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RESEARCH ARTICLE

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Key Points:

- MODIS Dark Target Ocean aerosol retrieval algorithm has a long-standing retrieval bias for dust aerosol
- This study implements a nonspherical dust aerosol model to better represent dust optical properties
- The new algorithm reduces biases in AOD, AE, and FMF in dusty pixels and causes an overall reduction in AOD in major dust regions over ocean

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Dust Aerosol Retrieval Over the Oceans With the MODIS/VIIRS Dark Target Algorithm: 2. Nonspherical Dust Model

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Abstract The Dark-target (DT) aerosol algorithm retrieves spectral Aerosol Optical Depth (AOD) and other aerosol properties from Moderate-resolution Imaging Spectrometer (MODIS) reflectance observations. Over the ocean, the DT algorithm is known to contain scattering-angle-dependent biases in its retrievals of AOD, Angstrom Exponent (AE), and Fine Mode Fraction (FMF) for dust aerosols. Following a two-step strategy to improve the DT retrieval of dust over ocean, for which the first step is to identify dusty pixels (reported in “Part 1”), in this “Part 2,” we report on construction of a new dust model lookup table (LUT) and the strategy for applying it within the existing DT algorithm. In particular, we evaluate different characterizations of dust optical properties from a variety of frameworks and databases, and compare them with the current DT retrieval assumptions. Substituting the standard operational LUT with a spheroid dust model with identified dusty pixels shows significant improvement when compared with collocated AERONET-identified dusty pixels. Specifically, the application of the new dust model to dusty pixels reduces their AOD bias from 0.06 to 0.02 while improving the fraction of retrievals within expected error from 64% to 82%. At the same time, the overall bias in AE is reduced from 0.13 to 0.06, and the scattering-angle-dependent AE bias is largely eliminated. In testing on two full months of data (April and July), the new retrieval will reduce the monthly mean AOD by up to 0.1 and 0.2 in the north Atlantic and Arabian seas, respectively. The average AE and FMF are also reduced in these dust heavy regions.

1. Introduction

Dust aerosols in the atmosphere can pose great impacts on weather, climate, air quality, public health, and ecosystems (Goudie, 2013; Griffin, 2007; Harrison et al., 2001; Kaufman et al., 2002; Kok et al., 2018; Miller & Tegen, 1998; Rosenfeld et al., 2011; Shao et al., 2013; Shell & Somerville, 2007; Yu et al., 2006; Zhao et al., 2011). Dust retrieval with satellite remote sensing is very difficult because dust particles have a wide range of sizes, shapes, and mineralogical compositions, depending on the source region and meteorological conditions. The size, shape, and chemical compositions of particles determine their absorption and scattering properties and the radiances received by remote sensing instruments that retrieval algorithms rely upon.

The primary instruments for aerosol observations (including dust) in space are multichannel spectroradiometers such as the Advanced Very-High-Resolution Radiometer (AVHRR) (Zhao et al., 2013), the Moderate-resolution Imaging Spectrometer (MODIS) (King et al., 1992), Visible Infrared Imaging Radiometer Suite (VIIRS) (Sawyer et al., 2020) in the United States, and the Medium Resolution Imaging Spectrometer (MERIS) (Antoine & Morel, 1999), Advanced Along-Track Scanning Radiometer (AATSR), the Sea and Land Surface Temperature Radiometer (SLSTR) by the European Space Agency, as well as instruments flying on geostationary satellites such as ABI (Advanced Baseline Imager) onboard on Geostationary Operational Environmental Satellite (GOES) and Advanced Himawari Imager (AHI) on HIMAWARI 8 and 9 (Gupta et al., 2019). The common feature of these instruments is that they observe at multiple channels in the visible and near-IR wavelengths that are sensitive to the scattering and absorption properties of aerosols. More than one aerosol retrieval algorithm exist for some of the instruments such as MODIS and AATSR (de Leeuw et al., 2011; Kokhanovsky & de Leeuw, 2009). Most of the retrieval algorithms have separate dust aerosol models but not many have considered nonsphericity of dust particles,

especially over the oceans (de Leeuw et al., 2015). Among these algorithms, the Deep Blue (DT) algorithm on VIIRS has recently added a nonspherical dust model (Lee et al., 2017) but NOAA's operational aerosol algorithm for VIIRS is based on the standard MODIS Dark Target (DT) Ocean algorithm (Remer et al., 2005) that does not implement a nonspherical dust model. Nonsphericity of dust is considered in the retrieval algorithm of the Multiangle Imaging Spectro Radiometer (MISR) aboard the Terra satellite platform. MISR measures radiance over nine camera angles at four wavelengths (446, 558, 672, 867 nm) and its multichannel and multiangle capabilities enable more accurate depiction of aerosol scattering properties, size information, and nonsphericity of aerosol particles (Kahn et al., 2007). However, the relatively narrow swath width (~360 km) of MISR results in a global coverage every 7 days, compared with 1–2 days by MODIS and 1 day by VIIRS. In addition, nonspherical dust models have been considered in the UV-based algorithm of the Ozone Monitoring Instrument (OMI) and Ozone Mapping and Profiler Suite (OMPS) (Torres et al., 2002, 2013) and the new PARASOL (Polarization & Anisotropy of Reflectances for Atmospheric Sciences coupled with Observations from a Lidar) retrieval, which utilizes the Generalized Retrieval of Aerosol and Surface Properties (GRASP) algorithm (Dubovik et al., 2006, 2014). Thus, the time is ripe for the Dark Target algorithm to implement a nonspherical option for retrievals of dust aerosol.

The MODIS instruments on board NASA's EOS Terra and Aqua satellites have been observing the Earth's cloud, aerosol, and surface since their launches in 2000 and 2002, respectively. The DT retrieval algorithm, which exclusively obtains aerosol properties over “dark targets,” provides several official MODIS aerosol products. It actually consists of two independent algorithms, one for ocean (DT-O; Levy et al., 2013; Remer et al., 2005; Tanré et al., 1997), and one for dark land surfaces such as vegetation (DT-L; Levy, Remer & Dubovik, 2007; Levy, Remer, Mattoo, et al., 2007; Levy et al., 2010, 2013). The DT algorithm (both land and ocean) follows a lookup table (LUT) approach, i.e., the top-of-atmosphere (TOA) spectral reflectance is precalculated using scattering and radiative transfer (RT) codes (Ahmad & Fraser, 1982; Dubovik et al., 2002; Evans & Stephens, 1991; Wiscombe, 1980) for predefined aerosol, surface, and atmospheric properties. Current ocean LUTs contain four fine aerosol models and five coarse aerosol models. The DT-O retrieval algorithm selects one fine mode and one coarse mode with an adjustable fraction from each mode to minimize the difference between the LUT reflectance and observed TOA reflectance in six wavelengths (0.55, 0.65, 0.86, 1.2, 1.6, and 2.11 μm). Dust aerosols are predominantly nonspherical in shape (Chou et al., 2008; Kalashnikova & Sokolik, 2002). However, none of the current DT-O aerosol models are nonspherical, including those meant to represent dust optical properties.

Furthermore, simply adding a nonspherical model or replacing an existing coarse model will not solve the problem, as there is no guarantee that the algorithm will reliably choose nonspherical models for dust and spherical models for nondust aerosol, possibly further degrading the results of the retrieval. Therefore, a two-step dust retrieval strategy is designed to first detect the dust and then apply a nonspherical dust model to identified dusty pixels. Part 1 of this series focused on dust detection (Zhou et al., 2020). In this paper, we focus on constructing a nonspherical model and LUT for dust aerosol retrievals.

The paper is organized as follows: section 2 describes the data used in this work. Section 3 describes construction of the dust optical models and LUT. Section 4 evaluates the performance of retrievals using the new dust models. Section 5 shows global aerosol retrieval in two dust heavy months. Section 6 summarizes the current work and plans for future improvement.

2. Data

The primary data used in the study are MODIS Level 1B calibrated TOA reflectance products (MxD02) (<http://mcst.gsfc.nasa.gov/content/l1b-documents>) and Collection C6 Level 2 aerosol products MxD04 (Levy et al., 2013) (where x is substituted by O for Terra and Y for Aqua). The MxD02 contains native resolution reflectance and radiance data (e.g., 0.25, 0.5, or 1.0 km at nadir, depending on band), whereas the DT algorithm aggregates these observations into NxN boxes. The MxD04_L2 products used here include Aerosol Optical Depth (AOD), Angstrom Exponent (AE), and Fine Mode Fraction (FMF) reported at 10 \times 10 km (nadir) resolution. In addition to the retrieved aerosol properties, a MxD04 data point contains the clear-sky reflectance values used in the retrieval, quality assurance (confidence) estimates, as well as other ancillary information such as 2-m surface wind speed from National Centers for Environmental Prediction (NCEP) reanalysis.

The “clear-sky reflectance” values in the MxD04_L2 files provide the average reflectance for the NxN aggregation of native-resolution data, after it has been screened for clouds, surface inhomogeneities (glint patterns, underwater sediments, etc.), and outliers. In addition, the clear-sky reflectances have been adjusted to remove the effects of trace gas absorption. In other words, the resulting clear-sky reflectance (in each of the six wavelength bands) is the gas absorption-corrected mean of the data that has not been masked or subject to statistical screening. The new dust retrieval uses the “clear-sky reflectance” from MxD04 as input to the DT algorithm. From this point forward, when we say “retrieved pixel,” we mean the 10×10 km resolution aggregation reported in the MxD04 product or the new retrieval at the same spatial resolution.

Aerosol measurements from the ground-based Aerosol Robotic Network (AERONET) sun photometers are commonly used for validating satellite aerosol retrievals (Gupta et al., 2018; Huang et al., 2016; Remer et al., 2005; Sayer et al., 2018; Torres et al., 2013). AERONET is a global network of Cimel Electronique CE-318 sun-sky radiometers with between four and nine spectral channels. In a common spectral configuration, AOD is obtained from direct sun measurements at 0.34, 0.38, 0.44, 0.50, 0.67, 0.87, and $1.02 \mu\text{m}$ with frequency of every 15 min; The high accuracy of AERONET AOD (estimated errors of ~ 0.01 – 0.02) makes it widely popular (Eck et al., 1999; Holben et al., 1998). In addition to direct sun measurement, the instruments measure the sky radiance in four spectral bands (0.44, 0.67, 0.87, and $1.02 \mu\text{m}$) along the solar principal plane up to nine times a day and along the solar almucantar up to eight times a day. Measurements from almucantar scans are used to retrieve aerosol particle size distribution, spectral complex refractive index, and single scattering albedo (SSA) (Dubovik et al., 2006; Dubovik & King, 2000). The majority of traditional AERONET stations are situated inland, with some located on islands or near the coast, and very few in the middle of the ocean. To better test our ocean algorithm, we supplement these land-based observations with Sun photometer measurements made from aboard ship cruises as part of the AERONET Maritime Aerosol Network (MAN; Smirnov et al., 2009). In this work, we use the latest Version 3 products, which provide improved cloud screening and better identification of heavy aerosol events (Giles et al., 2019).

