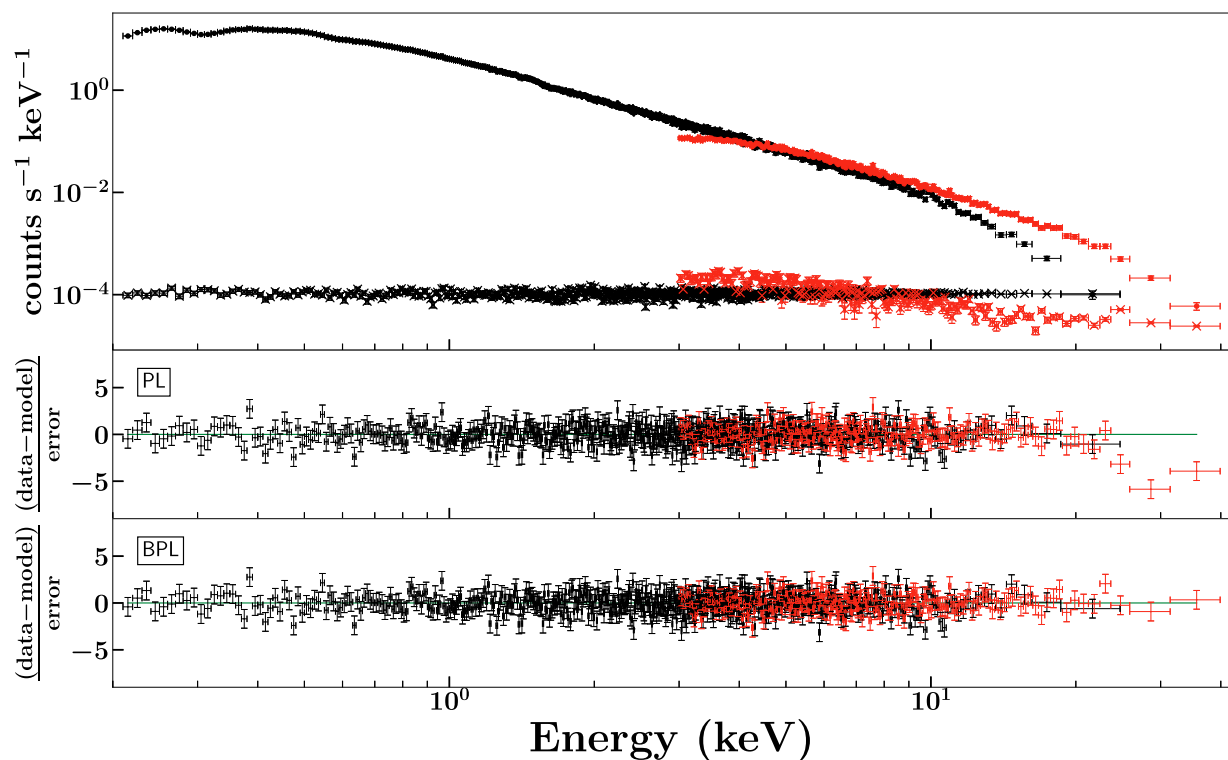


FIGURE 2

Simulated 100 ks observation of PKS 2155–304 with the LET (black data points) and the HET (red data points) assuming a broken power-law energy distribution of the underlying emission, with a break happening at  $\sim 20$  keV. The top panel also shows the background levels of the LET and the HET. The bottom panels show the residuals for a power-law and a broken power-law model. The low background of the HET would enable us to sample the distribution up to  $E \sim 40$  keV, hence disentangling the presence/absence of such a break at a  $\sim 5\sigma$  level. The higher NuSTAR background would not allow us to detect anything beyond  $E \sim 20$  keV for such a soft spectrum, even for a 1 Ms exposure.

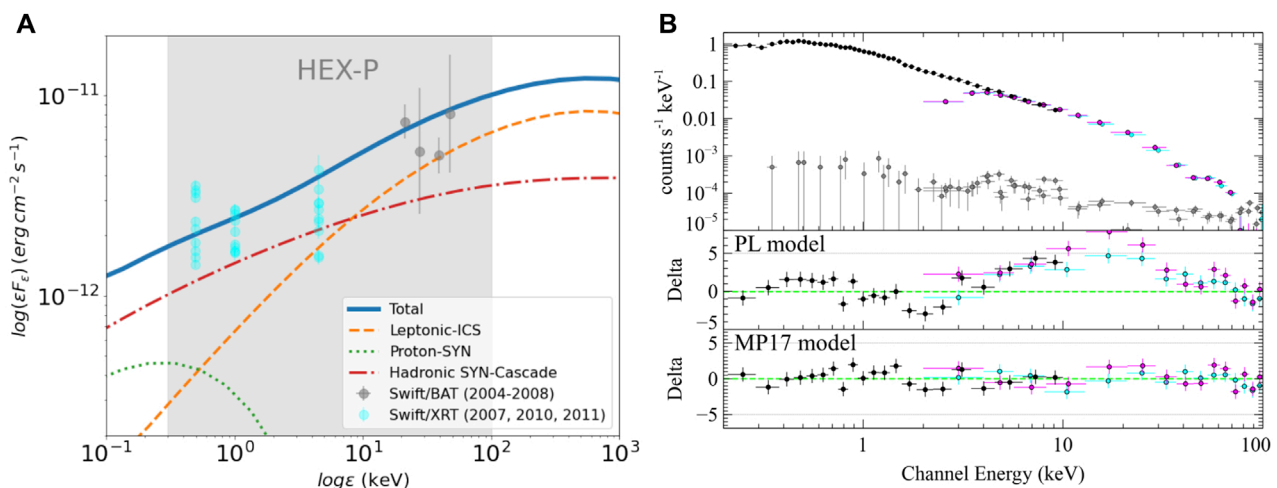
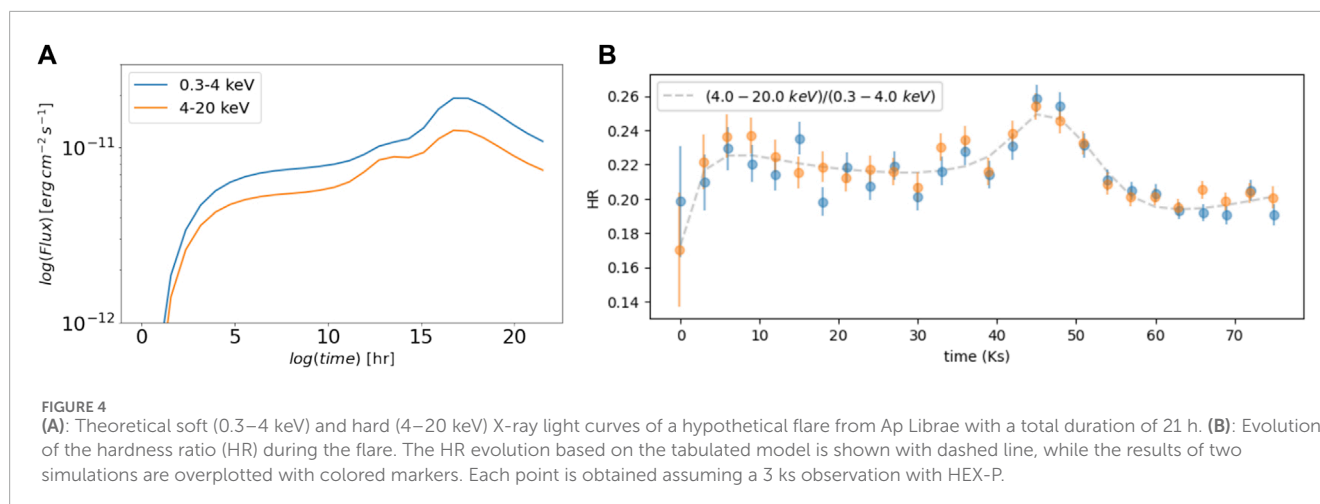


FIGURE 3

(A): Theoretical spectrum for the core emission of Ap Librae from [Petropoulou et al. \(2017\)](#) showing hadronic and leptonic contributions in the 0.2–100 keV range. Colored markers show non-simultaneous observations in soft and hard X-rays obtained with Swift/XRT and Swift/BAT, respectively. A smooth spectral break is evident at  $\sim 2$  keV due to the onset of leptonic inverse Compton scatter (ICS) over the hadronic synchrotron (SYN) cascade emission. (B): Simulated spectrum of Ap Librae for 50 ks exposure with HEX-P based on the theoretical spectrum shown on the left (top); background contribution is shown with gray shaded points. LET data are shown in black. Magenta and cyan colors correspond to the two high-energy detectors. Residual plots (i.e.,  $(\text{data}-\text{model})/\text{error}$ ) of a power-law fit (middle) and a fit with a tabulated leptohadronic model (bottom) to the simulated spectrum. The residuals indicate that the power law model is disfavored.





To demonstrate also the capabilities of HEX-P in detecting spectral changes during the course of an X-ray flare, we simulate a  $\sim 1$ -day long flare based on Model A of Petropoulou et al. (2017) by introducing variations in the injection rate of accelerated electrons and protons in the blazar core. The injection rate of both particle species is modeled with a Lorentzian as a function of time. The applied changes in the proton injection rate lag behind the variations in the electron injection rate by  $\sim 3$  hr. As a result, we expect to see first an increase in the hard X-ray flux, followed by an increase in softer X-rays where the hadronic cascade emission dominates (see left panel in Figure 3). The theoretical light curves of the flare in soft and hard X-rays are presented in the Figure 4A. We simulate 27 HEX-P spectra, 3 ks exposures each, and compute the hardness ratio (HR), defined as the ratio of counts in the 4–20 keV and 0.3–4 keV ranges. In Figure 4 we show the evolution of HR with time for two different simulations (blue and orange markers) that illustrate the statistical variance of the HR. Both simulations recover the expected changes in the HR shown by the dashed line.

## 4.2 Multimessenger follow-up: cosmic rays and neutrinos

Cosmic rays were first detected at the beginning of the 20th century (Wulf, 1909; Hess, 1912; Pacini, 1912), yet, more than 100 years later, the sources of the high and ultra-high energy cosmic rays remain unclear. Since protons are deflected in magnetic fields, they are not reliable tracers of their origin. Neutrinos that are produced in proton-proton and proton-photon interactions travel in a straight path, but are much harder to detect.

Blazars and AGN have long been suspected to be neutrino emitters. Material in the accretion disk and in the vicinity of the supermassive black hole includes partially or fully ionized gas, which is available for injection into the relativistic outflows. Protons in the jet can interact with photons from a stationary photon field (e.g., thermal UV emission from the disk), and produce neutral and charged pions. Charged pions decay into several products including several neutrinos, while neutral pions produce high-energy photons, detectable by HEX-P. These photo-pion processes are expected to happen in AGN jets (Biermann

and Strittmatter, 1987; Mannheim et al., 1992; Mannheim, 1993) and AGN cores (Eichler, 1979; Berezhinskii and Ginzburg, 1981; Protheroe and Kazanas, 1983; Begelman et al., 1990; Stecker et al., 1991). The photo-hadronic interactions likely responsible for the neutrino production in blazar jets inevitably lead to the development of electromagnetic cascades. These are initiated by synchrotron or inverse-Compton scattering by secondary electron-positron pairs produced in photo-pion and Bethe-Heitler pair production processes, and the subsequent decay of charged and neutral pions and muons. The radiative output of these cascades is expected to peak in the hard X-ray – soft  $\gamma$ -ray band (e.g., Petropoulou and Mastichiadis, 2015; Keivani et al., 2018; Murase et al., 2018; Gao et al., 2019; Zacharias, 2021).

Follow-up of IceCube neutrino alerts has been promising for identifying the source of astrophysical neutrinos and therefore cosmic rays. The associations of blazars with VHE neutrino events detected by IceCube have attracted much attention in the recent literature, with increasing evidence pointing towards blazars being responsible for at least part of the IceCube VHE neutrino flux (e.g., Krauß et al., 2014; Kadler et al., 2016; Garrappa et al., 2019; Plavin et al., 2020; Buson et al., 2022; Plavin et al., 2022). In 2017, a high-energy IceCube event (IceCube-170922A) was found in positional and temporal coincidence with a blazar flare from TXS 0506 + 056 (IceCube Collaboration et al., 2018), following the previous, lower-confidence coincidence of a PeV neutrino event and a blazar flare (Kadler et al., 2016). Surprisingly, an earlier neutrino flare consistent with the position of TXS 0506 + 056 between September 2014 and March 2015 is not accompanied by a  $\gamma$ -ray flare. In TXS 0506 + 056, X-ray observations in the soft and hard X-ray band were crucial in constraining the hadronic contribution in order to determine that the source is a likely neutrino counterpart. These cascading signatures were studied in detail for TXS 0506 + 056 by Reimer et al. (2019), who used the measured VHE neutrino flux from the neutrino flare coincident with the direction of TXS 0506 + 056 in 2014–2015 (IceCube Collaboration et al., 2018) to constrain the source parameters of the photo-hadronic interactions leading to the neutrino production. A crucial parameter in these calculations is the maximum optical depth to  $\gamma\gamma$  absorption within the source (typically peaking at TeV – PeV  $\gamma$ -ray energies), which determines the intensity of the cascades. In this case, the hadronic contributions

to the electromagnetic SED were found to be small, but consistent with the neutrino detection. Such small hadronic contributions are generally expected (Krauß et al., 2018), as electrons and positrons are much lighter, they are easier to accelerate and we therefore expect a high fraction of leptonic emission from AGN jets. The soft X-ray emission in the 2–10 keV band can be associated with either the synchrotron peak or the high-energy peak, depending on the peak frequencies. The purely leptonic synchrotron peak does not provide any information about hadronic fluxes. Only for the most extreme high-peaked blazars does the hard X-ray band not fall on the high-energy peak. In order to reliably constrain the contributions to the high-energy peak from pion cascades, it is necessary to have data in the hard X-ray band, which can only be provided by HEX-P with high angular resolution.

Previous detections of possible blazar counterparts were followed by the detection of possible Seyfert AGN counterparts, that can only be identified in the X-ray band: IceCube-190331A (Krauß et al., 2020) and the recently discovered neutrino source NGC 1068 (IceCube Collaboration et al., 2022). X-ray follow-up of likely-cosmic neutrino alerts has thus proven its utility for identifying possible neutrino counterparts, and most importantly for measuring hadronic contributions to the high-energy emission. HEX-P has a response rate of the order of 1 day (see Madsen et al., 2023). While we may detect neutrino from sources in quiescence, a short flare lasting less than a day or week does not increase the fluence of the source significantly to explain the neutrino event. As only the total flare fluence is relevant, sources in quiescence or long-lasting outbursts are the most promising source of neutrinos from blazars (Kreter et al., 2020b). The response time to triggers from HEX-P is therefore well suited to perform multimessenger follow-up observations.

Blazars have been ruled out as the dominant source of IceCube neutrino events in the TeV energy band through stacking analyses (e.g., Hooper et al., 2019). Blazars, especially BL Lacs are expected to produce neutrinos in the energy range of PeV to EeV (e.g., Murase et al., 2014; Padovani et al., 2015). These are energies that will be covered by the Askaryan Radio Array (ARA; Allison et al., 2016), currently under construction, as well as the IceCube-Gen 2 detector (Grant et al., 2019). IceCube Gen 2 with a 8 cubic kilometer volume is expected to be completed in 2032, which would coincide with the expected start of the HEX-P mission.

HEX-P will identify possible neutrino emitters and constrain the fluxes and emission mechanisms. One of the main challenges of neutrino follow-up is the large uncertainty regions of both cascade and track events. Track events are more promising for a direct identification of a counterpart source due to their smaller angular uncertainties,  $\geq 0.5^\circ$ . This currently requires Swift-XRT tiling observations with at least 7 pointings to fully cover the uncertainty area. This can often be reduced by using knowledge of previously catalogued X-ray sources. eROSITA will be crucial in providing a pre-identification of potential counterparts. HEX-P has a field of view of  $\approx 13' \times 13'$ . This will allow follow-up of known brighter X-ray sources in only a few pointings.

In order to test whether HEX-P observations are able to distinguish between a single power-law and the curved up-turns expected from proton-induced cascades, we performed detailed simulations. Three model cascade spectra (synchrotron-supported cascades) from Figure 12 of Reimer et al. (2019) (for  $\tau_{\gamma\gamma}^{\max} = 1, 10$ , and 100) were used

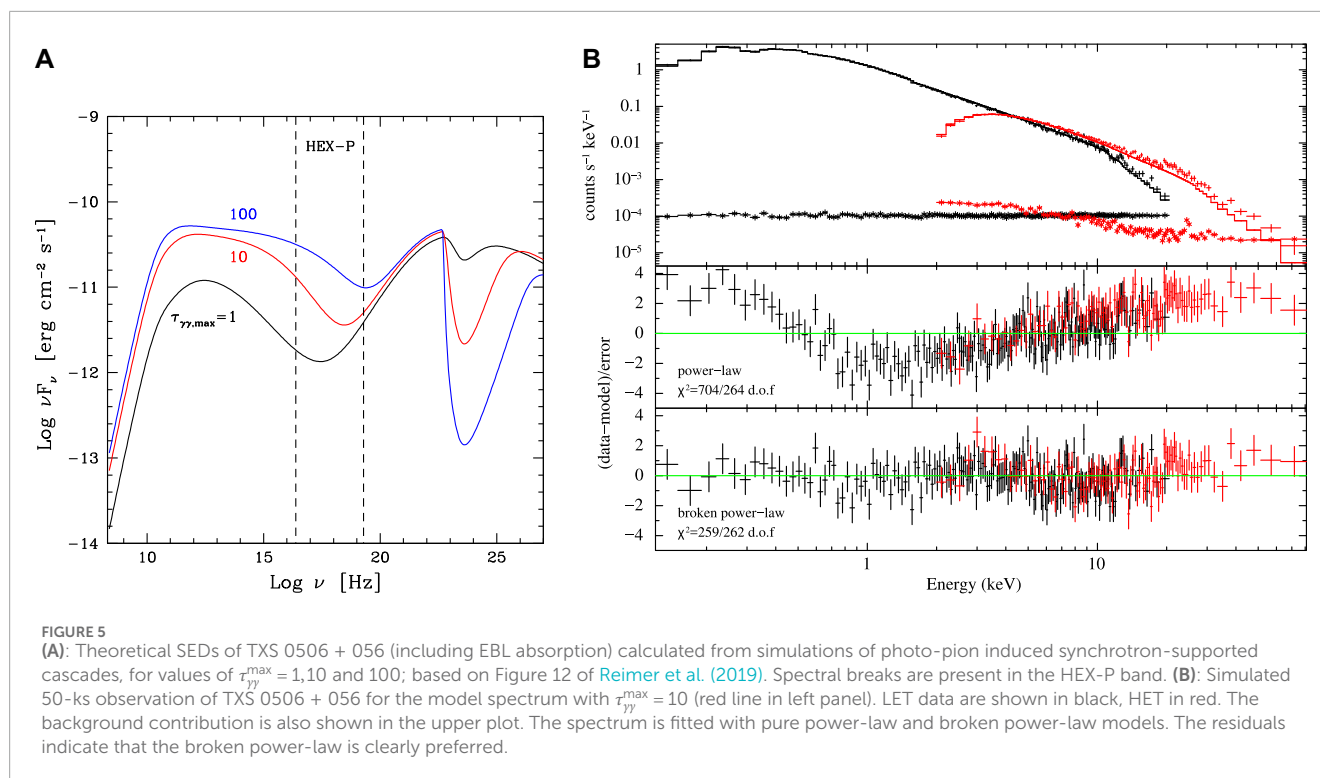
as input to simulate 50 ks observations with HEX-P, which were subsequently fitted with both a single power-law and a broken power-law model. An example for  $\tau_{\gamma\gamma}^{\max} = 10$  is shown in Figure 5. In all three cases, a single power-law could clearly be rejected, while a broken power-law was strongly preferred. Although this is not the only possible interpretation for such upturns, HEX-P can reveal and constrain this expected feature of the proton-induced cascades associated with photo-pion neutrino production.

### 4.3 Variability studies of blazars

Variability is a defining feature of all blazar subclasses, and studies have been conducted across the entire electromagnetic spectrum for many years, both with a focus on long-term (weeks to years) as well as short-term (days to hours) time scales. It is increasingly recognized that simultaneous broadband variability study is crucial to understand blazar physics. But the sample sizes of these studies depend largely on monitoring capabilities in the different energy bands. In the soft X-rays, monitoring campaigns can be performed for certain objects with Swift/XRT (Giommi et al., 2019; 2021), while observations with, e.g., XMM-Newton enable investigation of variability on time scales of hours (e.g., Zhang et al., 2006; Devanand et al., 2022; Noel et al., 2022). Long-term monitoring in the hard X-ray range is possible with Swift/BAT and MAXI (and in the past with RXTE), though, only for the brightest blazars. Variability studies above 10 keV on intraday time scales have been performed with RXTE (Wang et al., 2018) and later NuSTAR (Bhatta et al., 2018). Since no existing X-ray telescopes have sensitive broadband capabilities, simultaneous coverage of soft to hard X-ray variability has to be performed by carefully coordinating multiple X-ray telescopes pointing, which has only been arranged for a small number of observations.

The combination of high sensitivity and time resolution on the order of milliseconds and up to 80 keV will enable HEX-P to probe rapid variability, i.e., fluctuations within several minutes, at hard X-ray energies for a sample of bright blazars, which has not been feasible so far. In particular, we will be able to probe HSPs, for which minute timescale variability has been seen in gamma-rays and which could become evident only in hard-X-rays, since produced by higher-energy electrons, or a hardening at high energies as the result of the harder emission of freshly accelerated electrons on top of a more persistent, softer component.

In the case of FSRQs, HEX-P follow-up observations of flaring blazars will yield an unprecedented view up to 80 keV, i.e., full coverage of the rising half of the high-energy hump, and enable time-resolved spectral studies, which might be able to track changes in the high-energy emission related to, e.g., a change from SSC to EC emission over the course of a flaring event. This part of the spectrum can be a mixture of SSC and EC under the leptonic scenario, while under the hadronic scenario it may consist of SSC, hadronic cascades, EC, or proton synchrotron (Böttcher et al., 2013; Cerruti et al., 2015; Zhang et al., 2016). Theoretical studies have shown that these radiation processes can have different variability patterns, which can lead to distinguishable spectral and temporal signatures in the broadband X-rays. For instance, in the leptonic model EC often peaks at higher energy than SSC, thus the X-ray band can become dominated by a softer SSC spectrum from



higher-energy electrons (thus more variable) over a harder EC component by lower-energy electrons, leading to a “softer-when-brighter” behavior (see e.g., Ghisellini et al., 1999; Diltz et al., 2015; Zhang et al., 2015). Simultaneous broadband X-ray variability can therefore track the transition from SSC to EC, as well as unveil possible bulk-compton emission components, which are expected to be transient (Celotti et al., 2007). They can thus probe the underlying particle evolution, and put constraints on the origin of SSC and EC emissions.

Observations have shown that blazars of all types (low, intermediate or high-synchrotron-peaked objects) may experience epochs of extreme particle acceleration, resulting in its synchrotron peak shifting into or across the X-ray band. Most dramatic events have been shown for example, by Mkn 501 in 1997 (Pian et al., 1998; Tavecchio et al., 2001), 1ES 1959 + 650 in 2002 (Krawczynski et al., 2004) and more recently, again by Mkn 501 (Aliu et al., 2016; MAGIC Collaboration et al., 2020), Mkn 421 (Tramacere et al., 2009; MAGIC Collaboration et al., 2020), PKS 2005–489 (Tagliaferri et al., 2001; Chase et al., 2023) and 1ES 1215 + 303 (Valverde et al., 2020). Such phenomena are often referred to as the peak shift. The physical cause is generally understood as the emerging of a new hard emission component from freshly accelerated electrons, injected by a new dissipation event. The mechanism and location of this acceleration are instead still uncertain, being debated among shock acceleration (either in internal shocks or a standing shock), magnetic reconnection and/or hadronic origins (Böttcher et al., 2013; Zhang H. et al., 2022). Due to the significant spectral changes, observing these events needs broadband coverage. Temporal slicing of the SED evolution is therefore important to elucidate the formation of such extreme physical conditions. In Figure 6 we show as an

example of a simulated magnetic reconnection event for a high-peaked BL Lac object using the model of Zhang Z. et al. (2021). With very good broadband sensitivity and timing capability, HEX-P can track the peak shifts and spectral variations in greater detail than previously possible, revealing the physical driver of such extreme particle acceleration and probe potential hadronic origins. As such, HEX-P will serve as a space-based time domain astronomy instrument; a field that has been identified as one of the key research topics in the Astro2020 decadal survey<sup>2</sup> (National Academies of Sciences, Engineering, and Medicine, 2021).

For high-peaked BL Lac objects, spectral changes have been observed during flares, during which the spectrum often becomes harder when the source gets brighter. The study of light curves covering different energy ranges can reveal the existence of lags, which are indicators of the efficiency of the underlying acceleration and cooling processes. Interestingly, these lags are not consistent: for example, the well-studied source Mrk 421 has shown both soft and hard lags at different times, without, so far, a consistent pattern (Zhang et al., 2010; Kapanadze et al., 2016; Noel et al., 2022). HEX-P will be an essential instrument for these studies, in particular for flare follow-up observations, as it covers both the soft and hard X-ray range with high sensitivity.