To match the MODIS retrievals with the AERONET and MAN measurements, AERONET or MAN observations within ± 30 min of the MODIS overpass time are averaged and compared against the values for all the MODIS-retrieved pixels located within 0.3° of AERONET stations (Shi et al., 2019). In this way there are many collocations reported at each AERONET or MAN location at every overpass opportunity, since the AERONET or MAN observations are averaged to a single value but the many MODIS-retrieved pixels within the match-up circle are not. The time period of the MODIS-AERONET collocation data spans from the beginning of the MODIS Terra and Aqua missions, which began in 2000 and 2002, respectively, until 2014, while the MODIS-MAN collocation period covers the years from 2004 to 2014. AERONET AOD measurements at 0.44 and $0.67 \mu\text{m}$ are linearly interpolated in log-log space (i.e., fit by AE) so that they may be compared with MODIS AOD at $0.55 \mu\text{m}$. Comparisons are only performed on pixels likely to be dominated by dust aerosol, which are identified by means of thresholds applied to the AERONET or MAN measurements ($\text{AOD} > 0.3$, $\text{AE} < 0.6$).

3. Constructing a Dust Aerosol Model

Part 1 of this work reports on a novel dust detection algorithm. It combines spectral tests from near-UV (deep blue), visible, shortwave infrared (SWIR), and thermal infrared (TIR) wavelengths selected from a survey of existing dust detection algorithms. For each 10-km MxD04 retrieval pixel, there are 10×10 1-km pixels that the dust detection algorithm is applied to individually, using L1B reflectance and BT data from MxD02. The 10-km retrieval pixel is considered dusty if at least three 1-km pixels are identified as dusty by the new dust detection algorithm. Quantitative evaluation of this new dust detection algorithm using collocated ground-based sun photometer (AERONET) as well as space-based lidar (CALIOP on CALIPSO) finds it detects about 30% of dusty pixels in weak aerosol loading and more than 80% of heavy dust pixels, with very few false positive detections. The dust aerosol model derived in this paper is only applied to pixels that have been identified as dusty by this dust detection algorithm.

Once the dust-dominated pixels are identified, the next step is to select a proper dust model and compute the corresponding LUT that can be used in the retrieval. In the current DT framework, the vector RT code of Ahmad and Fraser (1982, AF82) was used to create the LUTs. AF82 provides an internal Mie scattering

code to compute spherical aerosol optical properties if particle size distribution and refractive indices are specified. However, alternative optical properties (scattering/extinction coefficients, phase matrices, etc.) can also be fed directly into AF82 in order to create new nonspherical LUTs. Our approach is to find a nonspherical aerosol model that is representative of most dust conditions in the real world. We compute the optical properties of the dust model outside of AF82 and then insert them into AF82 so that the rest of RT calculations are consistent with other DT aerosol models.

3.1. Dust Optical Properties

Dust in the atmosphere consists of a mixture of minerals with different sizes, shapes, densities, hygroscopicities, and chemical compositions (Sokolik & Toon, 1999). Optical properties (scattering and absorption) of dust are determined by these physical and chemical properties. In turn, the physical and chemical properties of the dust depend on the region of origin, how it was mobilized, and the transformation processes during dust transport. Since it is nearly impossible to prescribe a unique dust model for each dust sample or dust region, we seek a more generic dust model that represents the dust ensembles found in the real atmosphere.

The optical properties of an ensemble of particles can be integrated as

$$Q = \int_{r_{\min}}^{r_{\max}} \int_{\epsilon_1}^{\epsilon_n} C(r, \epsilon, r_e, r_i, \dots) \text{wt}(\epsilon) \frac{dN(r)}{dr} dr d\epsilon, \quad (1)$$

where C represents the relevant optical property (i.e., extinction, absorption coefficient, or phase function) of the individual particles and is a function of refractive index (r_e , r_i), size (r), and shape (ϵ); $\text{wt}(\epsilon)$ is the weighting function (here it is assumed to only change with shape); and $dN(r)/dr$ is the number size distribution (Dubovik et al., 2002) which we assume to take the form of a log-normal distribution:

$$\frac{dN(r)}{dr} = n(r) = \frac{1}{\sigma\sqrt{2\pi}} \frac{1}{r} \exp\left(-\frac{(\ln r - \mu)^2}{2\sigma^2}\right). \quad (2)$$

Here μ and σ are the mean and standard deviation of $\ln(n(r))$, respectively, which follows a normal distribution by definition. For a nonspherical particle, r can be taken as the equivalent radius of a spherical particle that has equal surface area or equal volume. It should be noted that, for a fixed population of particles, the value of u will change depending on whether it is represented in terms of volume or area equivalency (Wyser, 1998).

Dust particles are often approximated by spheres (i.e., in the DT ocean algorithm, Levy et al., 2010, 2013; Remer et al., 2005, 2008) or spheroids (Dubovik et al., 2002, 2006). Dubovik et al. (2006) derived an ensemble of 25 spheroid particle shapes with aspect ratios ranging from 0.3 to 3.0 to improve the AERONET retrieval. The same particle shape distribution is used in the GRASP (Generalized Retrieval of Aerosol and Surface Properties; Dubovik et al., 2014) model, which is developed from Dubovik et al. (2002, 2006), but is more flexible and capable of characterizing surface and atmospheric properties from a variety of remote sensing observations. In addition, recent advances in RT theory and computational methods have made possible the computation of optical properties of more complicated particle shapes (Bi et al., 2010; Liu et al., 2013; Tang et al., 2019). The Texas A&M single-scattering properties database consists of precomputed optical properties of a wide range of triaxial ellipsoidal particles with different refractive indices, sizes, and shape parameters (Meng et al., 2010) that allow users to compute the optical characteristics of a specific aerosol ensembles of their choice. We have conducted many sensitivity studies to determine the size distribution, shape combination, and optical database most suitable for our dust model.

An initial study has been conducted to examine the sensitivity of optical properties to particle shape for different aerosol sizes. We start by replacing each of the spherical particles in the fine (Model 1, Figure 1a) and coarse (Model 8, Figure 1b) mode models of the standard DT algorithm with a nonspherical particle of equivalent projected area. Furthermore, in this new aerosol model, each nonspherical particle maintains the same refractive index as its spherical counterpart (Table 1). Since the Texas A&M (TX) database provides a wide range of complex particle shapes, we chose to test three different shapes from the database: (1) spheres (TX-Sphere), (2) a mixture of spheroid shapes that are consistent with what has been used in the GRASP model (TX-Sphd), and (3) a mixture of ellipsoids (TX-Ellipd). The ellipsoid particles can be

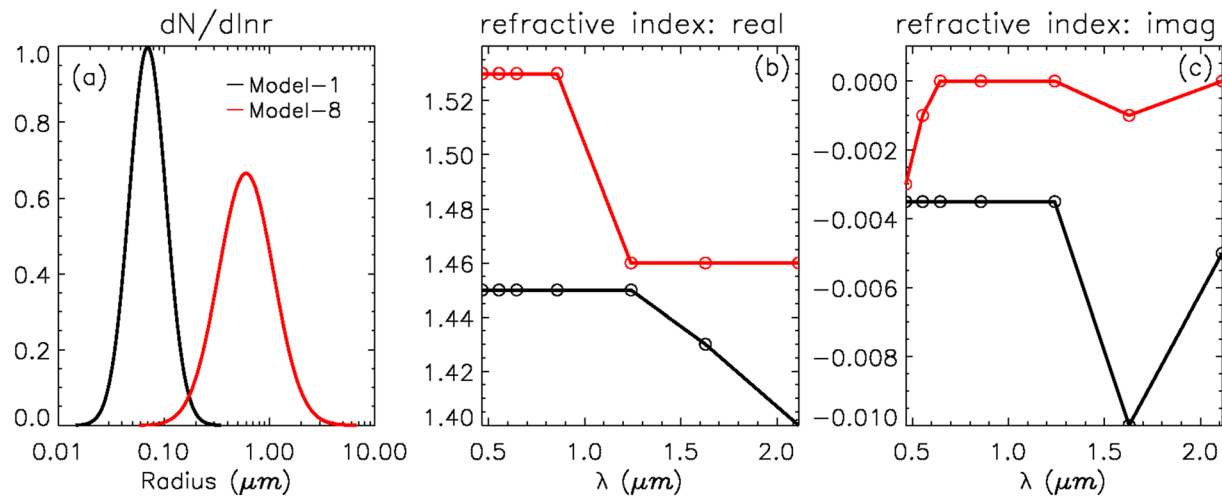


Figure 1. The number density distribution (a) and real (b) and imaginary (c) parts of the refractive index of DT aerosol Model-1 (black curves) and Model-8 (red curves).

described with three pairwise perpendicular axes of symmetry and two aspect ratios with spheroid being a special case of ellipsoid where two of the axes are equal. The shape and weight of TX-Sphd and TX-Ellipd are listed in Table 2. As mentioned earlier, the current DT algorithm assumes all particles to be spherical and excellent agreement was found between the single scattering properties of the DT and TX-sphere cases (Figure not shown).

For the fine mode, Figure 2 shows that particle shape does not have a large impact on the optical properties at the seven MODIS wavelengths (0.47, 0.55, 0.65, 0.86, 1.24, 1.63, and 2.11 μm) used in the DT aerosol retrievals. The extinction cross sections (CEXT) from the TX-Sphere (not shown) and original DT are almost identical and only slightly smaller than the values corresponding to the spheroid (TX-Sphd) and ellipsoid (TX-Ellipd) particle shapes (Figure 2a). Furthermore, the spectral shapes of CEXT are almost indistinguishable among the three particle morphologies (Figure 2b). The phase functions (P_{11}) are also very similar, with TX-Ellipd showing slightly more backscattering than the TX-Sphere and TX-Sphd shapes (Figure 2c). The only relatively large difference appears in the degree of linear polarization ($-P_{12}/P_{11}$) near its peak at scattering angles around 90° (Figure 2d). Since the MODIS instruments do not explicitly measure the polarization state of the observed scene and P_{12}/P_{11} has only a small impact on the total reflectance LUT, it is not necessary to add new nonspherical dust particles to the existing fine mode aerosol model.

However, for the coarse mode, CEXT, phase function, and the degree of linear polarization all have significant sensitivity to particle shape (Figure 3). A large difference can be observed between the TX-sphere and nonspherical particle phase functions for scattering angles greater than 80° (Figure 3c). The particles with nonspherical shapes have much flatter phase functions in these scattering angle ranges than the spherical particles, resulting in larger phase function values between 80° and 140° and smaller phase function values at scattering angles greater than 140° . These phase function differences between spherical and nonspherical particles are the main cause of scattering-angle-dependent-bias in DT retrievals of dust aerosol (shown later).

Table 1
Aerosol Models (DT Model-1 and Model-8, Model-10)

	$\lambda = 0.47 \geq 0.87 \mu\text{m}$	$\lambda = 1.24 \mu\text{m}$	$\lambda = 1.63 \mu\text{m}$	$\lambda = 2.11 \mu\text{m}$	$r_g (\mu\text{m})$	$\sigma (\mu\text{m})$	$r_{\text{eff}} (\mu\text{m})$
Model-1	1.45–0.0035i	1.45–0.0035i	1.43–0.01i	1.40–0.005i	0.07	0.40	0.10
Model-8	1.53–0.003i (0.47) 1.53–0.001i (0.55) 1.53–0.000i (0.65) 1.53–0.00i (0.86)	1.46–0.000i	1.46–0.001i	1.46–0.000i	0.60	0.60	1.48
Model-10	Same as M8				0.5	0.6	1.23

Table 2
The Shape and Relative Weighting of Particles in TX-Sphd and TX-Ellipd Experiments Used in Figures 2 and 3

Spheroid particles in TX-Sphd		Ellipsoid particles in TX-Ellipd		
ϵ	$N(\epsilon)$	$\epsilon_{a/c}$	$\epsilon_{b/c}$	$N(\epsilon_{a/c}, \epsilon_{b/c})$
0.3349	0.066185	0.83	0.83	5.3639256E-04
0.3669	0.065025	0.37	0.37	8.9411497E-02
0.4019	0.063635	0.56	0.67	0.2316431
0.4403	0.06205	0.31	0.36	0.1678001
0.4823	0.05872	0.37	0.56	0.1187796
0.5283	0.05335	0.41	0.75	0.2485939
0.5787	0.0477625	0.47	1.00	2.2177355E-02
0.6339	0.042953	0.31	0.73	1.0744088E-02
0.6944	0.0403205	0.31	0.82	5.5961255E-03
0.7607	0	0.31	1.00	5.5961255E-03
0.8333	0			
0.9129	0			
1	0			
1.0954	0			
1.2	0			
1.3145	0			
1.44	0.0403205			
1.5774	0.042953			
1.728	0.0477625			
1.8929	0.05335			
2.0736	0.05872			
2.2715	0.06205			
2.4883	0.063635			
2.7258	0.065025			
2.986	0.066185			

Note. ϵ is the aspect ratio of spheroid particles and $\epsilon_{a/c}$ and $\epsilon_{b/c}$ are two aspect ratios of ellipsoid particles. N is the weight for each particle shape and adds up to 1 in each experiment.

We also notice large differences in the linear polarization state of the scattered light, with the spherical and nonspherical particle classes having opposite signs at most scattering angles (Figure 3d). Lastly, we see large differences in CEXT, both in magnitude and spectral dependence, among the different shapes (Figure 3a). The TX-Ellipsoid has the largest CEXT and DT-Sphere has the smallest CEXT. Due to the way the DT algorithm is designed, the spectral dependence of CEXT rather than the absolute value of CEXT at 0.55 μm matters in the retrieval of AOD at 0.55 μm , and at the other wavelengths (Equations 4–6). The weaker the spectral dependence of CEXT, the weaker will be the spectral dependence of the retrieved AOD and the smaller the magnitude of the retrieved AE. The DT algorithm is particularly sensitive to the spectral dependence of the observations and the LUT values that will be matched to the observations.

Between 0.47 and 0.86 μm , we notice a “spectral dip” appearing near 0.65 μm for the spheroid and ellipsoid shapes from the TX database. We believe that this dip is unphysical, given the small change in refractive indices among these channels (Table 1). Further investigation of the TX database revealed that the dips are caused by significant computational noise in some of the database entries, specifically those representing prolate particles with aspect ratios larger than 2.2 (personal communications with Jiacheng Ding, 5 January 2019). Unfortunately, these artifacts prohibit the use of the TX database for our application because realistic modeling of the spectral dependence is central to the performance of the DT algorithm.

For this reason, we ultimately derive our new optical properties from the database of single scattering properties furnished with the GRASP toolset. The green curve in Figure 3 (GRASP-Sphd) represents the same particle microphysics as TX-Sphd but with the optical properties pulled from the GRASP database. The phase function from GRASP is very similar to those

from TX-Sphd and TX-Ellipd, but the GRASP-derived values of CEXT fall between the spherical and nonspherical particles in the TX database. Most importantly, GRASP does not have the undesired spectral dip. The spectral variation of CEXT from GRASP-Sphd is also smaller than that from DT (Figure 3d), which will impact the retrieved AE as will be shown later.

Besides the shape of the particles, the spectral refractive indices and size distribution are very important in determining the optical properties of dust. There are many laboratory and field measurements of dust optical properties (Dubovik et al., 2002; Espinosa et al., 2018; Formenti et al., 2011; Hess et al., 1998; Levoni et al., 1997; Muñoz et al., 2007; Petzold et al., 2009). Many of these measurements are not directly applicable to our needs because specific locations or samples are not representative of global dust properties, or instrument geometry, wavelength, and calibration differ from the MODIS instrument. For example, refractive indices derived from AERONET stations have to be adjusted when applied to DB algorithm for the Visible Infrared Imaging Radiometer Suite (VIIRS) instrument due to calibration issues (Lee et al., 2017). We tried several dust models from the literature (see above), including the dust models used by Lee et al. (2017), two “dust-like” coarse mode models from the current DT algorithm (replacing the spherical-shaped particles with a mixture of spheroids and ellipsoids), and perturbations of different refractive indices and size distributions based on these models. The tests are somewhat “ad hoc” as only retrieval results can indicate whether the chosen dust model is “good enough” by comparison with AERONET retrievals. Ultimately, the model chosen was similar to the “dust-like” Model-8 of the operational DT algorithm but with a slightly reduced median radius. The new model is named Model-10 in the following description and in Table 1.

3.2. Dust LUT

The LUT calculates the relationship between AOD at 0.55 μm ($\tau_{0.55}$) and TOA reflectance at a given wavelength and geometry

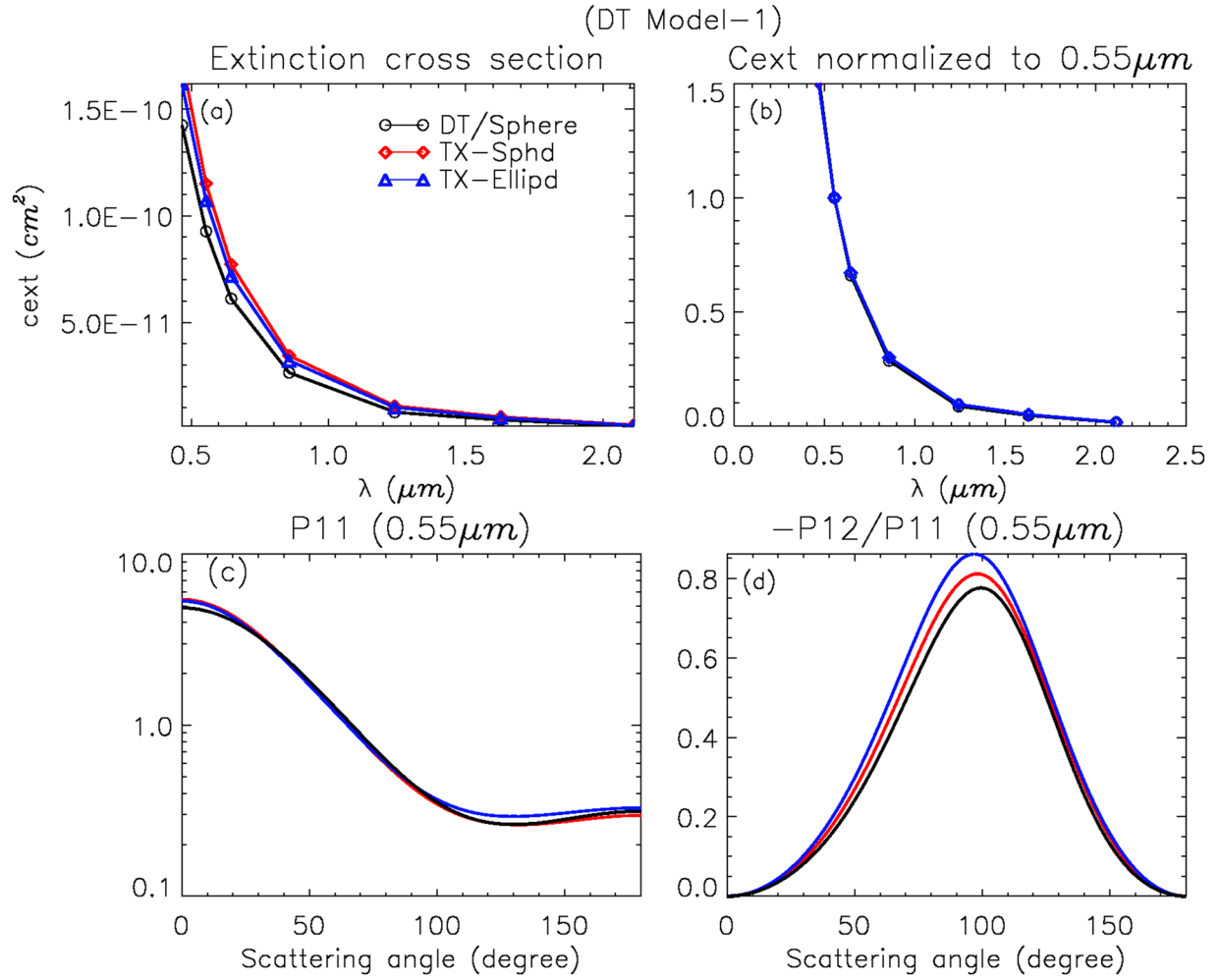


Figure 2. Sensitivity of optical properties to particle shapes for DT Model-1. Optical properties for spheres (DT/Sphere) are obtained from the DT Lookup Table, while a mixture of spheroids (TX-Sphd) and ellipsoids (TX-Ellipd) are obtained from the Texas A&M optical database. (a) Extinction cross section (Cext) and (b) extinction cross section normalized at 0.55 μm are shown for each wavelength used in the DT aerosol retrieval. (c) Phase function at 0.55 μm (P11) and (d) degree of linear polarization at 0.55 μm ($-P12/P11$) are shown as a function of scattering angle.

$$R = f(\tau, \lambda, \theta_0, \theta, \phi, \text{wind}). \quad (3)$$

The value $\tau_{0.55}$ is used as the anchor of the LUT and the RT code seeks the right concentration of particles to satisfy this optical loading. After the ensemble dust absorption and single scattering optical properties are computed, they are used in the AF82 RT model along with wind dependent surface reflectance and the Rayleigh scattering contribution of the atmosphere (dependent on wavelength) to simulate the TOA reflectances of the LUT.