One key discovery in the last decades has been flux variations close to or below the light-crossing time scale of the super-massive black hole (SMBH, Aharonian et al., 2007b; Aleksić et al., 2014b; Ackermann et al., 2016b), which revealed extreme particle acceleration in very localized regions. These rapid flux variations

<sup>2</sup> <https://cor.gsfc.nasa.gov/Astro2020/26141.pdf>

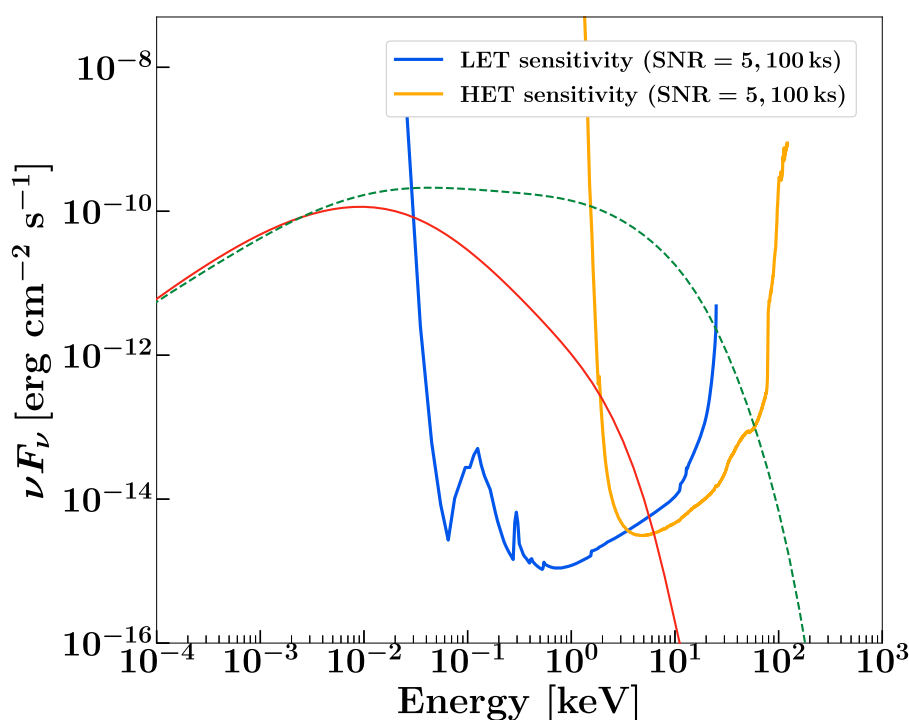


FIGURE 6

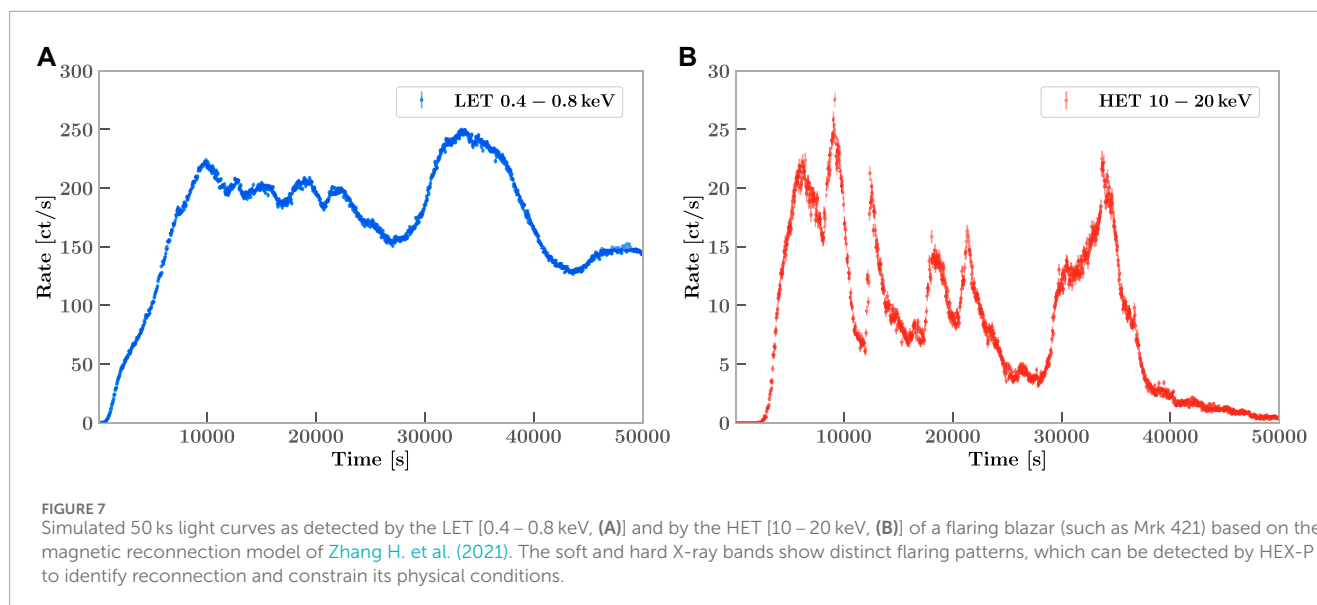
Simulated synchrotron peak of a high-peaked BL Lac object in quiet state (red line) versus the shifted synchrotron peak from the same source during flaring activity (dashed green line). Both the quiescent and the flaring state are produced by the magnetic reconnection in the same current sheet using the numerical simulation from Zhang H. et al. (2021). Both the quiescent and flaring state spectra result from plasmoids in the magnetic reconnection region, which has a fixed amount of initial magnetic energy. However, the flare is due to a major plasmoid merger event, which leads to significant dissipation of magnetic energy and efficient acceleration of nonthermal particles. The newly accelerated particles reach a much higher energy and harder spectrum than particles in other plasmoids, producing a peak shift. The sensitivity curves for both the LET and HET of HEX-P illustrate the excellent capability of the X-ray telescope to precisely measure changes in the position and shape of the synchrotron peak of blazars.

provide invaluable information about the particle acceleration mechanisms in jets and are broadly connected to the capability of the jet of accelerating protons and producing neutrinos. Observed, so far, in a few exceptional flares, such rapid gamma-ray variability has shown the limitations of standard emission models (Begelman et al., 2008; Ghisellini and Tavecchio, 2008; H. E. S. S. Collaboration et al., 2019) and prompted more sophisticated scenarios (see, e.g., Aharonian et al., 2017), such as turbulent, multi-zone emission (Marscher, 2014), magnetic reconnection inside the jet or in the jet launching region (see, e.g., Cerutti et al., 2012; Giannios, 2013; Kadowaki et al., 2015; Christie et al., 2019; Zhang H. et al., 2021), jet-star interactions (Barkov et al., 2012), or acceleration in the black-hole magnetosphere (Levinson and Rieger, 2011). These scenarios predict different signatures in both spectral and temporal behaviors, but require broadband time-domain observations to further understand the jet physics. HEX-P would be able to detect a break in the broad band X-ray spectra that is predicted by some models (jet-star interaction) due to the emergence of an additional synchrotron component, which would rise in the hard X-ray band. Alternatively, simulations of processes dominated by magnetic reconnection and turbulence predict distinct temporal patterns in the X-ray band (e.g., Zhang H. et al., 2020; Zhang H. et al., 2023). Those require highly resolved timing data across the full X-ray band in order to be detected, which HEX-P would be able to provide. At VHE,

CTAO is expected to probe broadband (sub-TeV to  $>100$  TeV) minute-scale variability in several other sources with less exceptional flares, as well as on sub-minute timescales for events like the one detected from the blazar PKS 2155–304 in 2006 (Aharonian et al., 2007b; Sol et al., 2013).

Current theories suggest that such short timescale variability should also be present in X-ray synchrotron emission, but so far the few simultaneous data show that, even when tightly correlated, the (soft) X-ray variations have typically been of much lower amplitude than at VHE (Aharonian et al., 2009a) and in general on longer timescales (tens of minutes). While X-ray synchrotron emission may suffer from a stronger dilution of variable flux by brighter persistent emission, a key issue in observation is that current X-ray instruments lack the same broadband sensitivity as the VHE telescopes. The synchrotron emission of the rapidly-variable TeV electrons could peak at higher X-ray energies, in a band not yet investigated because of the lack of large collection areas above 5–8 keV. Additionally, several models, such as magnetic reconnection, predict distinct variability patterns in soft and hard X-ray bands (Figure 7), which require broadband coverage with good temporal sensitivity (Zhang H. et al., 2021). HEX-P will be able to probe, for the first time, such short timescales simultaneously at both soft and hard X-ray energies.

Another aspect is the different inter-band variability among X-ray bands. The interplay between acceleration, cooling, and



escape timescales of electrons at a shock front or dissipation region can produce different and characteristic patterns in the variation of spectral index vs. flux. For example, such variations can have different “hysteresis loops” (e.g., Kirk et al., 1998; Kirk and Mastichiadis, 1999; Böttcher and Chiang, 2002; Abeysekara et al., 2017; Wang et al., 2018; Böttcher, 2019; Chandra et al., 2021). The timescales of the flux rise and fall provide information on the acceleration mechanisms, if nearly instantaneous or progressive, and these could be different between hard and soft X-rays. So far, such studies have been limited to a few objects observed during large flares over 1 decade in energy. HEX-P will be able to investigate these relations over more than 2 decades in energy, and at lower fluxes and/or shorter timescales (see Madsen et al., 2023, for HEX-P repointing capabilities).

## 4.4 Synergies with other telescopes

### 4.4.1 CTAO

In HSP blazars, electrons accelerated to TeV energies emit by synchrotron typically in the X-ray band and by SSC in the VHE band. The possibility to perform extended simultaneous X-ray–VHE observations therefore provides a powerful diagnostic tool to understand particle acceleration and emission, allowing the evolution of a single freshly-accelerated electron population to be sampled by two different emission mechanisms. This is especially important during flares with hour and sub-hour variability. To perform such studies, an X-ray satellite must have a large visibility window at any given epoch for most of the sky and must be able to point in an antisolar direction. A large visibility window increases the chance that a source can be observed when it is found flaring, while the antisolar direction enables simultaneous observations at night with ground-based telescopes when the target culminates at the highest altitudes above the horizon, and thus can be followed by Cherenkov telescopes continuously under the best conditions for at least several hours. Such studies are not possible with XMM-Newton due to its strong solar panel constraints (Costamante, 2004).

Chandra showed the potential for such studies with the detection of large flares by PKS 2155–304 in 2006, albeit with a much lower collecting area than XMM-Newton (Aharonian et al., 2009a). At an L1 orbit, HEX-P will fulfill all these operational requirements, allowing us to perform such studies for the first time with a sensitive, flagship-class X-ray telescope.

### 4.4.2 IXPE & COSI

High-energy polarization, a recently opened new window to the Universe, will play a very important role in future efforts to understand the matter content and high-energy emission processes in jets. In a leptohadronic scenario, emission at the trough can be dominated either by proton synchrotron or synchrotron from secondaries (Zhang H. et al., 2019), and hence it can be highly polarized in the X-ray and MeV bands. Recent IXPE observations have demonstrated that determining the shape of the X-ray spectrum is crucial to correctly recover the polarization parameters, regardless if X-rays are part of the synchrotron or high-energy spectrum (Ehlert et al., 2022; Di Gesu et al., 2022; Liodakis et al., 2022; Middei et al., 2023a; b). This is particularly important for sources like TXS 0506 + 056, i.e., ISP blazars, where the X-rays lie in the trough between the low- and high-energy humps of the SED (Peirson et al., 2022; 2023). Additionally, IceCube observations of TXS 0506 + 056 suggest that the hadronic spectral component, potentially peaking in the MeV band, can be variable (IceCube Collaboration et al., 2018). Therefore, both the spectral shape and variability of the hadronic contribution can have significant effects on the X-ray to MeV polarization degree and angle (Zhang et al., 2016), therefore requiring broadband coverage in X-ray and MeV bands. However, IXPE and future X-ray polarimeters, like the enhanced X-ray Timing and Polarimetry mission (eXTP), have limited energy ranges. For example, IXPE is sensitive in the 2–8 keV band. On the other hand, future MeV missions, such as the Compton Spectrometer and Imager (COSI, Tomsick et al., 2019), observe in the MeV band (0.2–5 MeV), which cannot fully capture the rising part of the hadronic component (in the X-ray bands) and its variability. In the X-ray band, the



spectrum of the HSP blazars often shows significant curvature up to tens of keV, often modeled with a log-parabola model (see, Donato et al., 2005; Bhatta et al., 2018; Middei et al., 2022). For some ISP blazars the spectrum is often characterized best with the sum of a soft and a hard powerlaw or a power law with an exponential cut-off and a hard powerlaw. HEX-P's broad energy range can uniquely allow for the accurate modeling of the broadband X-ray spectral shape and transition region in a large number of blazars, paving the way for detecting polarization from the jet's high-energy component. Coordinated observations with missions current or future X-ray polarimeters (e.g., IXPE, eXTP etc.), and future MeV missions (COSI, ASTROGAM, AMEGO-X; De Angelis et al., 2017; Tomsick et al., 2019; Caputo et al., 2022) will enable measurements that can pinpoint or rule out hadronic processes by comparing the predicted and observed polarization degrees.

#### 4.4.3 SPT (CMB)

As discussed in Section 4.3, truly simultaneous, broad-band monitoring is key in unraveling blazar jets structure as well as the emission mechanisms and thus particle acceleration processes operating in blazars. One important spectral band is the millimeter regime, where, for the high-luminosity (FSRQ-type) blazars, the low-energy part of the SED peaks: blazars are generally strong mm sources and their mm emission can even be used to find counterparts of un-identified Fermi-LAT gamma-ray sources (see e.g., Zhang L. et al., 2022). This is the spectral regime where the blazar synchrotron emission is believed to become optically thin, and thus reveals the entire volume of the jet rather than its surface.