We take CEXT_{λ} , SSA_{λ} , and $\text{P11}_{(\lambda, \theta)}$ to represent the extinction coefficient, SSA, and phase functions, respectively, at wavelength λ and scattering angle θ . For example, $\text{CEXT}_{0.55}$ is the extinction coefficient at 0.55 μm . If N is the total number of scattering particles per cross-sectional area of atmospheric column the following equation holds,

$$N \propto \tau_{0.55} / \text{CEXT}_{0.55}. \quad (4)$$

In the single scattering limit, the Rayleigh-corrected reflectance at the sensor over a dark surface can therefore be approximated by $R'_{(\lambda, \theta)}$ where

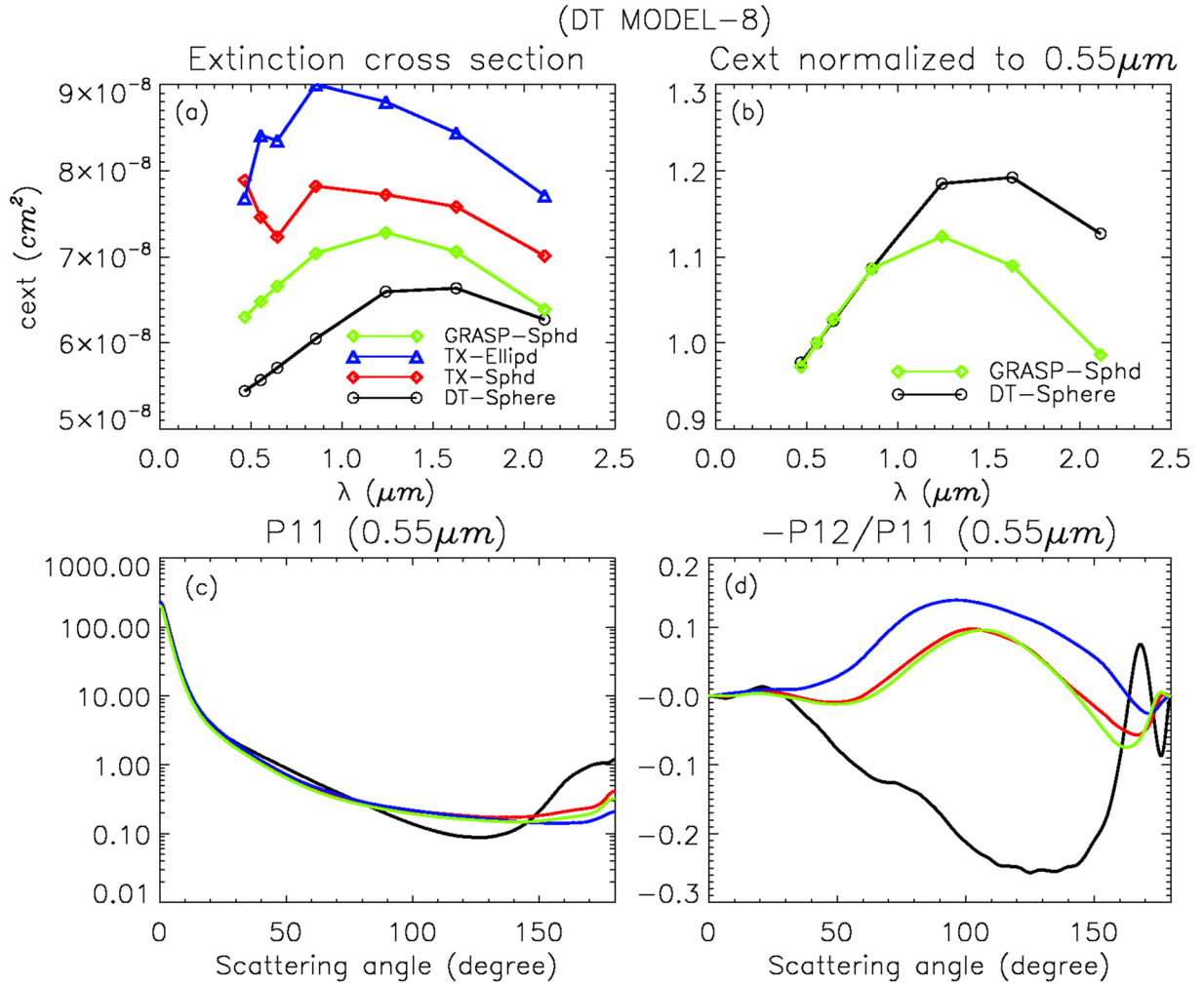


Figure 3. Sensitivity of optical properties to particle shapes for DT Model-8. Optical properties for spheres (DT-Sphere) are obtained from the DT Lookup Table, while a mixture of spheroids (TX-Sphd) and ellipsoids (TX-Ellipd) are obtained from the Texas A&M optical database. The GRASP-Sphd uses the same mixture of spheroid particles as TX-Sphd but computes the optical properties with the GRASP GRASP's internal database instead of the Texas A&M database. (a) Extinction cross section (Cext) and (b) extinction cross section normalized to the value at $0.55\mu\text{m}$ are shown for each wavelength used in the DT aerosol retrieval. (c) Phase function at $0.55\mu\text{m}$ (P11) and (d) degree of linear polarization at $0.55\mu\text{m}$ ($-P12/P11$) are shown as a function of scattering angle.

$$R'_{(\lambda, \theta)} \propto P11_{(\lambda, \theta)} \times SSA_{\lambda} \times CEXT_{\lambda} \times N$$

$$\propto P11_{(\lambda, \theta)} \times SSA_{\lambda} \times \tau_{0.55} \times \frac{CEXT_{\lambda}}{CEXT_{0.55}}. \quad (5)$$

At $0.55\mu\text{m}$, we have

$$R'_{(0.55, \theta)} \propto P11_{(0.55, \theta)} \times SSA_{0.55} \times \tau_{0.55}. \quad (6)$$

From Equation 6 it can be seen that, given a fixed reflectance $R_{(0.55, \theta)}$, the retrieved optical depth is very strongly dependent on our choice of $P11_{(0.55, \theta)}$ and $SSA_{0.55}$. Overestimation of τ could be due to an unrealistically small value of a $P11$ at the observed scattering angle ($P11$ is primarily dependent on the real refractive index, particle size, and shape) or too much absorption (depends primarily on the imaginary refractive index). It is clear that the $P11$ values corresponding to spherical particles used in the current DT algorithm are too high at scattering angles $>140^\circ$ and too low at $80\text{--}140^\circ$ for dust particles, which will likely make the retrieved AOD too low at scattering angle $>140^\circ$ and too high at $80\text{--}140^\circ$. The optical depth at wavelengths other than $0.55\mu\text{m}$ depends on the ratio of extinction coefficient at that wavelength relative to

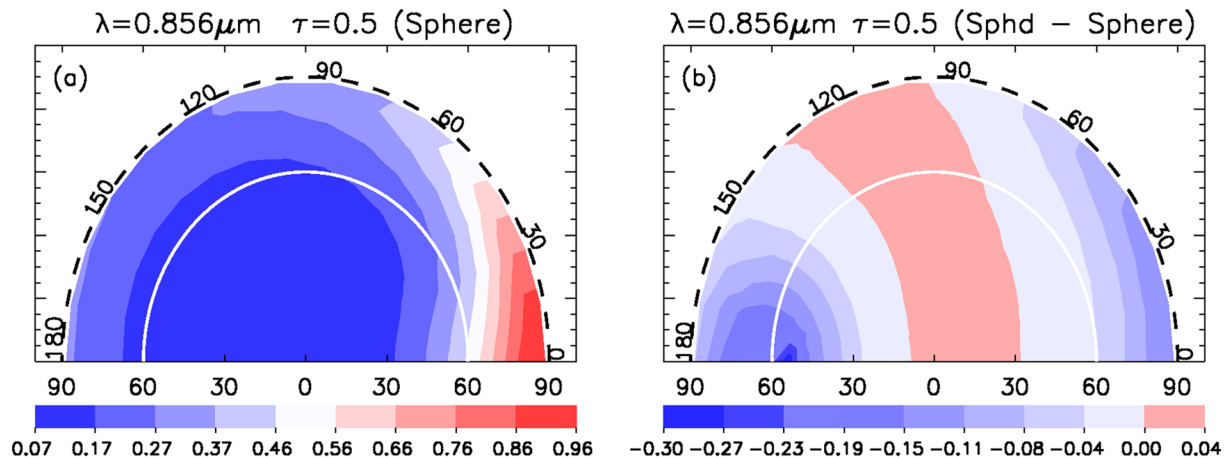


Figure 4. (a) LUT reflectance at $0.856\ \mu\text{m}$ from MODEL-10 with spherical shape, (b) the reflectance difference when the particle shape in MODEL-10 is replaced with mixed spheroids. Reflectance is shown in polar coordinates. The values are plotted for LUT variables of $\tau = 0.5$, surface wind = $10\ \text{m/s}$, and $\text{SZA} = 54^\circ$. View zenith angle ranges from 0 – 90° and relative azimuth angle ranges (counterclockwise) from 0 – 180° . The inner (white) arch encircles the view zenith angle approximates the range of MODIS observations.

extinction at $0.55\ \mu\text{m}$ as we mentioned earlier. The flatter spectral variation of nonspherical particles (Figure 3d) will reduce the spectral steepness of optical depth and reduce the AE.

Figure 4 shows the angular distribution of reflectance, i.e., bidirectional reflectance function (BRF), of spherical particles and the difference produced by changing the particle shapes from spheres to the mixed spheroid shapes of Model-10 at a given optical depth and wavelength. The negative (positive) BRF difference means a higher (lower) optical depth would be required to match the measured reflectances at the different sun-view geometries. We notice that more angular regions present negative difference than positive difference at $0.87\ \mu\text{m}$. Figure 5 (a–f) shows the reflectance differences as a function of scattering angle at six wavelengths. Focusing on the range of scattering angles normally viewed by the MODIS sensor (greater than 80°), we notice that the reflectance differences are negative in all channels for scattering angle greater than 140° . This is expected from Figure 3c, as the phase functions of nonspherical particles at these angles are lower than those of the spherical particles; therefore, we could expect that the retrieved AOD (at $0.55\ \mu\text{m}$) at these angles to be higher under the spheroidal particle assumption. For scattering angles ranging from 80° to 140° , reflectance differences tend to be positive at shorter wavelengths ($0.87\ \mu\text{m}$ and shorter) while negative at the longer wavelengths. Since the reflectance at short wavelengths tends to be larger than at the NIR channels, their overall impact on the retrieved AOD is larger and the retrieved AOD (at $0.55\ \mu\text{m}$) is more likely to be smaller with the spheroid particles. Meanwhile, we could also expect the retrieved spectral AOD to be higher at longer wavelengths and smaller at shorter wavelengths in these scattering angles, resulting in a reduction of AE.

4. Dust Retrieval With Standalone DT-O Algorithm

Before operational implementation, we applied the new dust model (Model-10) to the standalone DT-ocean algorithm, which takes cloud-screened and grid-averaged clear sky reflectance, surface wind, and geometric angles from MxD04 output (e.g., Levy, Remer, & Dubovik, 2007). The standalone algorithm uses the same “guts” as the standard algorithm (spectral observations compared with LUTs iteratively to find the best fit); however, it only does one retrieval per call (Levy & Pinker, 2007). This has a fast turnaround because it does not have to do the time-consuming preprocessing (reading L1B data, gas absorption corrections, performing pixel screening, cloud masking, etc.) and computes only the specified pixels. However, there is a small difference in the retrieved AOD between the operational algorithm and the standalone algorithm due to numerical truncation of the reflectance values in the MxD04 output (Falguni Patadia personal communication, 9 January 2018). The standalone model can be easily modified to only allow the new dust Model-10 as its “chosen” coarse model while the choice of fine model remains as usual for specified dusty pixels. We conducted extensive validation of the new dust model with many dusty granules and collocated AERONET dusty pixels.

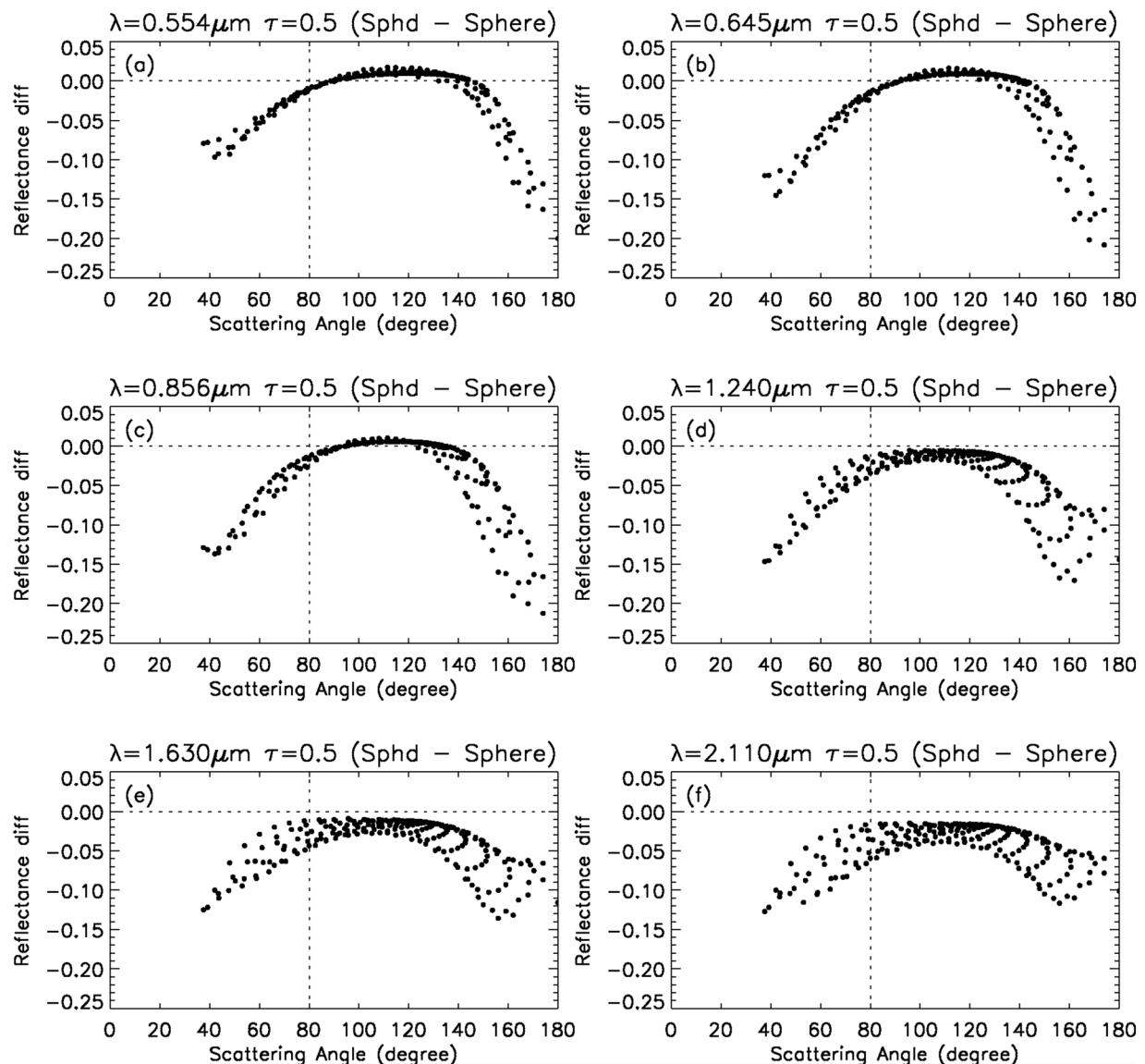


Figure 5. Reflectance difference in six DT aerosol retrieval wavelengths (a-f) of Model-10 with spheroid and spherical particles as a function of scattering angle. Results shown are from with $\tau = 0.5$, surface wind = 10 m/s, and SZA = 54° .

4.1. Evaluation With Dusty Granules

We applied our new dust detection algorithm from Part 1 to hundreds of MODIS granules. When dust pixels are detected, we run the retrieval of that pixel with the modified standalone algorithm that forces the new dust model to be picked. The standalone model retrieves AOD, FMF, AE, and reports retrieval error ERR as in the operational retrieval.

Figure 6 shows one such granule. It illustrates retrieved AOD, FMF, and AE1 from the DT (spherical, Figures 6d, 6g, and 6j), dust-detected (nonspherical, Figures 6e, 6h, and 6k), and the joint (Figures 6f, 6i, and 6l) aerosol retrievals for the MODIS-Aqua granule from 13 July 2005 over the Arabian Peninsula (Figure 6c). The joint retrievals assume original DT retrievals for nondust pixels while replacing the values with standalone model retrievals for identified dusty pixels. Areas of high AOD and low AE, as determined by the C6 DT-retrieval, are shown in Figure 6a, with areas detected as dust, by the detection algorithm (of Part 1) described in section 3, are shown in Figure 6b. Only pixels detected as dust are retrieved using the nonspherical LUT. Included in each panel are contours of scattering angle (Θ). For this granule,

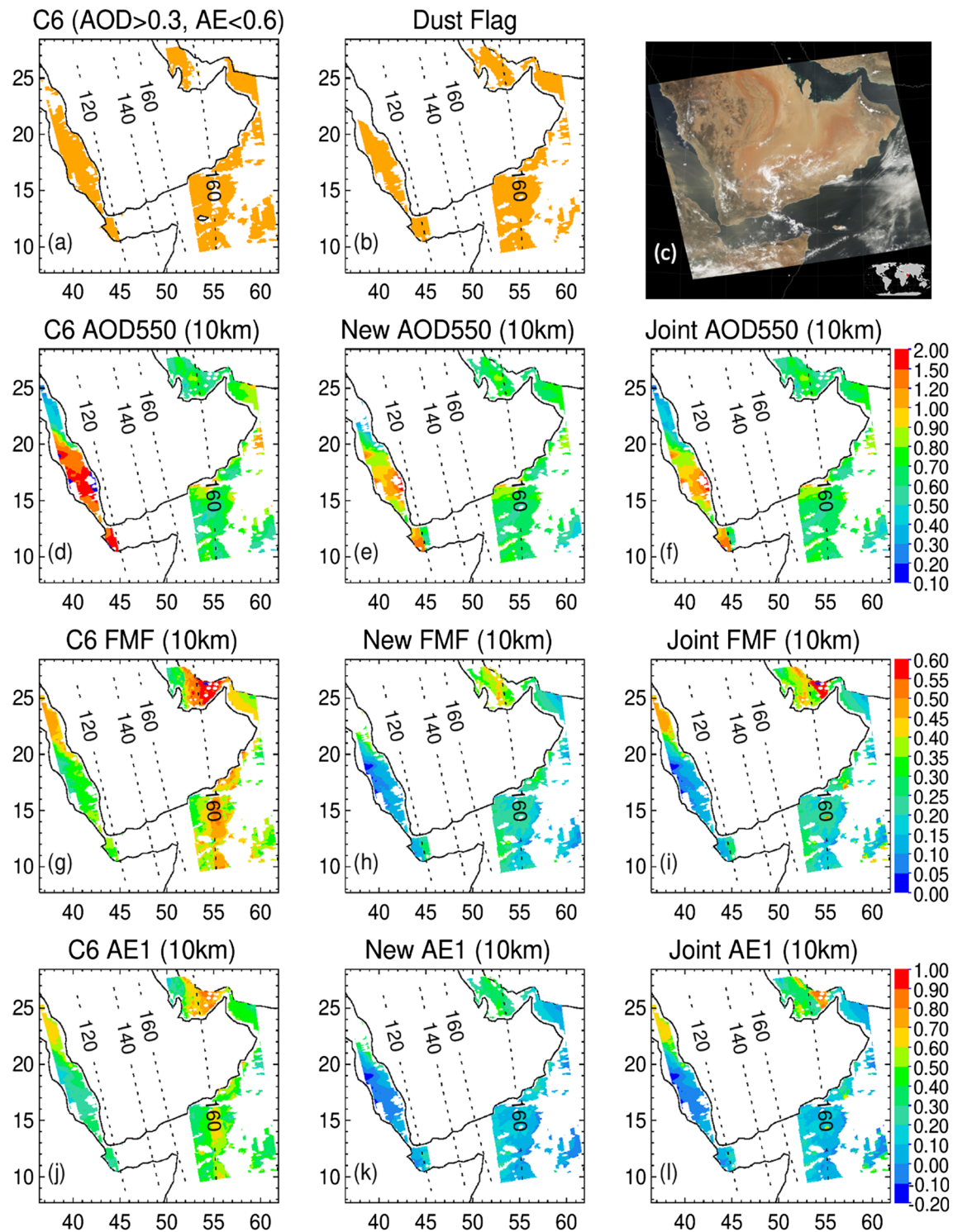


Figure 6. Sample retrievals for the MODIS-Aqua granule from 13 July 2005 over the Arabian Peninsula. (a) Dust areas indicated by high AOD and low AE from the DT C6-retrieval, (b) dust areas detected by the new dust detection algorithm. (d, g, j) AOD, FMF, and AE1 from DT C6 (spherical LUT) retrievals; (e, h, k) retrievals using new dust model (for detected dust pixels only); (f, i, l) joint products. Included in each panel are contours of scattering angle (Θ).

Figure 6f shows that the AOD is significantly reduced by applying the nonspherical model in the Red Sea (smaller Θ), with little or no reduction in the Persian Gulf (larger Θ). FMF and AE are reduced for most dust-detected pixels. Compared with the C6 retrievals, the AOD, AE, and FMF from the nonspherical

Table 3
Time Period and Number of Collocated Dusty Pixels From MODIS and AERONET/MAN Sites

	AERONET		MAN	
	Terra	Aqua	Terra	Aqua
Period	2000–2014	2002–2014	2004–2014	2004–2014
Dusty pixels	18,775	19,387	2,955	3,032

Note. Note that this number reflects the number of individual MODIS retrievals within a 0.3° radius of the AERONET/MAN site and that there is no spatial averaging involved in the collocation.

retrievals show less abrupt changes spatially of the dust plumes over the ocean. But some rough transitions between the nonspherical and spherical retrievals in FMF and AE1, especially in the Persian Gulf, are still visible. It appears that while there might be “dust” everywhere in the region, the magnitude may not be large enough to trigger the dust detection in some areas, hence the rough transition between nonspherical and spherical.