Time-resolved blazar studies in the mm band were usually conducted by individual radio/mm telescopes, generally dedicated to a single object at a time, and clearly revealed the large amplitude variability on time scales of days. Such studies also revealed a strong degree of polarization (in some cases, more than 20%), which is key in determining the orientation and uniformity of the magnetic field. Current and planned Cosmic Microwave Background (CMB) facilities will provide such data as a part of the CMB measurements: while the astrophysical sources are “noise” for CMB maps, the data – obtained “for free” – are very valuable for studies of such sources. This holds particularly true for blazars, since the current and planned CMB facilities collect data by continually scanning relatively large regions of the sky, and subsequently combining the data to form sky maps. Such time-resolved data result, generally, in a relatively well-sampled time series for many sources. Importantly, these CMB facilities, in addition to flux, also collect polarization information, thus providing time series in both intensity and polarization. One recent “proof of concept” example is observations of the blazar PKS 2326–502 with the South Pole Telescope (SPT; Hood et al., 2023) which revealed correlated flaring in the mm (at 150 GHz) and GeV-range gamma-ray bands, measured with the Fermi-LAT. For bright blazars in the millimeter band, SPT will provide a time-resolved flux and polarization time series in the SPT bands. Correlating such time series with hard X-ray flux will provide additional constraints on the origin of hard X-ray flux in blazars. Moreover, for HSP objects, the combination of SPT, HEX-P, and other observatories at IR/optical wavelengths would allow us to sample the energy distribution of the accelerated particles from the lowest to the highest energies, simultaneously sampling the

SED over 5 orders of magnitude and informing us about the full particle populations in the jet.

The next-generation CMB facility is known as “CMB Stage 4” and is planned for deployment in the early 2030s. This matches well with the proposed timeline for HEX-P. CMB Stage 4, described in Carlstrom et al. (2019) and Abazajian et al. (2022)<sup>3</sup> is significantly more sensitive than the current CMB facilities and, in its full form, will cover well over 50% of the sky. The facility will have detectors sensitive over the range of 30 GHz–280 GHz. While the details of the instrumental sensitivity are still under development, it is expected that in 24 h worth of sky scanning, CMB Stage 4 should be able to measure a source flux of 30 millijansky ( $5\sigma$ ). This complements the blazar studies undertaken with HEX-P *without altering the strategy of the anticipated CMB studies*, and such synergistic studies (with various classes of astrophysical sources, not just blazars) are already part of the instrument design and the data management plan.

## 5 Source classes (MeV blazars, extreme BL lacs, and masquerading BL lacs)

### 5.1 The accretion flow to jet power connection

Flat Spectrum Radio Quasars are defined as blazars with strong broad emission lines. They are generally associated with strong jet emission, peaking at low frequencies both in synchrotron and inverse Compton components. Similar to low-peaked BL Lacs, their high-energy component dominates the X-ray emission, particularly the rising section of the hump. Depending on redshift, different parts of this emission are mapped with specific telescopes: among the most widely used X-ray instruments, Swift/XRT, XMM-Newton and the Chandra X-ray Observatory map the softer X-ray emission ( $< 10$  keV), while Swift/BAT reaches 500 keV, and INTEGRAL reaches 10 MeV (see e.g., Giommi et al., 2019; Marcotulli et al., 2022). NuSTAR provides a sort of “link”, by covering the energy range between 3 and 80 keV (Ghisellini et al., 2019). The two instruments onboard HEX-P guarantee an energy range that can overlap with all X-ray telescopes currently available to the astronomical community. With a single observatory, the inverse Compton component of low-frequency peaked blazars, including FSRQs, can be mapped throughout cosmic time, by focusing on their evolution and possible variability, spanning across the whole inverse Compton emission of this AGN class. On the soft X-ray side, LET will be able to shed light on the amount of possible absorption or intrinsic breaks at soft energies. The nature of such a soft break is still under scrutiny, and the recent analysis performed on different samples of FSRQs do not allow us to constrain whether they are due to excess absorption or intrinsic jet features, such as peculiar energy distributions of the electrons involved in jets' high-energy emission (see e.g., Gaur et al., 2021, for a thorough discussion). The significant increase in sensitivity given by the LET and HET combination will likely make a difference in discriminating between the two possibilities. The HET instead will extend spectroscopic studies

<sup>3</sup> <https://cmb-s4.org/>

currently done by Swift/XRT, Chandra and XMM-Newton up to significantly higher energies, getting simultaneously close to the inverse Compton peak. The wide spectral coverage will allow us for the first time to uniformly characterize this section of FSRQ SEDs, crucial for the study and interpretation of jet physics and emission processes.

Even though FSRQs have been known and studied for more than 60 years, many open questions remain about their relativistic jets: how do they relate to the accretion flow, whether their launching system can be associated with it, and how the accretion-jet system evolves across cosmic time. The better insight on the electron energy distribution, and a better mapping of the spectral evolution of the inverse Compton emission in time will already be a significant advancement in relativistic jet physics. Nonetheless, HEX-P's wide energy coverage and deep sensitivity has the potential to make even bigger improvements on our knowledge.

The connection between accretion flow and the relativistic jet in blazars is in fact a key ingredient to understand the process behind jet launch and acceleration. Relativistic jets are generally associated with geometrically thick accretion flows when their launching mechanism is modelled: an accretion flow funnel close to the central black hole, along with a strong magnetic field, are effective ingredients to accelerate plasma along the polar direction (Tchekhovskoy et al., 2011). In contrast, a geometrically thin accretion disk has not historically been considered a solid basis for an effective jet collimation. FSRQs show that a relativistic jet can coexist with an optically thick, highly UV-luminous accretion flow, since they show a distinct and luminous optical-UV thermal continuum that cannot be justified with jet emission, with strong and broad emission lines, along with powerful relativistic jets. Such optical-UV features are most likely interpreted as coming from a geometrically thin, optically-thick accretion disk ionizing gas clouds in its vicinity (e.g., Padovani et al., 2017). Moreover, the direct correlation between jet and accretion power in blazars suggests a significant role in the jet launching process by the accretion disk, or at least in providing power for the overall process (Jolley and Kuncic, 2008; Ghisellini et al., 2014). The coexistence of these two structures is thus under deep scrutiny: coordinated multiwavelength observations of high-energy jet emission and optical-UV radiation from the disk are crucial to test the actual relation between disk and outflows. In particular, analogous to the cyclic X-ray binaries evolution, FSRQs might show variations in geometry and emission physics of the central part of their accretion flows, corresponding to jet launching or rebrightening. In case of an accretion rate fluctuation, a thin accretion disk might thicken towards its inner radii, providing a temporary funnel, allowing for significant magnetic field advection and ultimately accelerating a new emitting region in the relativistic jet. Coordinated fluctuations would also be expected during flux variations, and they would be easily identified during coordinated optical and X-ray monitoring. On the other hand, no thick component might be necessary: in this case, the launching mechanism might be different from what we currently expect, and this specific multiwavelength simultaneous or correlated variability would not be observed. Coordinated optical-UV and X-ray observations can thus provide unique insights into the most powerful engines in the Universe. HEX-P is crucial in this sense, being able to thoroughly map the jet emission profile and its evolution over a broad energy range. Changes in jet power

or acceleration would be easily detectable thanks to changes in the inverse Compton profile, well explored by the combined effort of LET and HET.

A good knowledge of the emission profile of high-energy blazar SEDs is crucial to nail the connection between jet and accretion flow in these sources. A reliable representation of the median blazar SED is necessary to identify variations in the jet emission, possibly linked to accretion fluctuations. The current state-of-the-art is mainly focused at low redshift. Blazars with different jet powers are known to differ in their pre-peak inverse Compton slope and inverse Compton peak frequency. Even though the overall SED blazar shape is always characterized by the two prominent synchrotron and inverse Compton humps, their relative dominance and peak positions vary significantly depending on the jet power. Whether physically motivated or just an observing bias, blazars appear to peak at lower frequencies and have a harder X-ray spectrum when they are intrinsically more powerful. This has been tested among the known  $\gamma$ -ray detected blazars, that are mainly concentrated below  $z \sim 3$ , but a detailed study of the inverse Compton emission evolution across cosmic time is still incomplete. A significant limit is posed by the narrow X-ray coverage of the currently available sensitive X-ray observing facilities: hard X-rays are needed to explore the inverse Compton peak position, and no current telescope reaches deep enough to detect high redshift sources ( $z > 3$ ). In this sense, HEX-P will be transformative. For the very first time, the astronomical community will have the possibility to explore the evolution of AGN high-energy jet emission. At  $z > 3$ , blazars appear to host among the most powerful jets known, with very high inverse Compton luminosities ( $L_\gamma > 10^{47} \text{ erg s}^{-1}$ ) and MeV-peaking inverse Compton humps. This is likely a selection bias: we are currently observing and classifying only the most powerful sources (Marcotulli et al., 2020; Sbarrato, 2021). The possibility to consistently study most of the pre-peak inverse Compton SED over a wide redshift range ( $1 < z < 4$ ), offered by HEX-P, will probe inverse Compton evolution across cosmic time. Limited (in numbers) but self-consistent (in method) results could already be drawn from a 1 Ms wide-field survey (see Section 5.5). Figure 8 shows how the expected HEX-P sensitivity limits and energy ranges will allow a comprehensive study of the most powerful FSRQs over a large redshift range.

## 5.2 Supermassive black holes at the dawn of time

Due to their high luminosities, we are able to trace blazars, particularly FSRQs, back to the early Universe. The entire broadband spectrum of these objects is shifted towards lower energies, and because of their high-energy hump peaking in the MeV band, they have been coined "MeV blazars" (Bloemen et al., 1995). These sources are more luminous in the hard X-ray band (Sbarrato et al., 2012), and studies made use of the NuSTAR satellite to analyse and characterise the emission processes responsible for the high-energy hump, in combination with Fermi-LAT at gamma-ray energies (Sbarrato et al., 2013; Tagliaferri et al., 2015; Marcotulli et al., 2017; 2020). The most puzzling findings about SMBHs hosted by quasars in the early Universe is that some of them exceed one billion solar masses, detected out to redshift of  $z > 7$  (Fan et al., 2001; Mortlock et al., 2011; Wu et al., 2015). Obviously, there is a bias

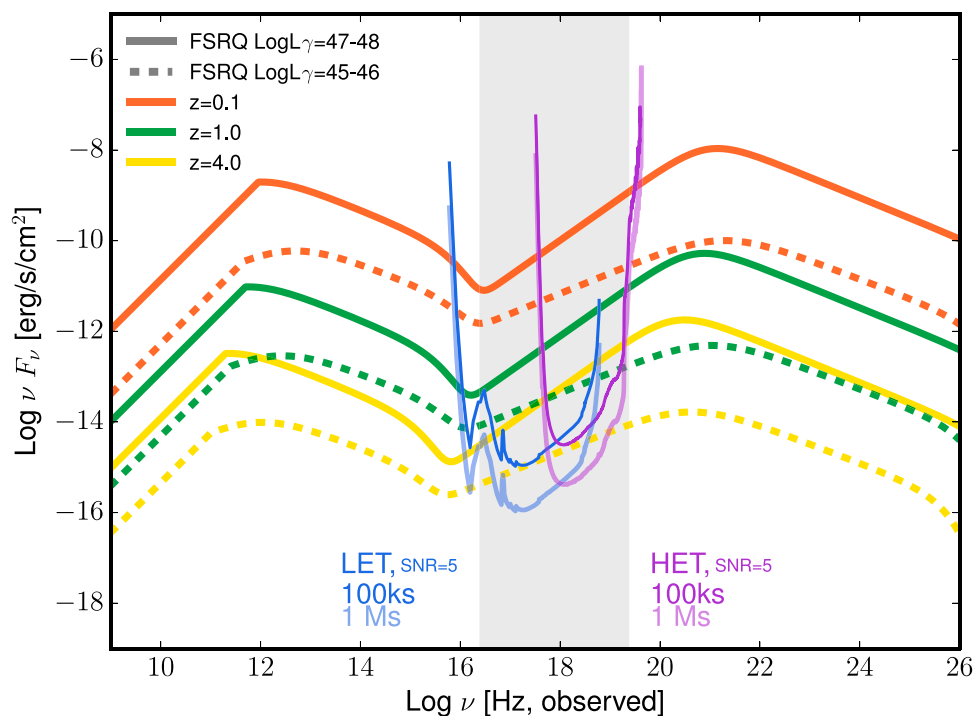


FIGURE 8

FSRQ SEDs with different powers and redshifts, compared with LET (blue) and HET (purple) differential sensitivities, calculated with 100 ks (thin, bright) and 1 Ms (thick, faint) exposures. The SEDs assume the Ghisellini et al. (2017) phenomenological models, corresponding to  $L_\gamma = 10^{47} - 10^{48} \text{ erg s}^{-1}$  (solid lines) and  $L_\gamma = 10^{45} - 10^{46} \text{ erg s}^{-1}$  (dashed lines). Different colours correspond to different redshifts. Note that the most powerful blazars will be easily detectable up to 80 keV and up to  $z \sim 4$ . HET detections will be more challenging for less powerful FSRQs or sources in a low jet state.

towards the most luminous objects being detected, but the discovery of such massive black holes questions our understanding of how SMBHs form and how fast they can grow. Current theories involve either the creation of  $\sim 100 M_\odot$  BH seeds as remnants of Population III stars (Madau and Rees, 2001; Schneider et al., 2002), which then require sustained super-Eddington accretion growth periods (Begelman and Volonteri, 2017; Yang et al., 2020; Johnson and Upton Sanderbeck, 2022), or the direct collapse of gas into  $10^4 - 10^6 M_\odot$  SMBHs in pre-galactic discs or halos (e.g., Bromm and Loeb, 2003; Begelman et al., 2006; Lodato and Natarajan, 2006). Another hypothesis is the formation of massive BH seeds within star clusters with masses of  $\sim 1000 M_\odot$ , which then grow due to the copious supply of matter in the star cluster (Devecchi and Volonteri, 2009; Alexander and Natarajan, 2014).