4.2. Evaluation With Collocated AERONET and MAN Observations

To conduct a more quantitative evaluation, we used dust observations from Version 3 AERONET and MAN observation. As mentioned in section 2, we selected a subset of dusty pixels from MODIS-AERONET and

MODIS-MAN collocation data sets based on AERONET/MAN measurements ($AOD > 0.3$, $AE < 0.6$). The total number of AERONET/MAN-identified dusty pixels from each collocation set can be found in Table 3, while the spatial distribution of available dusty pixels (including both AERONET and MAN observations) in each $10^\circ \times 10^\circ$ box is shown in Figure 7. The number of available dusty pixels in each grid box depends on both the number of AERONET/MAN stations and dust frequency in the region as we can see

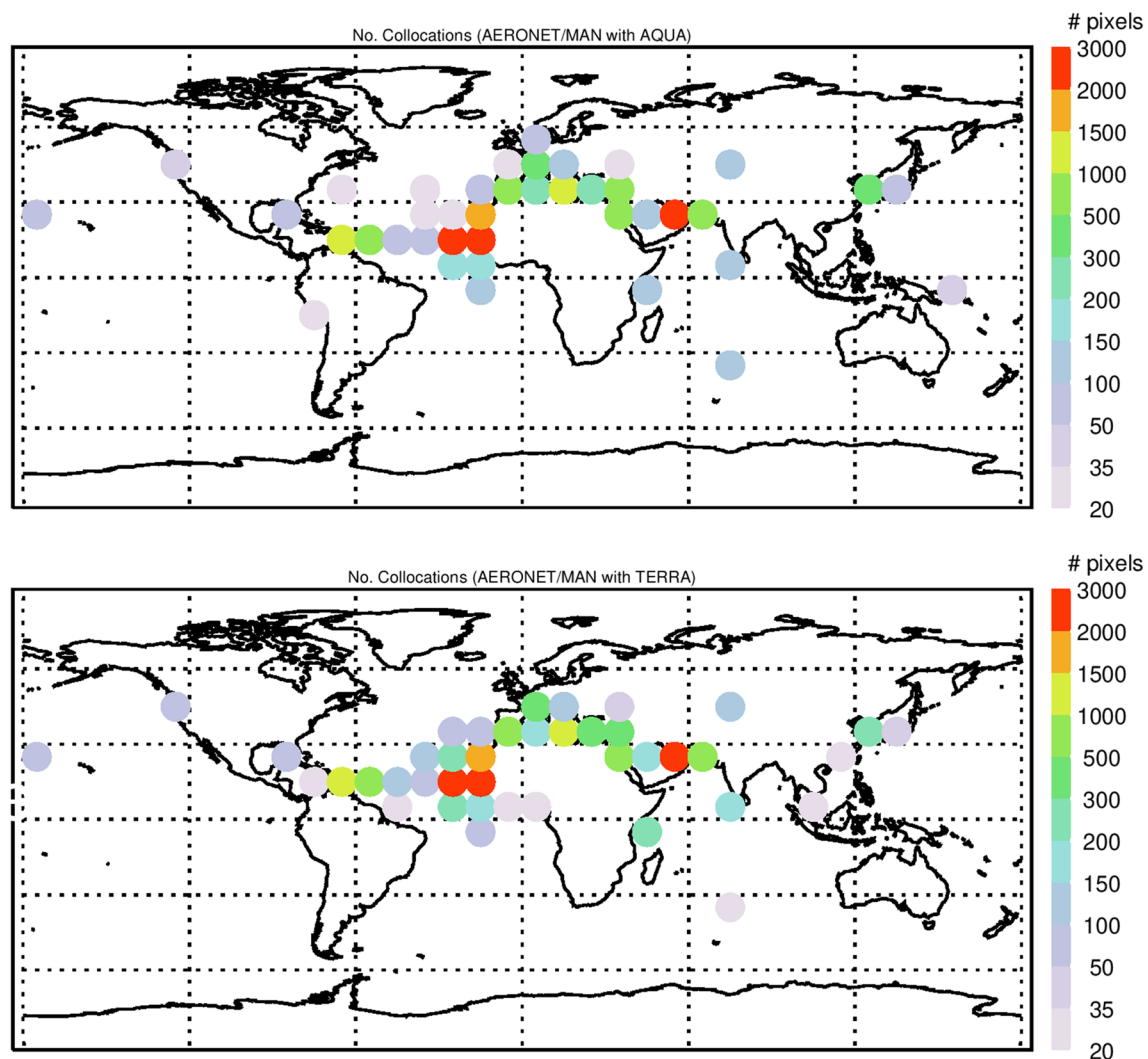


Figure 7. (a) Number of AERONET and MAN dusty pixels collocated with AQUA MODIS in 10° latitude \times 10° longitude boxes. (b) The same as (a) but for collocation with TERRA MODIS.

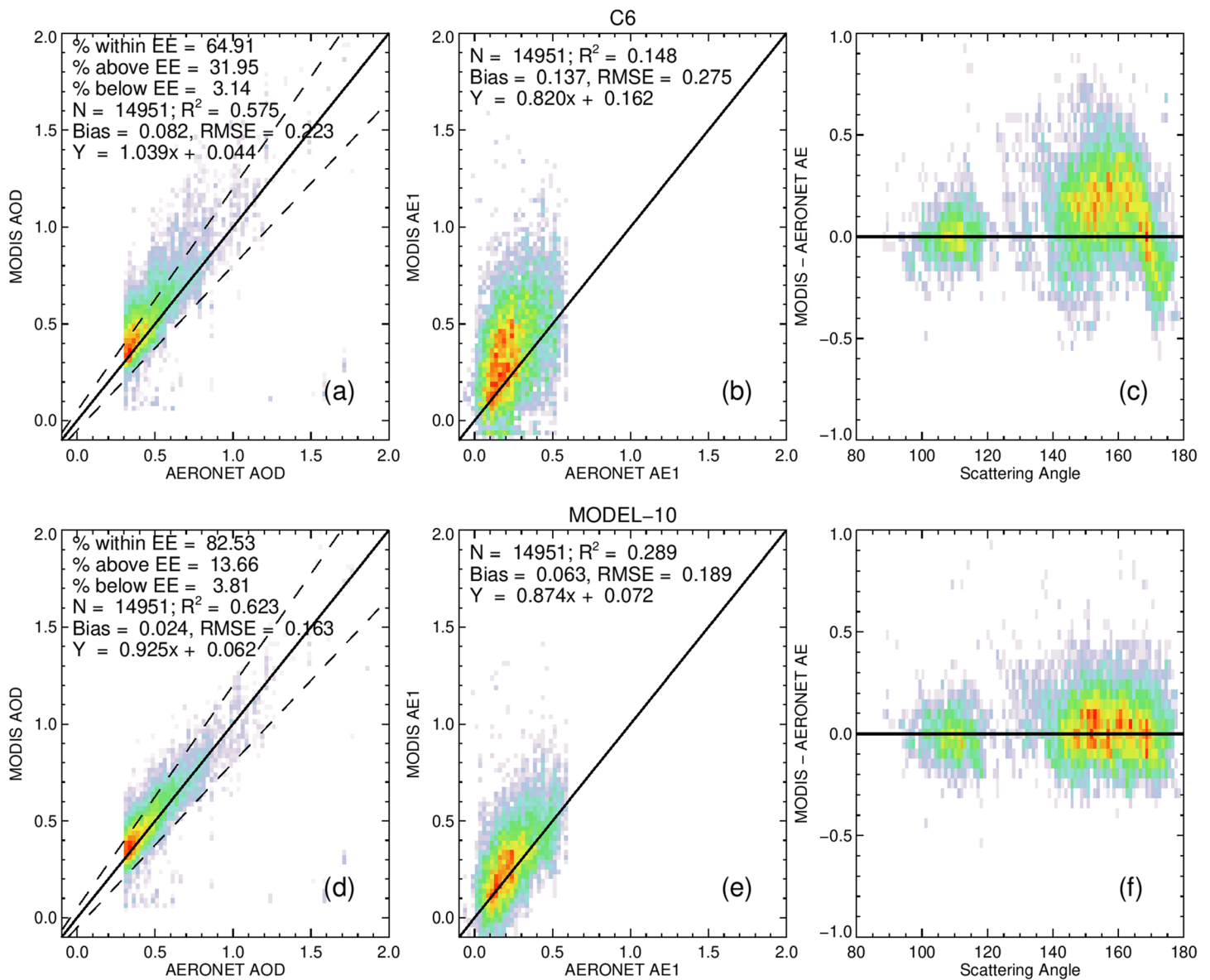


Figure 8. Comparison of dust retrievals from DT version C6 and the new dust model (Model-10) for collocations between AERONET and Aqua dusty pixels. (a) Scatter plot of AOD at $0.55 \mu\text{m}$ between C6 and AERONET; (b) scatter plot of AE1 between C6 and AERONET; (c) MODIS-AERONET AE1 bias as a function of scattering angle. Bottom panel (d, e, f) is the same as the top panel except for the new retrieval with Model-10.

that most validation dusty pixels are concentrated in the tropical north Atlantic, Mediterranean, Persian Gulf, and off the coast of East Asia. For selected MODIS pixels, we apply the retrieval with the modified standalone DT-O algorithm using the new dust LUT.

Figure 8 shows the AOD and AE retrievals with the new dust model as compared with the AERONET measurements for dusty pixels selected from the AQUA-AERONET collocation data set. The C6 retrievals are also displayed for comparison. The new retrieval has significantly reduced the high bias in AOD at $0.55 \mu\text{m}$ from C6 (0.08) to 0.02, percentage of retrievals above expected error (EE) ($\pm (0.05 + 15\%)$) from 32.5% to 13.7% and RMS from 0.23 to 0.17. The percentage within EE has increased from 64.3% to 82.3%. The high bias and RMS in AE from C6 are reduced from 0.13 to 0.06 and 0.27 to 0.19, respectively. It has been a long-standing problem of the DT algorithm that AEs are overly flat (biased high in low AE and low in high AE) (Schutgens et al., 2013). Here we improved AE in the low end for the coarse mode aerosol. In addition, the

Table 4

Performance Metrics of C6 and the New Retrievals With a Nonspherical Dust Model as Compared With AERONET and MAN Measurements for Dusty Pixels From the Collocated Data Sets

		Aqua-AERONET		Terra-AERONET		Aqua-MAN		Terra-MAN	
		C6	New	C6	New	C6	New	C6	New
AOD	bias	0.083	0.024	0.056	0.005	0.025	−0.008	0.055	0.013
	RMS	0.228	0.168	0.191	0.145	0.098	0.073	0.221	0.127
	% within EE	64.35	82.30	69.76	85.61	86.18	95.09	70.66	88.53
	%Above EE	32.48	13.73	25.28	8.66	13.29	3.03	27.17	9.58
	% Below EE	3.17	3.97	4.95	5.73	0.53	1.88	2.17	1.90
AE	bias	0.134	0.060	0.211	0.013	0.229	−0.005	0.073	0.070
	RMS	0.272	0.187	0.324	0.181	0.319	0.179	0.216	0.162

Note. EE is defined as $\pm(0.05 + 15\%)$.

scattering-angle-dependent AE bias in C6 (Figure 8c) is also largely eliminated in the new retrieval (Figure 8f). There are similar improvements from the other collocated data sets, i.e., Terra-AERONET, Terra-MAN, and Aqua-MAN. The detailed statistics are listed in Table 4. In all the cases, the new algorithm reduces bias and RMS of the AOD and AE significantly. The percentages of AOD retrievals within the EE have risen from 64–71% to 82–89% with the exception of Aqua-MAN collocated data set, in that case the EE raised from 86% to 95%. The new retrieval has a mean bias less than 0.02 in AOD and less than 0.07 in AE in all the cases.

To examine whether the new retrievals perform uniformly across the regions, we computed mean dusty pixel AOD and AE for each 10° latitude \times 10° longitude degree box from observations (combined collocated measurements from AERONET and MAN), C6 and the new nonspherical retrievals for Aqua (Figure 9) and Terra (Figure 10). The purpose of this test is to identify regions with consistently poor performance that may need a specific regional dust model. The results indicate that AOD has been improved across the board (more white circles and very few colored circles), with the exception of one location in the central Asia (over Lake Balkhash, Kazakhstan) where both C6 and the new retrievals are consistently higher than the ground observations. For AE, there are overall improvements in matching AERONET/MAN by using the new Model-10 versus the C6 retrieval, but there remain a few spots with persistent large deviations from the ground truth. The discrepancies in AE in these regions, despite the improvement of AOD in these same regions, speak to the sensitivity of AE to subtle changes in retrieved spectral AOD, versus the relatively insensitivity of AOD retrieval at a single wavelength.

Note that here we chose to include AERONET/MAN measurements with smaller aerosol loading to test the nonspherical dust model. The results found in this section are similar if we restrict our dust samples to general dust criteria used in Part 1 (AOD > 0.5, AE < 0.6).

5. Impact on Global Dust Estimates

To estimate the global impact of the new dust retrieval on AOD, AE, and FMF, we applied the new dust retrieval globally to two dust heavy month-long periods: 12 April to 11 May 2011 (Figure 11) and 1 July to 31 July 2011 (Figure 12). The two periods are selected to capture the prevalence of dust in northwest Asia in April (spread over Northwest Pacific, Yu et al., 2012) and the heaviest period of African dust occurring in July (Huang et al., 2010; Yu et al., 2015). Dust pixels were first detected with our new dust detection algorithm (shown in Figure 9 and Figure 10 in Part 1). The new dust model (Model-10) is applied to detected dusty pixels with the standalone algorithm and the retrievals from these dusty pixels are merged with the standard C6 the retrievals. In Figure 11, the left column shows the period-averaged retrievals of AOD at $0.55 \mu\text{m}$, AE1 and FMF from C6 retrievals. The right column shows the difference between C6 retrievals and retrievals using new dust model after dust detection (New-C6). We notice that major impacts occur in the oceans near the major dust sources: Tropical North Atlantic, the Mediterranean and Arabian seas, and the northeast China sea. There is an overall reduction in AOD in the aforementioned regions of about 0.02, and also generally reduced values of AE and FMF. In July 2011, the dust events in northeast Asia subsided to some degree but dust events in West Africa and the Arabian Peninsula increased significantly

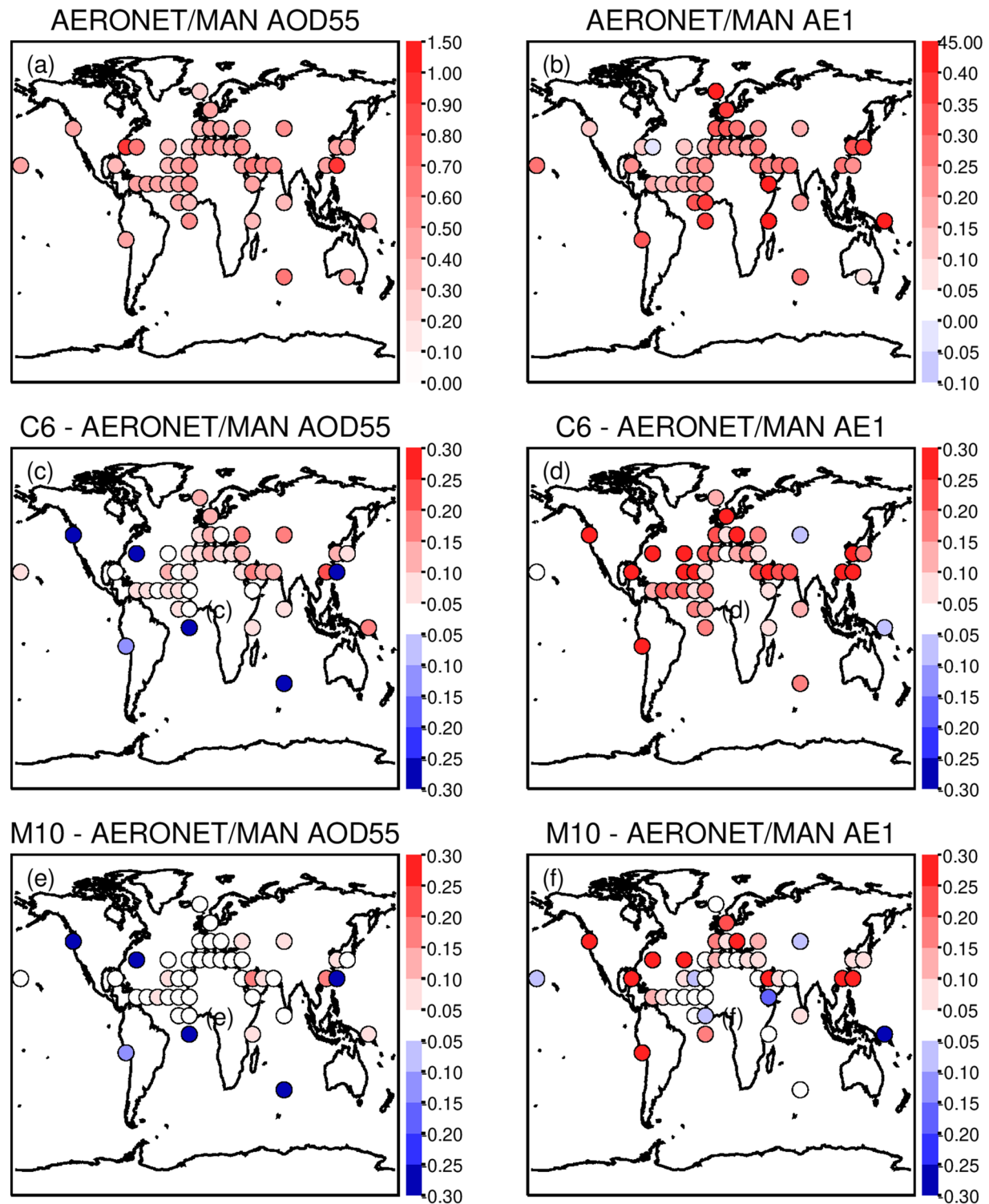


Figure 9. (a) Long-term mean AOD at 0.55 μm and (b) AE1 for 10° latitude \times 10° longitude boxes from AERONET/MAN measurements using the collocated data set with Aqua. Mean bias of C6 DT retrievals as compared to AERONET/MAN in each 10° latitude \times 10° longitude box for (c) AOD at 0.55 μm and (d) AE1. Mean bias of retrievals using the nonspherical dust model (Model-10) as compared to AERONET/MAN in each 10° latitude \times 10° longitude box for (e) AOD at 0.55 μm and (f) AE1.

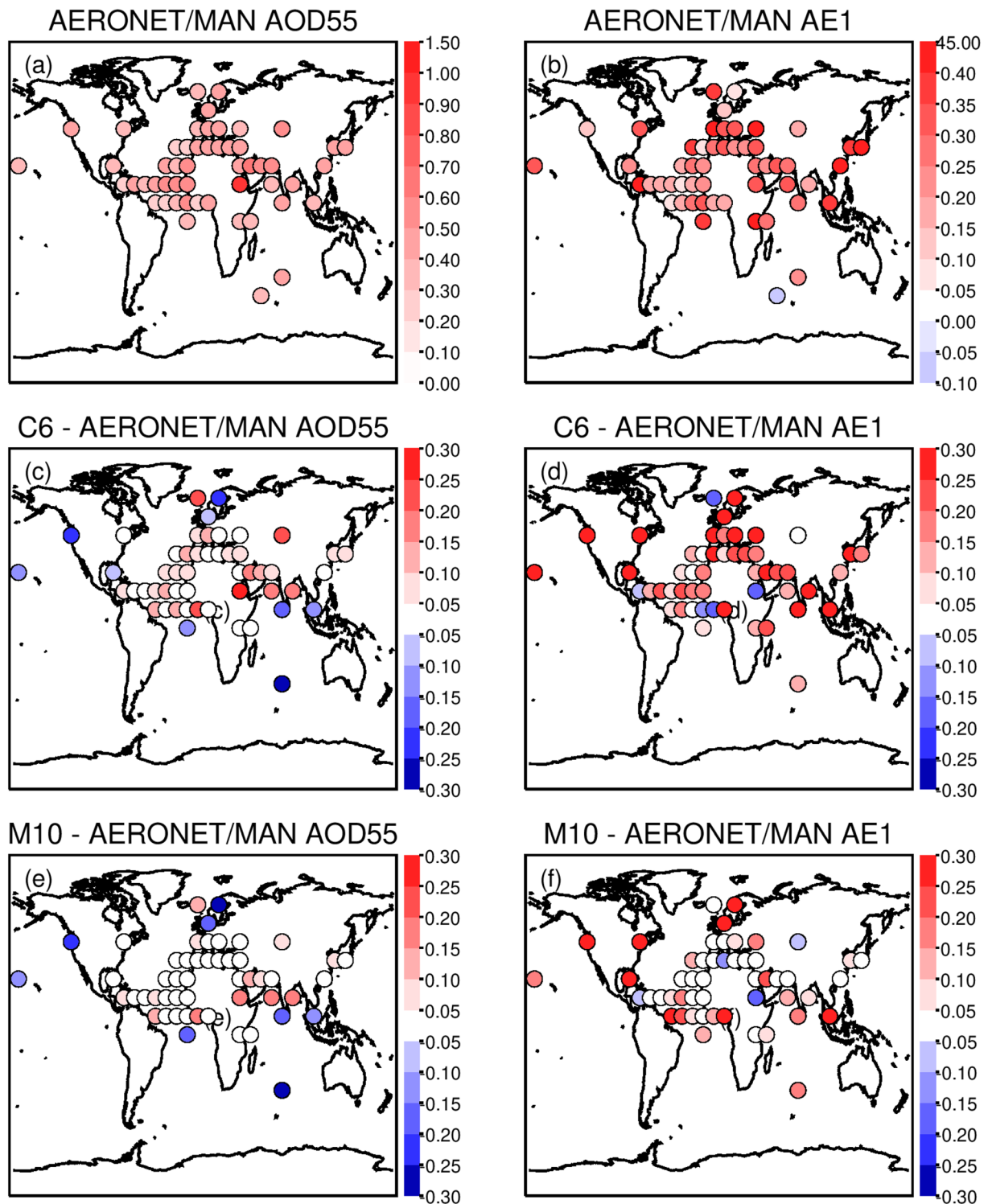


Figure 10. Same as Figure 9 but for the collocated Terra and AERONET/MAN data set.

leading to increased dust occurrences over the North Atlantic Ocean, Persian Gulf, and Arabian Sea (Figure 12). The result is a much decreased monthly averaged AOD (up to 0.2), AE, and FMF in these regions using the new dust model.