A fraction of these SMBHs will power jets that, if pointing towards us, should appear as MeV blazars. In particular, the detection of these sources in the gamma-ray band is limited by the sensitivity of current instruments, paired with a very low flux due to the immense distances. Less than 30 high-redshift ( $z \geq 2.5$ ) blazars have been detected at gamma-ray energies with *Fermi*-LAT (Acero et al., 2015; Ackermann et al., 2017; Liao et al., 2018; Abdollahi et al., 2020; Kreter et al., 2020a), with the most distant source having a redshift of  $z = 4.72$ . In addition, EBL absorption attenuates high-energy gamma-ray emission above  $\sim 10 \text{ GeV}$  (H. E. S. S. Collaboration et al., 2013; Ackermann et al., 2017; Domínguez et al., 2023). The large majority of high- $z$  blazar detections at gamma-ray energies were only possible due to flaring

episodes (Paliya et al., 2019a; Kreter et al., 2020a), which biases our studies towards more variable sources. At X-ray energies, instruments like *ROSAT* and *Swift*/BAT have been able to perform unbiased all-sky surveys, but have a limited sensitivity. Pointing observatories like *Chandra*, *XMM-Newton*, and *NuSTAR* have been used to study a selected sample of high-redshift blazars (e.g., Bassett et al., 2004; Young et al., 2009; Ighina et al., 2019). In particular, *NuSTAR*, with its sensitivity reaching up to 80 keV, has significantly contributed to the study of high- $z$  blazars by providing hard X-ray data that is essential to constrain the shape and peak of the high-energy hump of the blazars' SEDs (e.g., Marcotulli et al., 2017; Paliya et al., 2020; Middei et al., 2022). Similarly, HEX-P will enable us to continue studying blazars in the early Universe.

AGN create feedback mechanisms in their host galaxies and can be linked to their evolution (Silk and Rees, 1998). On large scales, jets influence their surroundings by creating shocks in the surrounding gas (Schoenmakers et al., 2000; Rawlings and Jarvis, 2004; Shabala et al., 2011; Raouf et al., 2017), leading to truncation of star formation. Studying AGN-galaxy co-evolution in the early Universe has been difficult and so far has only been done up to redshift  $z = 4$  for jetted AGN (Poitevineau et al., 2023). Due to its superior sensitivity compared to current X-ray telescopes, HEX-P will detect sources with fluxes as low as  $10^{-16} \text{ erg cm}^{-2} \text{ s}^{-1}$  (Civano et al., 2023). With this sensitivity, and the anticipated survey area, we used the prediction on the source count distribution of blazars from Marcotulli et al. (2022) employing a density and a luminosity evolution model to predict how many MeV blazars

could be detected. This number ranges from 50 for the luminosity evolution and 1,000 for the density evolution (see Figure 11). This will include a large number of newly detected extragalactic objects, including high-redshift blazars. The discovery of new sources will be less biased towards the most luminous sources, and enable us to get a better understanding of the cosmic evolution of blazars (see Section 5.5), while hard X-ray data above 20 keV will help distinguish the different high-energy processes in blazar jets, and allow comparisons of local and early Universe objects.

### 5.3 Extreme BL lacs

Observations in the last 2 decades have shown that BL Lacs can reach much higher SED peak energies than previously thought, surpassing by 1–3 orders of magnitude or more the values of the bulk of the HSP population. They have been thus called *extreme BL Lacs* (Costamante et al., 2001). They come in two flavours: extreme synchrotron and extreme-TeV objects, depending on whether they show extreme peak energies in the synchrotron or gamma-ray emission. By convention, an HSP blazar starts to be called extreme if the  $\nu_{\text{peak}}$  is above 1 keV or 1 TeV, respectively, but these sources have displayed peak energies in excess of 100 keV (both in flaring and quiescent states; Pian et al., 1998; Costamante et al., 2001) and 10 TeV (Aharonian et al., 2007a), respectively.

Conventional one-zone leptonic models (see also Section 4.1.1) struggle to model SEDs peaking at such high-energies, especially in gamma-rays. They generally require extreme conditions with very low magnetic fields, several orders of magnitude below equipartition, unusually hard electron distributions and very low radiative efficiency (Costamante et al., 2018; Biteau et al., 2020; Zech and Lemoine, 2021). The relation between extreme synchrotron and extreme TeV emission is still not well established: they should be correlated in a standard SSC scenario, but there are objects which are extreme in synchrotron but not in gamma-rays, and *vice versa*, although extreme-TeV sources do tend to display also extreme synchrotron peaks more often than the other way round (Biteau et al., 2020). Some objects show large peak variability, both in amplitude as well as in energy, with shifts by more than 3 orders of magnitude during flares (see, e.g., Mkn 501, Ahnen et al., 2018; MAGIC Collaboration et al., 2020). Others instead show rather steady extreme features over several years, a fact that could hint at the contribution of different mechanisms to the gamma-ray emission, in addition to inverse Compton emission by ultrarelativistic electrons (Biteau et al., 2020).

The location of the synchrotron peak in the SED and the slope of the spectrum around it are essential parameters to constrain the emission models and to study the origin of the peak itself (if due to cooling, escape, or the end of the accelerated particles' distribution). So far, this has been possible only for peaks up to 2–5 keV with great sensitivity, or up to 30–50 keV for only the brightest objects.

Thanks to the unique combination of broad-band coverage, low background, and large collecting area from 0.2 to 80 keV with a single observation, HEX-P is ideally suited to identify and study these objects: it can directly measure or constrain the peak emission at flux levels or timescales one to two orders of magnitude lower than any current instrument (see Figure 9).

In this respect, the synergy with the upcoming CTAO—which will bring a similar improvement of one to two orders of magnitude in sensitivity at multi-TeV energies—cannot be understated. Besides characterizing VHE spectra over a wide energy band, simultaneous observations can test if and when these are indeed the same electrons emitting by two different mechanisms (synchrotron and SSC), as indicated by the tight correlations observed in HSPs down to minute timescales in Mkn 421 (Maraschi et al., 1999; Fossati et al., 2008; Baloković et al., 2016; MAGIC Collaboration et al., 2021) and PKS 2155–304 (Aharonian et al., 2009a). CTAO and HEX-P coordinated observations will also test if the same relations that hold in extreme BL Lacs are observed in normal HSPs.

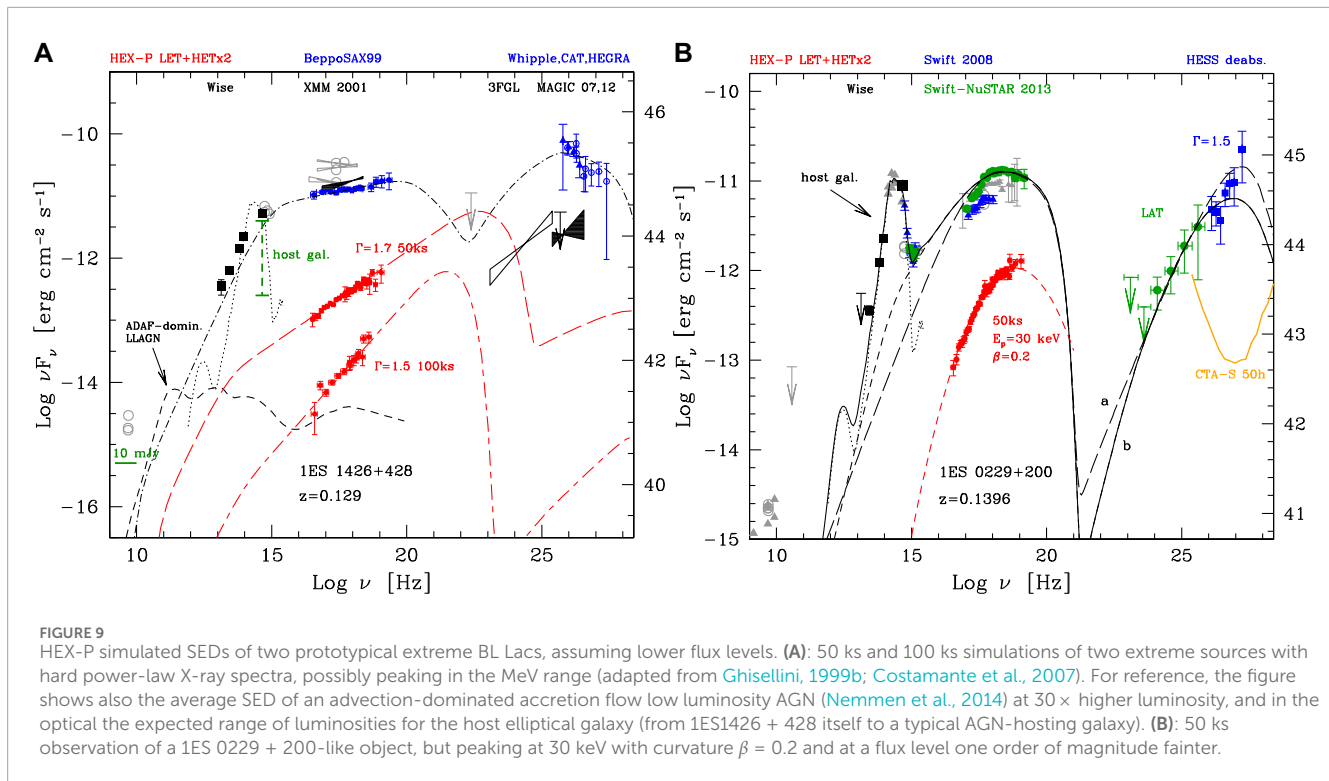
A crucial advancement enabled by HEX-P is in population studies of extreme blazars. Their overall luminosities tend to be low ( $\sim 10^{44-45}$  erg s<sup>-1</sup>) compared to regular HSPs, more or less in agreement with the blazar-sequence framework (Ghisellini et al., 1998; 2017). So far, they have only been found at low redshift (up to  $z \sim 0.3$ ). However, it is not clear if this is an intrinsic characteristic or simply a bias due to the low fluxes and limited observations.

With 50 ks exposures, HEX-P can discover and characterize such sources down to a 2–10 keV flux of  $F_{2-10 \text{ keV}} = 10^{-13}$  erg cm<sup>-2</sup>s<sup>-1</sup> out to  $\sim 30$  keV. For example, with a photon index  $\Gamma = 1.8$ , the 1-sigma error on the slope would be  $\pm 0.04$  full band, and  $\pm 0.20$  above 3 keV (vs.  $\pm 1.0$  with NuSTAR, for the same exposure time—and twice the clock time given the inefficiency of NuSTAR's low-Earth orbit compared to HEX-P's L1 orbit). With  $\Gamma = 1.6$  and  $F_{2-10 \text{ keV}} = 5 \times 10^{-14}$  erg cm<sup>-2</sup>s<sup>-1</sup>, an extreme BL Lac can still be detected up to  $\sim 20$  keV with a slope uncertainty of  $\pm 0.08$  full band and  $\pm 0.25$  above 3 keV (compared to  $\pm 1.2$  with NuSTAR). HEX-P can therefore discover and characterize sources similar to 1ES 1426 + 428 (i.e.,  $L \sim 10^{44}$  erg s<sup>-1</sup>) down to a luminosity of  $1 \times 10^{43}$  erg/s, or up to  $z = 0.7$ .

For objects with peak fluxes around  $10^{-12}$  erg cm<sup>-2</sup>s<sup>-1</sup>, HEX-P can also fully characterize the curvature of the synchrotron emission. For example, with a SED peaking at 30 keV and a curvature parameter  $\beta = 0.2$  (typical of HSPs), HEX-P can constrain the peak energy to better than 30% and the curvature to within 10% (Figure 9).