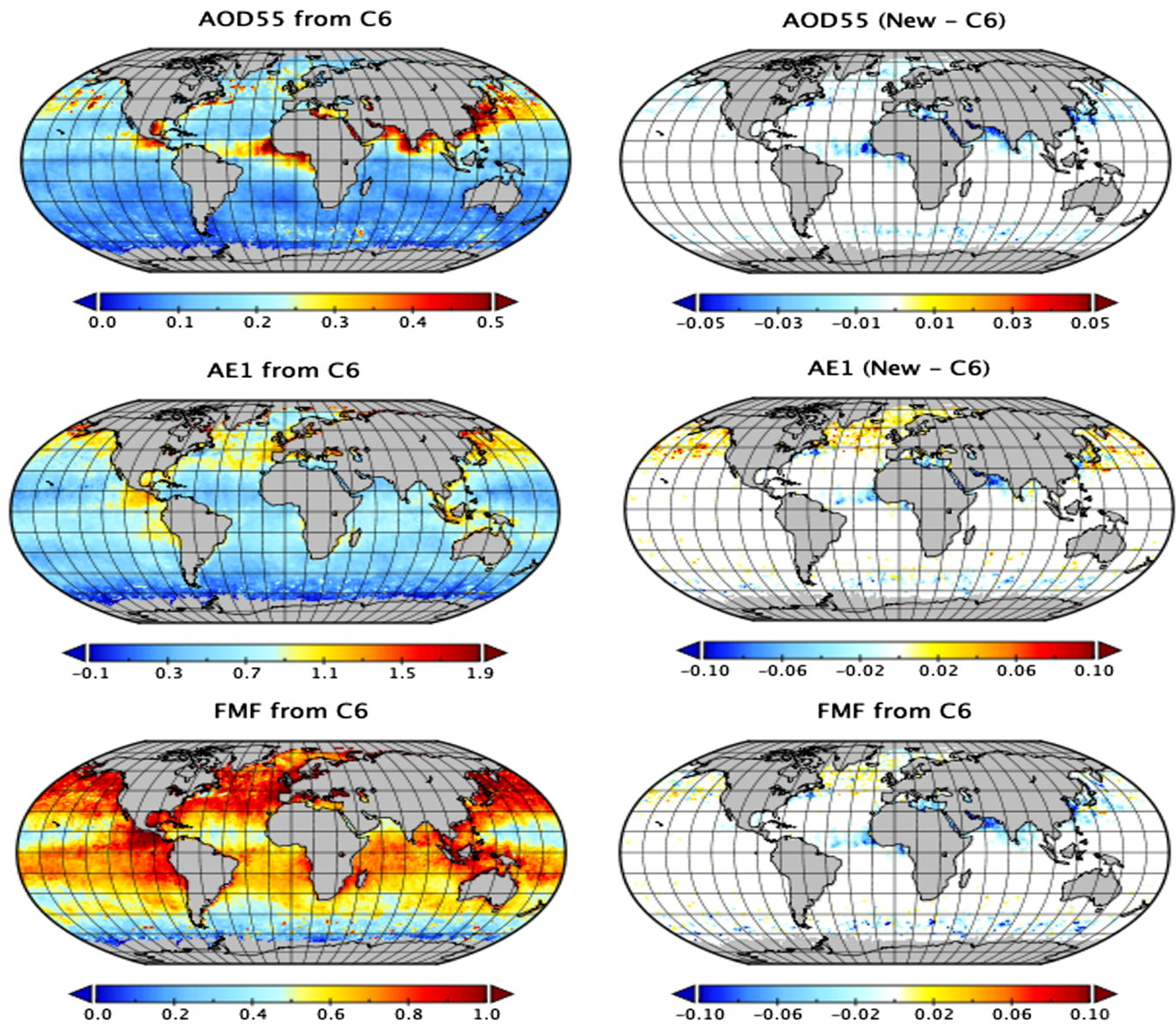


Figure 11. Mean AOD, AE1, and FMF from the C6 DT algorithm during 12 April to 11 May 2011 (left column) and the mean difference between C6 and retrievals with new dust model for detected dusty pixels (New-C6).

The TOA shortwave radiative efficiency in the northern Atlantic region is estimated to be $-49.7 \text{ W m}^{-2} \text{ AOD}^{-1}$ and $-36.5 \text{ W m}^{-2} \text{ AOD}^{-1}$ based on CERES radiation data and AOD from MODIS and CALIOP, respectively (Song et al., 2018). The monthly AOD reduction in these regions could be as much as 0.05–0.2 based on the new retrieval, which means previous estimates of the monthly shortwave radiative forcing in this region could be biased high by about $2\text{--}8 \text{ W m}^{-2}$. The full-scale impact of the new dust retrieval on global and regional AOD and its direct radiative effect will be investigated after the new scheme is implemented operationally and the entire MODIS data record reprocessed.

6. Summary and Discussion

The DT aerosol retrieval is an operational algorithm of the MODIS instruments on the Terra and Aqua satellites that has retrieved spectral AOD over land and ocean since 2000 and 2002, respectively. Recently, DT has been ported to the VIIRS instrument aboard Suomi-NPP (Sawyer et al., 2020), the Advanced Baseline Imager

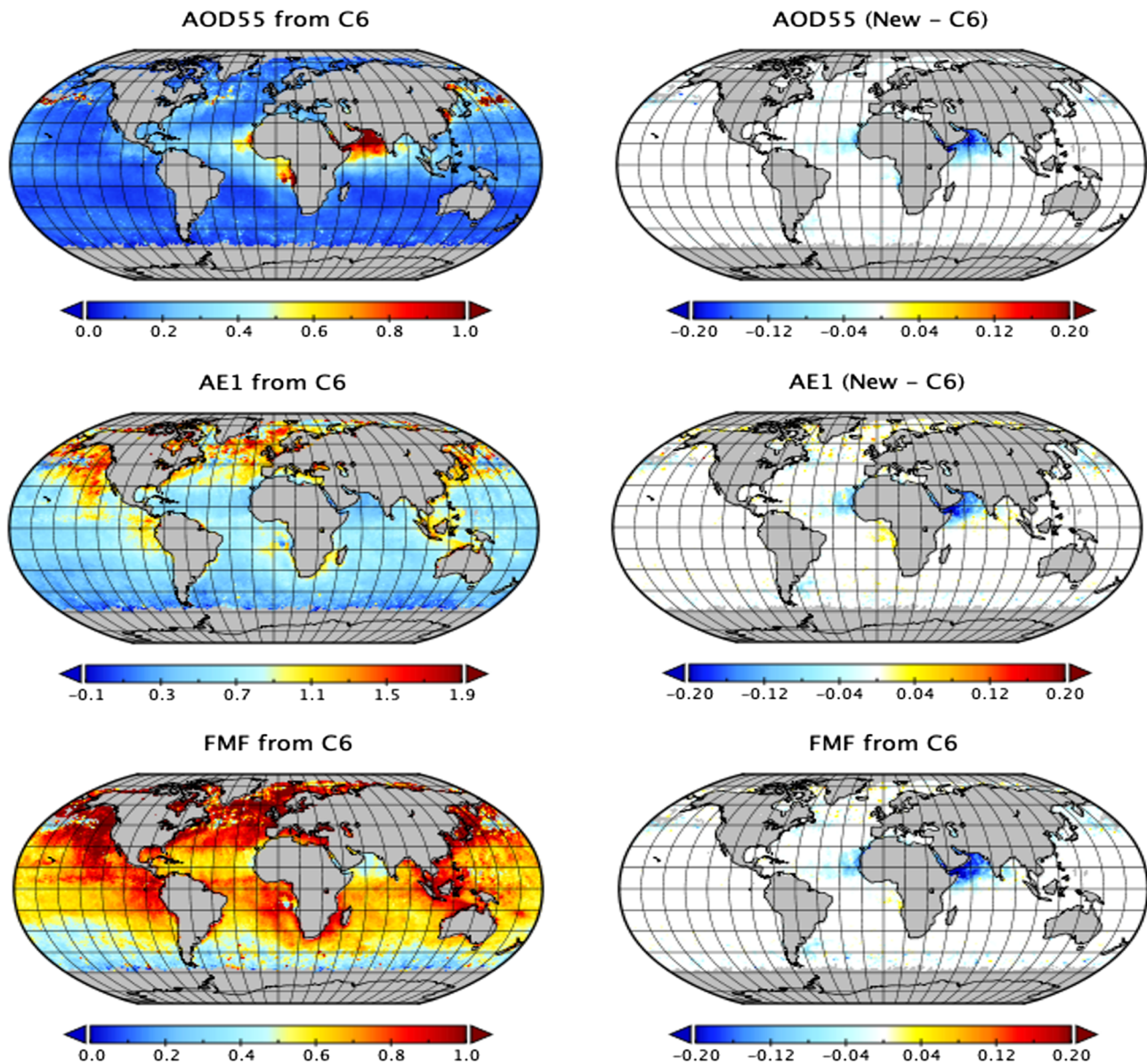


Figure 12. Same as Figure 11 but for July 2011.

(ABI) on GOES-East and GOES-West, the Advanced Himawari Imager (AHI) on Himawari (Gupta et al., 2019). The continued application of the DT algorithm to observations made by the VIIRS instrument, both on the current Suomi-NPP platform and on future NOAA series satellites, will ensure a seamless long-term aerosol climate record. Moreover, the algorithm's application to geosynchronous satellites will enable better monitoring of major aerosol events, such as fires, dust storms, and pollution.

Over the ocean, the DT algorithm is known to produce biased retrievals of AOD, AE, and FMF in pixels containing significant dust aerosol because the current version of DT-O does not have an aerosol model capable of faithfully representing these nonspherical dust particles. In this work, we designed and evaluated a two-step dust aerosol retrieval strategy in which we first detect dust pixels and then utilize a new nonspherical dust model in the retrieval of the detected dusty pixels. The dust detection methodology is reported in Part 1 (Zhou et al., 2020), while this paper describes the construction of the nonspherical dust model and LUT. Sensitivity tests are conducted to determine a dust size and shape distribution and set of refractive

indices that best represent the spectral and angular single-scattering characteristics most relevant to a MODIS-like sensor. In particular, we examined ensemble dust optical properties from the Texas A&M scattering databases for a variety of shapes (sphere, spheroid, and ellipsoid) and size distributions. A final dust model (Model-10) is created that resembles the current Model-8 coarse mode in the DT algorithm but with slightly smaller size and a mixture of spheroidal particles. The optical properties of the new dust model are ultimately computed using the GRASP toolset.

Both the new dust detection and dust models have been tested extensively with manually selected dusty granules. Aerosol measurements from AERONET and MAN are used for quantitative evaluation of the new dust retrievals. Results indicate the new algorithm reduces bias and RMS in the retrieved AOD and AE significantly. The percentages of AOD retrievals within the EE rose from 64–71% to 82–88% with the exception of the Aqua-MAN collocated data set, where the EE increased from 86% to 95%. In all the cases, the new retrieval has lowered the mean bias to less than 0.02 in AOD and less than 0.07 in AE in all the cases. Even more encouragingly, the new retrievals have eliminated all scattering-angle-dependent biases in dusty pixels through the incorporation of phase functions corresponding to nonspherical particles. We estimate that implementation of the new dust retrieval method in the operational algorithm will result in a reduction of monthly mean AOD, AE, and FMF in dust heavy regions such as the northeast Pacific, North Atlantic, Persian Gulf, and Arabian Sea. The magnitude of the resulting reductions in AOD can be as large as 0.05 to 0.2, corresponding to a decrease in TOA shortwave radiative forcing of up to $2\text{--}8\text{ Wm}^{-2}$.

We should point out that neither the dust detection nor the dust aerosol model developed in this work is comprehensive due to the extreme spatial and temporal variability in dust's physical sizes and chemical compositions. The need to rely on manual tuning of detection thresholds hinders the accuracy of the new dust detection algorithm. We expect that machine learning may be able to determine better thresholds and decision trees (Cho et al., 2013; Wang et al., 2020). As for the dust model, using one set of microphysical properties to represent the wide variety of dusts found globally is a significant limitation and the development of regionally applicable dust models may yield significant improvements.

Data Availability Statement

The AERONET/MAN-MODIS collocation data are provided by Yingxi Shi from GSFC. MODIS level 1 (MxD02) and level 2 (MxD04) data are available from NASA DAAC. AERONET products can be obtained from AERONET website (<https://aeronet.gsfc.nasa.gov>).

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