CTAO is expected to discover many more TeV-emitting HSPs, as allowed by a 10-fold improvement in sensitivity at 1 TeV with respect to present arrays. Just from the planned CTAO extragalactic survey, which will cover  $\sim 25\%$  of the sky, about 110–180 new detections are expected (Cherenkov Telescope Array Consortium et al., 2019), about 1/4 of which could be extreme-TeV according to the present number of extreme vs. standard objects among TeV-detected HSPs with known redshift (see e.g., Biteau et al., 2020; Costamante, 2020). Above a radio flux of 3.5 mJy at 1.4 GHz, the inferred surface density of  $\sim 4.5 \times 10^{-3}$  deg<sup>-2</sup> for extreme-synchrotron objects corresponds to a total number over the sky of  $\sim 200$  (Biteau et al., 2020). Searches for these extreme blazars will benefit from the eROSITA all-sky survey. This survey is expected to have a point-source sensitivity of  $6 \times 10^{-15}$  erg cm<sup>-2</sup>s<sup>-1</sup> in the 0.2–2 keV band (Brunner et al., 2022), making it more than ten times deeper than the ROSAT all-sky survey (Boller et al., 2016). Additionally, eROSITA will offer enhanced sensitivity in the 2–5 keV band, enabling the identification of further extreme blazar candidates. This will be particularly advantageous for low-luminosity





extreme HSP BL Lac objects, which can then be confirmed and characterized by HEX-P.

### 5.3.1 MeV-synchrotron BL lacs

In principle, blazars could reach even higher energies than shown by extreme BL Lacs. Expressing the acceleration time  $t_{acc}$  in terms of the gyroradius  $r_g$ ,  $t_{acc} = \eta(E) r_g/c$  (see e.g., Aharonian, 2004), where the factor  $\eta(E) \geq 1$  characterizes the energy-dependent rate of acceleration (where unity corresponds to the maximum possible rate in shock acceleration, when a particle gains all its energy in 1 gyroradius). Therefore, even extreme BL Lacs (where  $h\nu_{peak}/\delta \sim 1-10$  keV, with  $\delta$  being the beaming factor) are in fact *not* extreme accelerators ( $\eta \geq 10^4 - 10^5$ ). In the Crab nebula, for example, unbeamed synchrotron emission up to 10–20 MeV is observed, meaning electrons are accelerated at  $\sim 10\%$  of the maximum rate.

There might be blazars where the acceleration reaches even higher rates than observed so far, yielding a synchrotron peak in the MeV range (Ghisellini, 1999b; a). For example,  $h\nu_{peak} \sim 6-300$  MeV for  $\eta \approx 20-1000$  and  $\delta = 10-40$ . The luminosities are expected to be at the low end among blazars, if these objects follow the blazar sequence trend. Such objects could have easily escaped detection so far. At radio wavelengths, the flux is expected to be weak, likely below the typical level of large-area sky surveys or the selection cuts for blazar samples. In the optical band, the non-thermal continuum could be hidden one to two orders of magnitude below the host-galaxy thermal emission. In soft X-rays, the jet emission is expected to be low, thus possibly hidden under (or mixed with) the emission of the corona. In the VHE band, the SSC flux is likely too faint for detection, with the emission severely inhibited by Klein-Nishina effects. In summary, they could be faint at all frequencies except the

hard X-ray and MeV bands, where they emit all their power and could reach luminosities  $\sim 10^{43-45}$  erg s $^{-1}$ . As a result, even if already detected in present or future deep X-ray surveys, such sources might not be recognized as blazars, having been more likely classified as radio-quiet AGN.

The brightest among such objects could be detectable by COSI, but only HEX-P can pin-point and reveal the exceptional jet emission and extreme accelerator inside otherwise unremarkable objects. HEX-P can identify such a source with a peak luminosity of  $10^{43}$  erg s $^{-1}$  up to  $z = 0.5$ , with the observational signature of a hard X-ray spectrum ( $\Gamma = 1.5-1.7$ ), connecting by power-law to the possible COSI (or even Fermi-LAT) detections.

## 5.4 Masquerading BL lacs

Masquerading BL Lacs are objects that are classified as BL Lacs because of the lack of broad emission lines in their optical spectra, but are very likely FSRQs in disguise (Giommi et al., 2013). In these objects the jet is well aligned with our line of sight and its synchrotron emission overwhelms the disk and broad line region (BLR) emission, making the latter undetectable. Several spectroscopic and photometric campaigns have started collecting accurate redshift measurements of BL Lac sources detected in  $\gamma$ -ray catalogs, unveiling a few objects at high redshifts. ( $z > 0.8$ ; Shaw et al., 2013; Massaro et al., 2014; Paggi et al., 2014; Ricci et al., 2015; Álvarez Crespo et al., 2016a; b; Paiano et al., 2017; Marchesi et al., 2018; Desai et al., 2019; Kaur et al., 2019; Rajagopal et al., 2021; Sheng et al., 2022; Rajagopal et al., 2023). The community has shown renewed interest in them after TXS 0506 + 056, a BL Lac associated with the high-energy neutrino



IceCube-170922A, was revealed to be a masquerading BL Lac (Padovani et al., 2019). Indeed, neutrino emission in blazars via photo-hadronic interactions requires external photon fields, something that a masquerading BL Lac can provide. In the last few years, more masquerading BL Lacs potentially associated with IceCube neutrinos have been discovered (Padovani et al., 2022a; Padovani et al., 2022b). Distinguishing a BL Lac from a masquerading one is often complex and requires SED modeling with high quality data. This becomes somewhat easier when sources are detected at high redshift ( $z > 1$ ). In that case, the large radio power and Compton dominance gives away their true nature (Rau et al., 2012). High-redshift masquerading BL Lacs have also been dubbed ‘blue FSRQs’ (Ghisellini et al., 2012) because of their blue SED which typically shows synchrotron peak frequencies greater than  $10^{15}$  Hz and very hard  $\gamma$ -ray spectra. This implies that the electrons experience weak radiative cooling, a condition satisfied if the location of the emission region is outside the BLR (Ghisellini et al., 2012). The hard  $\gamma$ -ray spectra and large distances make masquerading BL Lacs powerful probes of the extragalactic background light (Ackermann et al., 2012; Fermi-LAT Collaboration et al., 2018), the integrated emission of all star-forming activity in the Universe.

HEX-P’s deep, broadband sensitivity is a unique advantage compared to other instruments for studying masquerading BL Lacs. By measuring the synchrotron component to high energies, HEX-P can put constraints on the maximum Lorentz factor ( $\gamma_{\max}$ ) of the electron population, at least up to the energy where the inverse Compton emission starts to dominate. This is something that is generally unconstrained given the limited sensitivity of X-ray instruments above 40 keV. In Figure 10 we show an example of a simulated SED for the high-redshift ( $z = 1.34$ ) BL Lac 4FGL J2146.5–1,344 studied by Rajagopal et al. (2020). In just 40 ks, HEX-P will be able to characterize the X-ray range in the SED of the source from 0.2 keV up to 30 keV, enabling us to accurately determine parameters such as the electrons’  $\gamma_{\max}$  and to detect a possible break in the high-energy spectrum of the source. We note that to achieve the same quality spectrum  $> 10$  keV, it would require  $> 100$  ks NuSTAR observations (and  $> 200$  ks clock-time, given NuSTAR’s low-Earth orbit). Even in sources which are more typically ISP much of the time, like TXS 0506 + 056, HEX-P’s broadband sensitivity will be crucial to identify epochs and states when the synchrotron emission enters the X-ray band, which typically yields enhanced gamma-ray activity in the VHE band through SSC emission. Constraining both the spectrum and the lower limit of  $\gamma_{\max}$ , HEX-P will provide fundamental information to elucidate the origin of possible VHE gamma-ray flares to be compared with epochs of higher neutrino fluxes.

## 5.5 Blazar evolution

As a source class, blazars have been shown to follow different evolutionary paths in the X-ray band. The powerful FSRQs evolve positively (i.e., there is an increasing number of sources at higher redshift, and/or the sources’ luminosity increases with redshift) up to redshift  $z > 4$  (e.g., Ajello et al., 2009; Marcotulli et al., 2022). BL Lac sources, fainter and for which redshifts are often tenuous,

appear not to evolve, or to evolve negatively in the X-ray band (Beckmann et al., 2003; Ajello et al., 2009).

The latest blazar evolutionary model in the hard X-ray band has confirmed that the peak of the evolution of the most luminous FSRQs ( $L_X > 10^{47}$  erg s $^{-1}$ ) lies beyond redshift 4 (Marcotulli et al., 2022). These sources are powered by the most massive black holes (see Section 5.2), and hence closely trace the evolution of the most massive black holes in the Universe. However, owing to the sensitivity limits of current X-ray surveys (e.g., Oh et al., 2018), we have yet to detect a turnover in the luminosity function at the lowest X-ray luminosities. This break, expected both as a result of beaming (Urry and Shafer, 1984) and due to the preferred evolutionary path, hampers the determination of how these sources evolve through cosmic time, i.e., whether there are more (density evolution) or more luminous (luminosity evolution) sources at earlier times. If we could tell the two evolution scenarios apart, this would have a profound impact on our current understanding of SMBH evolution. Moreover, studies based on radio-selected samples, found that blazar evolution follows the trend observed in the overall AGN population with a density peak at  $z \sim 2 - 3$  (Caccianiga et al., 2019; Ighina et al., 2021; Diana et al., 2022). Only by finding the bulk of the X-ray blazar population at  $z > 4$  would enable us to confirm the location of the peak, which in turn will inform us of the SMBH space density at those distances.

Using the latest hard X-ray luminosity function, we calculated the number of blazars detectable at different redshift bins in a 1 deg $^2$  survey, as planned for the HEX-P mission (Civano et al., 2023) with the predicted source sensitivity in the 8–24 keV band for the HET instruments. The right panel of Figure 11 shows these predictions based on the luminosity (pink hatched boxes) and density evolution model (green solid boxes). As can be seen, depending on the chosen evolution model, the number of sources detected at different redshifts is significantly different, with the luminosity evolution prediction extending up to  $z > 8$ . With its sensitivity, HEX-P would be able to answer these questions: are FSRQs evolving in density or luminosity? Where is the peak of the population? We note that blazars, or possibly nascent jets, have been already detected beyond  $z = 6$  (e.g., Zhu et al., 2019; Belladitta et al., 2020). The left panel of Figure 11 shows the predicted source count distribution ( $\log N - \log S$ ; i.e., number of sources as function of energy flux) of X-ray blazars in the 8–24 keV band as extrapolated from the X-ray best-fit luminosity evolution function in the same 1 deg $^2$  survey. With the capabilities of the HET, we will finally be able to detect the break in  $\log N - \log S$  and break the degeneracy that would lead us to understand for the very first time the type of evolution taking place in this source class.

Moreover, these sources, though contributing only  $\sim 2 - 10\%$  to the cosmic X-ray background above 14 keV, are predicted to be the major contributors to the MeV background. Assuming a luminosity evolution scenario, estimates are of the order of 70%, while for the density evolution, they are up to 100% (Marcotulli et al., 2022). However, such numbers are derived from extrapolations of the best-fit X-ray luminosity function to the MeV band, assuming an average high-energy (X-to- $\gamma$ -ray) SED of the MeV blazar class derived from time-averaged Swift/BAT and Fermi-LAT data. As shown in (Marcotulli et al., 2020), (Figures 3–5), quasi-simultaneous soft (Swift/XRT) and sensitive hard (NuSTAR) X-ray data, combined with  $\gamma$ -ray ones, are necessary to constrain the location (both

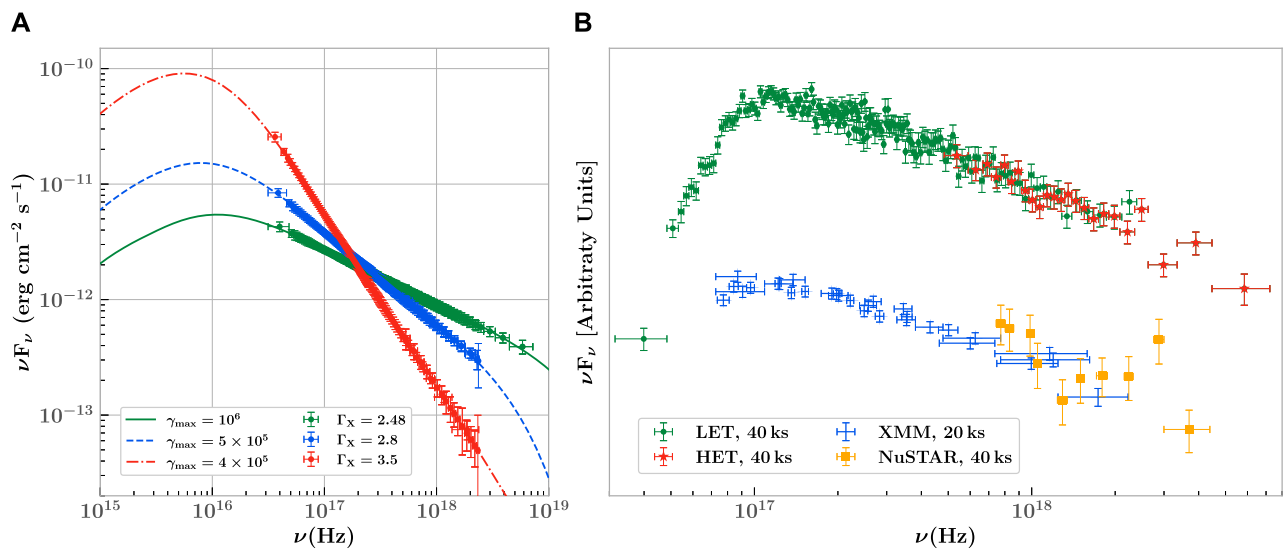


FIGURE 10

(A): Simulated SED of the high-redshift BL Lac 4FGL J2146.5–1,344 ( $z = 1.34$ ; Rajagopal et al., 2020). The left panel shows the simulated HEX-P SED with different spectral indices ( $\Gamma_x$ ). In only 40 ks, HEX-P will characterize the source spectrum from 0.2 keV to 30 keV. This will provide crucial constraints on the shape of the high-energy particle population (e.g., the maximum Lorentz factor,  $\gamma_{\text{max}}$ ), the total jet power, and the true nature of the emitter (whether it is a high-luminosity BL Lac or a masquerading FSRQ). (B): comparison of 40 ks HEX-P simulations to observations carried out by Rajagopal et al. (2020), which entailed simultaneous observations of 4FGL J2146.5–1,344 with XMM-Newton (20 ks) and NuSTAR (40 ks). HEX-P data quality for a similar exposure will highly enhance our S/N and improve determination of important spectral parameters (such as photon index and flux) as well as enable us to detect spectral curvature beyond 10 keV, which would require a > 100 ks NuSTAR exposure (and > 200 ks clock-time, given NuSTAR's low-Earth orbit).

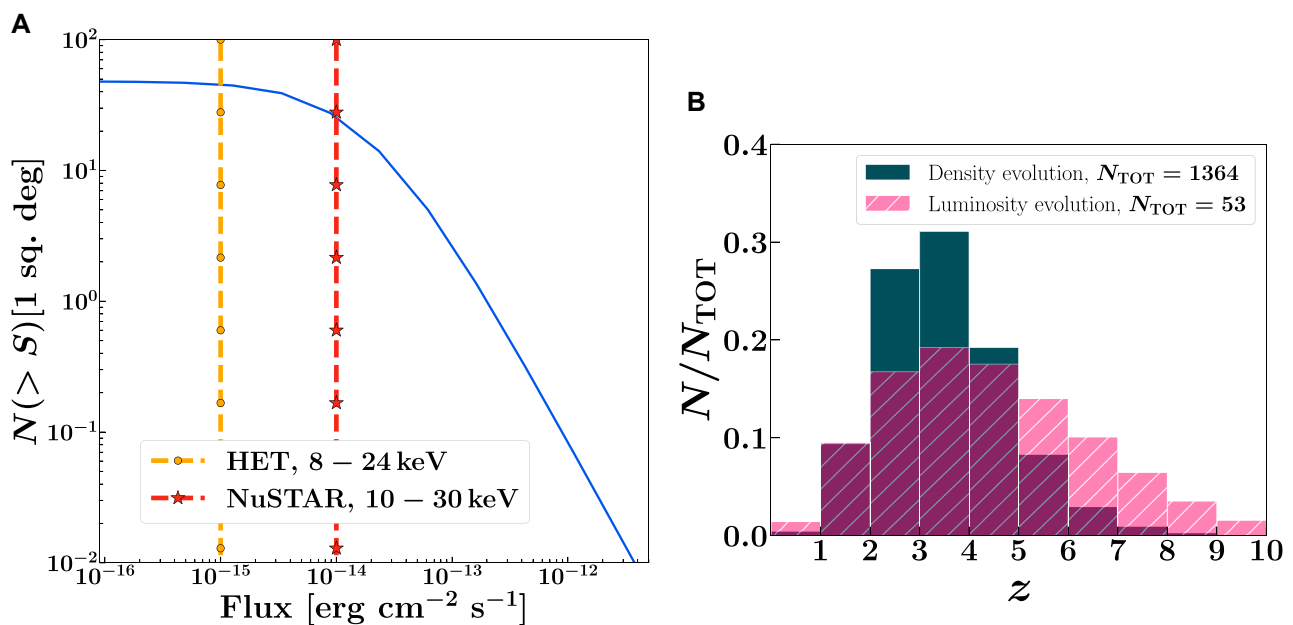


FIGURE 11

(A): Blazar source count distribution,  $\log N$ - $\log S$  (blue line), derived from the best-fit luminosity evolution model of Marcotulli et al. (2022) expected for a 1 deg<sup>2</sup> survey. The yellow dash-circle line indicates the HET 8–24 keV source sensitivity for the planned 1 deg<sup>2</sup> wide survey (Civano et al., 2023). For comparison, the red dash-star line represents the NuSTAR source sensitivity in the 10–30 keV band. HEX-P sensitivity would allow us for the very first time to detect the break of the luminosity function (which is yet to be seen by current facilities). This would lead us to definitely constrain which kind of evolution X-ray detected FSRQs have undergone, as measuring such break would solve the degeneracy between a density versus a luminosity evolution. (B): Predicted relative number of FSRQs ( $N/N_{\text{TOT}}$ ) that would be detected by HEX-P in the planned 1 deg<sup>2</sup> wide survey as a function of redshift. Prediction using the best-fit density evolution model are shown as solid dark green bars, and using the best-fit luminosity evolution model as hatched pink bars (Marcotulli et al., 2022). Depending on the number of sources detected by HEX-P as a function of redshift, we would be able to determine for the very first time the evolution properties of the population.

in flux and energy) of the high-energy SED peak, translating to constraints on the contribution of this class to the background. The average Swift-BAT uncertainty on the spectral index is  $\bar{\Gamma}_{\text{BAT, err}} = 0.2$  (Oh et al., 2018). A difference of 0.2 on the hard X-ray spectrum implies a factor of 10 difference in the contribution of blazars to the MeV background. Besides allowing us to detect the bulk of the high-redshift population, HEX-P, with its lower background at hard X-rays, will allow probing the high-energy SED very close to its peak ( $E \sim 50 - 80$  keV at  $z = 4$  corresponds to  $E_{\text{rest}} \sim 250 - 400$  keV and at  $z = 5$  to  $E_{\text{rest}} \sim 300 - 480$  keV wrt a peak at  $E_{\text{rest, peak}} = 3 - 17$  MeV, Marcotulli et al., 2022) pinpointing its location and the shape of the spectrum with uncertainties of the order of  $\Gamma_{0.2-70 \text{ keV, err}} = 0.01$ . This, together with the knowledge of the preferred evolution model, will provide tighter constraints on the contribution of MeV blazars to the cosmic MeV background, allowing us to constrain the uncertainties on the level of contribution of blazars to the MeV background from an order of magnitude down to 20%.

Similarly, the prospects to finally disentangle the preferred evolutionary paths for BL Lacs in the X-ray band are bright with HEX-P. Current photometric and spectroscopic follow-up of  $\gamma$ -ray detected BL Lacs are increasing the redshift completeness of this challenging source class. A sensitive X-ray telescope such as HEX-P would enable a complete sample of X-ray selected BL Lacs, even to lower luminosity (see Section 5.3.1), and determine for the first time their evolution in the full 0.2 – 80 keV band.

## 6 Resolving nearby jets

### 6.1 Blazars

Jets of active galaxies can be hundreds to thousands of kiloparsecs long (e.g., Oei et al., 2022). The extended structure is most easily observed in the radio band, where the low-energy synchrotron emission can be observed along the entire jet. However, many jets also show extended components (extended jets and/or hotspots along the jet flow) in other energy bands, including the X-ray domain. Observations with Chandra increased the number of resolved X-ray jets considerably and rekindled the debate on the X-ray emission processes in extended jets (see Harris and Krawczynski, 2006, for a review). Especially the explanation of hard X-ray spectra (harder than the radio-through-optical extrapolation) was strongly debated. Two main models emerged: Inverse-Compton scattering of CMB photons by low-energy electrons (which would produce the radio through optical emission through synchrotron emission; Tavecchio et al., 2000) or synchrotron emission of a second, high-energy electron population (e.g., Hardcastle, 2006; Jester et al., 2006). While the IC/CMB model was favored for some time, several studies over recent years have shown that this may not be the correct interpretation for many sources due to a lack of associated  $\gamma$ -rays (e.g., Breiding et al., 2017; Meyer et al., 2017).

For example, an extended X-ray jet has been observed in the flat-spectrum radio quasar 3C 273 (Marshall et al., 2001), revealing a  $\sim 10''$  extended structure with knots. No cut-off is observed in the X-ray spectrum. Multi-wavelength data suggest that the X-ray component is inverse-Compton emission (Uchiyama et al., 2006) or a second synchrotron component.

However, in a notable counterexample, the synchrotron explanation probably does not work for the blazar AP Librae (Zacharias and Wagner, 2016; Roychowdhury et al., 2022), one of the few BL Lac objects with a resolved,  $10''$  long X-ray jet (Kaufmann et al., 2013). The shape of the SED requires at least two emission zones. In fact, the VHE  $\gamma$ -ray emission observed with H.E.S.S. (H. E. S. S. Collaboration et al., 2015), has been interpreted as inverse-Compton scattering of dust emission from the inner 1kpc of the jet (Roychowdhury et al., 2022).

A recent study by Reddy et al. (2023) on jet knots revealed that in most knots the X-ray emission peak is located upstream of the radio emission peak. This might suggest a cooling effect on the particle flow across a shock pointing towards the synchrotron process in both energy bands, as the IC/CMB model requires a one-zone interpretation, where the fluxes should peak at the same location.

Hence, the debate continues. Observations with HEX-P will be crucial for studies of the hard X-ray emission of nearby, resolved blazar jets. The X-ray spectra do not show any spectral feature; they are described well with a power-law (e.g., Haardt et al., 1999; Marshall et al., 2001; Kaufmann et al., 2013). This suggests efficient particle acceleration in large-scale jets potentially up to the synchrotron-burnoff limit (de Jager et al., 1996). Unfortunately, the spectral range of the synchrotron-burnoff limit (50 MeV modulo beaming) is currently not well accessible, as the Fermi-LAT instrument cannot resolve extended blazar jets (e.g., Breiding et al., 2017) in order to derive the spectrum at these energies. The observation of a break or a cut-off in the hard X-ray domain would provide important information on the acceleration and radiation processes inside the extended jets. This would constrain both synchrotron and inverse-Compton models.

Using the 47 sources labeled as *BL Lac* or *CDQ* in the Atlas-X<sup>4</sup> (Reddy et al., 2023) and the XJET<sup>5</sup> catalogs (most of which are blazars), we determine that HEX-P's angular resolution of  $\leq 5''$  is sufficient to resolve up to 20 known (blazar) jets at hard X-rays. Six sources have jets longer than  $10''$  including, for example, 3C 273 (cf., Figure 12), 9 sources range between  $5''$  and  $10''$ , while 5 jets are about  $5''$  long which might still be detectable. To showcase the excellent imaging capabilities of HEX-P to resolve jets at different separations from the core and with different jet/core flux ratios, we performed several simulations using as a base image the blazar 3C 273 (Chandra obsid:4,876). All the images shown in Figure 12 have been produced using a 100 ks exposure with the LET instrument and they cover the 8 – 25 keV range. From the top panel to the bottom one, we increased the jet separation from the core from  $5''$  to  $20''$ . From the left panel to the right one, we increase the ratio of fluxes between jet and core ( $F_{\text{jet}}/F_{\text{core}}$ ) from 0.01 to 1. The real 3C 273 system as seen by HEX-P would be the left panel, second row. As can be seen, HEX-P imaging capabilities would be key to resolving blazar jets for the first time at  $> 10$  keV even when the jet is a few times fainter than the core and at very close separation. We would also note that the highest redshift ( $z > 5$ ) resolved X-ray jets are found at a separation of  $5''$  and

4 <https://astro.umbc.edu/Atlas-X/>

5 <https://hea-www.cfa.harvard.edu/XJET/>

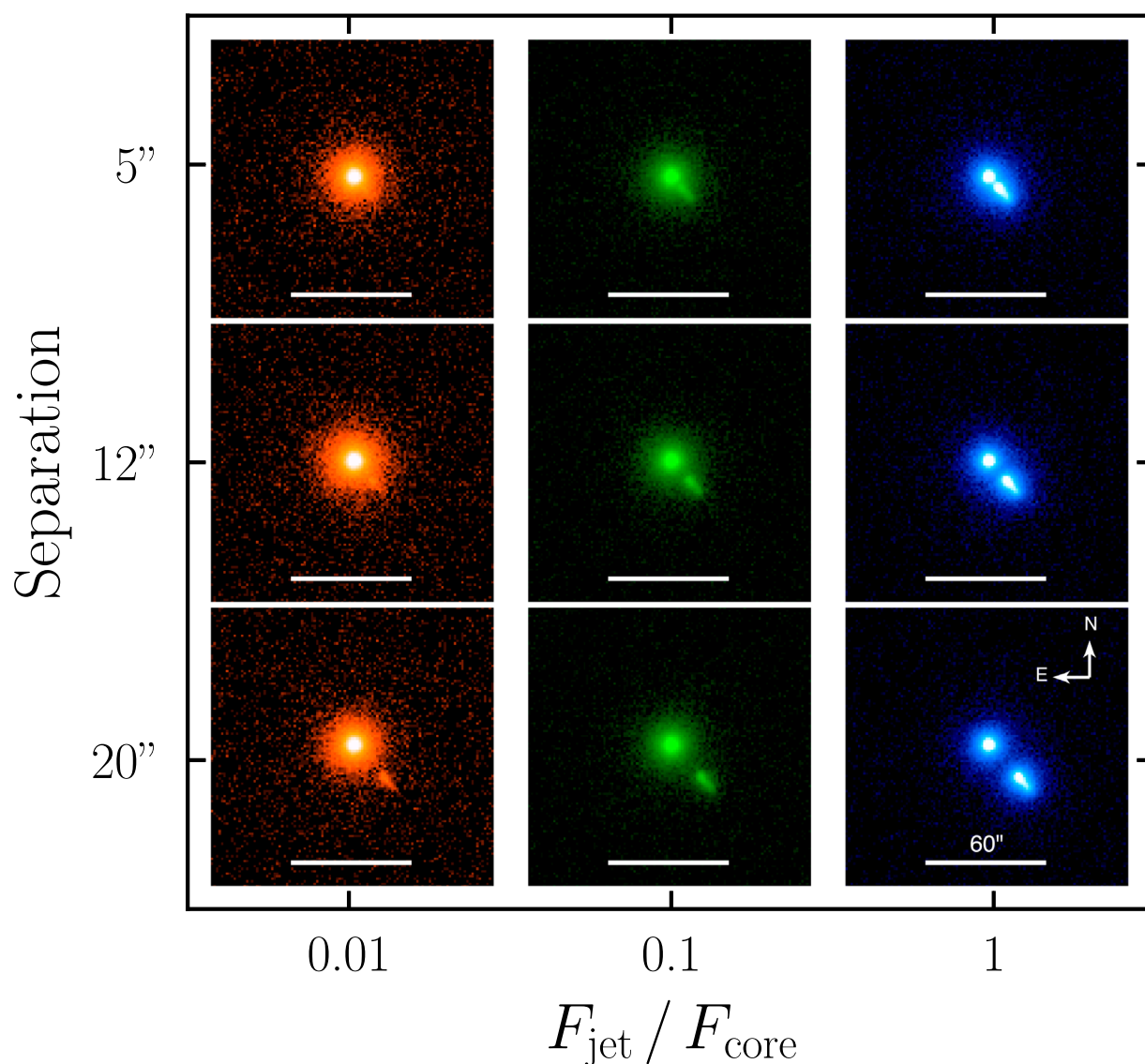


FIGURE 12

SIXTE simulations using as base image the blazar 3C 273 (Chandra obsid:4,876). All the images have been produced using 100 ks exposure with the LET instrument and they cover the 8 – 25 keV range. The field of view in all images is  $2' \times 2'$  and we have over-plotted a white line of length 60'' for scale. For reference, 60'' is the NuSTAR PSF. From top to bottom: the jet separation from the core increases from 5'' to 20''. From left to right: the ratio of fluxes between jet and core ( $F_{\text{jet}}/F_{\text{core}}$ ) increases from 0.01 to 1. The real 3C 273 system as seen by HEX-P would be the left panel, second row. All the other simulations showcase the excellent imaging capabilities of HEX-P to resolve jets at different separation from the core and with different flux ratios. As can be seen, HEX-P imaging capabilities would be key to resolve blazar jets for the first time  $> 10$  keV even when the jet is few times fainter than the core and at very close separation.

8'' from the core and a flux ratio  $F_{\text{jet}}/F_{\text{core}} \sim 0.06$  (Ighina et al., 2022) and  $F_{\text{jet}}/F_{\text{core}} < 0.1$  (Belladitta et al., 2020; Connor et al., 2021). Therefore, HEX-P would have the capability of resolving those jets.

This implies that HEX-P will help determine the spectral features (e.g., breaks or cut-offs, if present) at X-rays beyond 10 keV, which are so far unknown. In turn, the acceleration processes in extended jets will be strongly constrained, which will significantly enhance the modeling of these extended structures (Potter and Cotter, 2015; Lucchini et al., 2019; Zacharias et al., 2022), and which would greatly enhance the theoretical understanding of jets as a whole.

## 6.2 TeV-detected radio galaxies

Radio galaxies are typically non-aligned AGN with significantly reduced Doppler beaming. Depending on the X-ray production process, their extended X-ray jets may still be observable, as beaming is not as important depending on the particle energies in the jet. Here, we discuss a special class of radio galaxies, namely, those which are detected at VHE  $\gamma$ -ray energies. So far, six bona-fide radio galaxies are listed in the TeVCAT,<sup>6</sup> of which three have extended

<sup>6</sup> <http://tevcat2.uchicago.edu/>



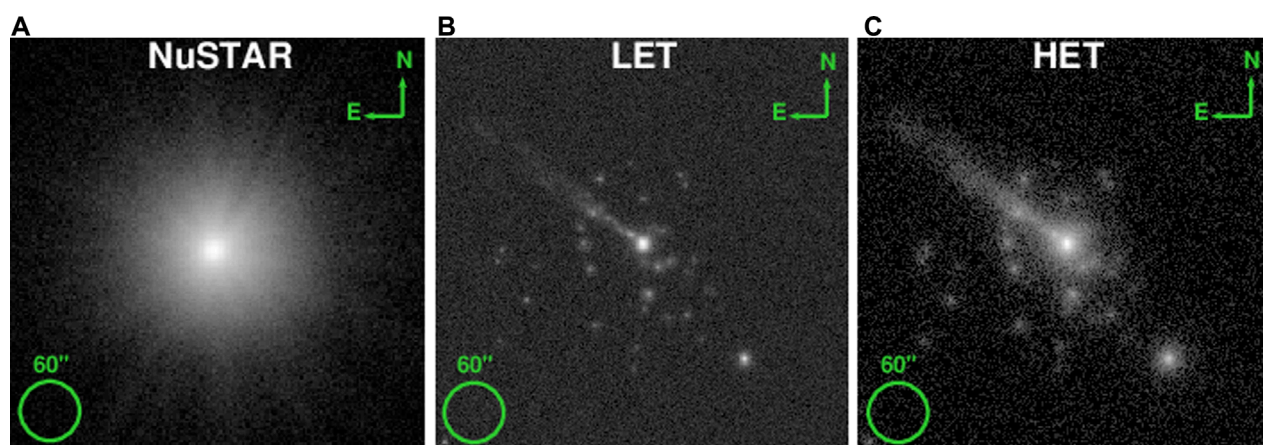


FIGURE 13

(A): NuSTAR FPMA 3–80 keV image of the radio galaxy Centaurus A (obsid: 60001081002). (B, C): SIXTE simulations of the same galaxy with HEX-P based on the Chandra image of the source (Middle: 500 ks observation in the 8–20 keV range with the LET; Right: 500 ks observation in the 20–50 keV range with the HET). HEX-P would enable us, for the first time, to resolve extended jets in the hard X-ray band as well as study point sources visible in the Chandra image at soft X-rays (see [Lehmer et al., 2023](#)). The images are all  $7' \times 7'$  and we have over-plotted a green circle of diameter  $1'$  for scale.

X-ray jets: M 87 ([Aharonian et al., 2006](#); [Acciari et al., 2009](#); [Abramowski et al., 2012](#)), Centaurus A ([Aharonian et al., 2009b](#); [H. E. S. S. Collaboration et al., 2020](#)), and 3C 264 ([Archer et al., 2020](#), however, the X-ray jet is less than  $5''$  long). Centaurus A is a special case, as it is not only the closest AGN to Earth, but also, so far, the only AGN where the TeV emission is spatially extended along the axis of the X-ray and radio jet ([H. E. S. S. Collaboration et al., 2020](#)). This could indicate the presence of high-energy particles on kpc scales that are inverse Compton scattering dust emission or starlight, or it could be associated with synchrotron radiation from ultra-high energy electrons (e.g., [Wang et al., 2021](#)).

Observations with HEX-P of these extended jets, coupled with the knowledge that these jets emit TeV  $\gamma$  rays at large distances from the black hole, will be crucial for constraining the location and physics of the  $\gamma$ -ray production. As mentioned above, cut offs or breaks in the X-ray spectra would reveal maximum particle energies (or at least the acceleration efficiency) indicating that such a particle population would probably not produce  $\gamma$  rays. Determining spatially resolved spectra of nearby jets (such as Centaurus A; see [Figure 13](#)) would allow us to pinpoint the regions where the extended  $\gamma$ -ray emission is produced.

Along with HEX-P and its broad X-ray coverage, multiwavelength observations with future high-angular-resolution instruments will be crucial to unlock the secrets of these objects. Most notably, collaborative observations by HEX-P with CTAO will be important to constrain the jet morphology and processes.

## 7 Conclusion

The HEX-P mission offers several key strengths for studying jetted AGN. Its broad-band sensitivity, spanning from 0.2 to 80 keV, provides deep coverage in relatively short exposure times. Additionally, HEX-P exhibits excellent angular resolution

at high energies and the unique capability to point in the antisolar direction. These remarkable capabilities set HEX-P apart from both current and currently planned X-ray observatories, making it an unparalleled asset for the study of blazars in the next decade. In particular:

- HEX-P's potential extends to time-domain and multi-messenger studies. HEX-P has the capability to independently discover the expected signatures of hadronic accelerators. In [Section 4.1](#), we illustrated how HEX-P can detect spectral changes during an X-ray flare as expected from the presence of a hadronic component. Moreover, by observing blazars that are detected by IceCube, HEX-P can identify the X-ray signatures expected from hadronic sources. In [Section 4.2](#), we demonstrated that HEX-P can detect the signature of proton-induced cascades associated with neutrino emission. It accomplishes this by distinguishing between power-law emission, typical of a leptonic synchrotron component, from broken power laws, expected in the case of proton-induced cascades. Note, however, that these signatures are not unique to hadronic processes and may also arise in multi-component leptonic emission scenarios.
- The deep broad-band sensitivity of HEX-P will enable the study of rapid variability on minute timescales. As shown in [Section 4.3](#), these investigations provide insight into the connection between the acceleration and cooling of charged particles. Previously, such studies were limited to only a few objects over an energy band spanning less than 1 decade. HEX-P will expand this research to a significantly larger number of objects and offer an energy band spanning more than 2 decades. Furthermore, HEX-P's ability to pinpoint extremely fast variability, comparable to the light-crossing time of the system, will reveal invaluable information about particle acceleration mechanisms. In [Section 4.3](#), we demonstrated how HEX-P can measure the lightcurve of a fast flaring blazar (e.g., Mrk 421) at both low ( $< 10$  keV) and high ( $> 10$  keV)



X-ray energies, thereby identifying the signature of magnetic reconnection.

- Section 5 showed how HEX-P will allow investigation of the accretion-jet connection by measuring the average X-ray SED of selected blazars and constrain variation in the SED linked to possible accretion fluctuations. Moreover, HEX-P will detect and study MeV blazars (the most luminous type of blazars) well above the current  $z \approx 3$  limit and possibly all the way to  $z = 6$ .
- As shown in Section 6, for the first time above 20 keV, HEX-P will be able to resolve jets of blazars and radio galaxies, thus pinpointing the regions where the extended emission is produced. HEX-P will be able to discriminate between the emission of the core and that of the jet allowing to constrain potential spectral cut-offs in the jet. This would yield an important constraints on the maximum energy of the particle population.

Lastly, HEX-P exhibits strong synergies with current and future planned facilities, including CTAO, IceCube, IXPE, COSI, and the current and future CMB facilities (Section 4.4). These collaborations enable the study of extremely fast flaring blazars, the detection and study of multi-messenger sources, and the investigation of hadronic emission components in blazars. An early 2030s launch of HEX-P will significantly enhance the science provided by those facilities and answer many key questions about the most powerful jets in the Universe.

## Data availability statement

The original contributions presented in the study are included in the article, further inquiries can be directed to the corresponding author.

## Author contributions

LM: Writing–original draft, Writing–review and editing. MA: Writing–original draft, Writing–review and editing. MB: Writing–original draft, Writing–review and editing. PC: Writing–review and editing. LC: Writing–original draft, Writing–review and editing. LG: Writing–review and editing. ME: Writing–original draft, Writing–review and editing. JG: Writing–review and editing. AG: Writing–original draft, Writing–review and editing. IL: Writing–original draft, Writing–review and editing. GM: Writing–original draft, Writing–review and editing. KM: Writing–review and editing. AM: Writing–review and editing. RM: Writing–review and editing. FM: Writing–original draft, Writing–review and editing. MP: Writing–original draft, Writing–review and editing. BR: Writing–review and editing. TS: Writing–original draft,

Writing–review and editing. DS: Writing–review and editing. GV: Writing–original draft. MZ: Writing–original draft, Writing–review and editing. HZ: Writing–original draft, Writing–review and editing.

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## Conflict of interest

The authors declare that the research was conducted in the absence of any commercial or financial relationships that could be construed as a potential conflict of interest.

The reviewer FT declared a shared affiliation with the author(s) AM, MR and TS to the handling editor at the time of review.

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